

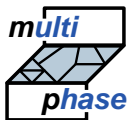
Approximation of dynamic boundary condition: The Allen–Cahn equation

Matthias Liero

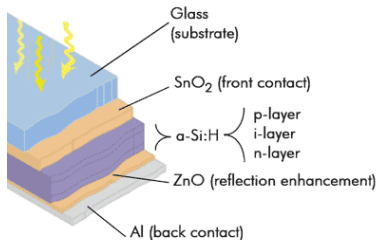
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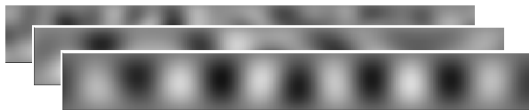


Interfacial dynamics in semiconductor models



Schematic cross-section of thin film a-Si:H photovoltaic cell

Interaction with domain walls in spinodal decomposition



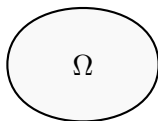
Cahn–Hilliard equation

Sprekels and Wu [2010]: $\Omega \subset \mathbb{R}^d$, $d \geq 2$, with C^2 boundary

$$\begin{aligned}\partial_t u - \Delta u + F'_b(u) &= 0, & \text{in } \Omega \times [0, T] \\ \partial_t \phi - \Delta_{\parallel} \phi + \partial_{\nu} u + \phi + F'_s(\phi) &= 0, & \text{on } \partial\Omega \times [0, T] \\ u &= \phi, & \text{on } \partial\Omega \times [0, T] \\ u(0, \cdot) &= u_0, & \phi(0, \cdot) = \phi_0\end{aligned}$$

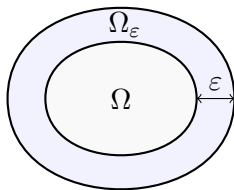
Δ_{\parallel} Laplace–Beltrami operator on $\partial\Omega$

- Cahn–Hilliard equation (R. Kenzler et al. [2001])
- Caginalp system (R. Chill, E. Fasangova and J. Pruess [2006])
- Nonisothermal Allen–Cahn equation (C. Gal and M. Grasselli, A. Miranville [2008])



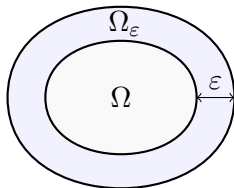
$$r_b \partial_t u - a_b \Delta u + F'_b(u) = 0, \quad \text{in } \Omega \times [0, T]$$

$$\Omega_\varepsilon \stackrel{\text{def}}{=} \{x \in \mathbb{R}^d : \exists(y, \eta) \in \partial\Omega \times (0, \varepsilon), \quad x = y + \eta\nu(y)\}$$



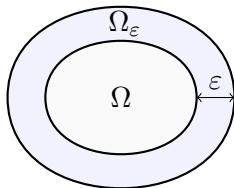
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$$\begin{aligned} r_b \partial_t u - a_b \Delta u + F'_b(u) &= 0, & \text{in } \Omega \times [0, T] \\ r_s \varepsilon^\delta \partial_t \phi - a_s \varepsilon^\alpha \Delta \phi + \varepsilon^\beta F'_s(\phi) &= 0, & \text{in } \Omega_\varepsilon \times [0, T] \\ u &= \phi & \text{on } \partial\Omega \times [0, T] \\ a_b \partial_\nu u - a_s \varepsilon^\alpha \partial_\nu \phi &= 0, & \text{on } \partial\Omega \times [0, T] \\ a_s \varepsilon^\alpha \partial_\nu \phi &= 0, & \text{on } \partial\Omega_\varepsilon \setminus \partial\Omega \times [0, T] \end{aligned}$$

$$\Omega_\varepsilon \stackrel{\text{def}}{=} \{x \in \mathbb{R}^d : \exists (y, \eta) \in \partial\Omega \times (0, \varepsilon), \quad x = y + \eta\nu(y)\}$$



$$\begin{aligned} r_b \partial_t u - a_b \Delta u + F'_b(u) &= 0, & \text{in } \Omega \times [0, T] \\ r_s \varepsilon^{-1} \partial_t \phi - a_s \varepsilon^\alpha \Delta \phi + \varepsilon^{-1} F'_s(\phi) &= 0, & \text{in } \Omega_\varepsilon \times [0, T] \\ u &= \phi & \text{on } \partial\Omega \times [0, T] \\ a_b \partial_\nu u - a_s \varepsilon^\alpha \partial_\nu \phi &= 0, & \text{on } \partial\Omega \times [0, T] \\ a_s \varepsilon^\alpha \partial_\nu \phi &= 0, & \text{on } \partial\Omega_\varepsilon \setminus \partial\Omega \times [0, T] \end{aligned}$$

Today: $\delta = -1, \beta = -1, \alpha \in (-\infty, 1)$

$$r_b \partial_t u - a_b \Delta u + F'_b(u) = 0, \quad \frac{r_s}{\varepsilon} \partial_t \phi - a_s \varepsilon^\alpha \Delta \phi + \frac{1}{\varepsilon} F'_s(\phi) = 0$$

- Spaces

$$\tilde{V}_\varepsilon = \{(u, \phi) \in H^1(\Omega) \times H^1(\Omega_\varepsilon) : u|_{\partial\Omega} = \phi|_{\partial\Omega}\}$$

$$\tilde{H}_\varepsilon = L^2(\Omega) \times L^2(\Omega_\varepsilon)$$

- Energy functionals

$$\tilde{\mathcal{E}}_\varepsilon(u, \phi) = \int_\Omega \frac{a_b}{2} |\nabla u|^2 + F_b(u) dx + \int_{\Omega_\varepsilon} \frac{a_s \varepsilon^\alpha}{2} |\nabla \phi|^2 + \frac{1}{\varepsilon} F_s(\phi) dx$$

- Dissipation functionals

$$\tilde{\mathcal{R}}_\varepsilon(\dot{u}, \dot{\phi}) = \int_\Omega \frac{r_b}{2} |\dot{u}|^2 dx + \frac{1}{\varepsilon} \int_{\Omega_\varepsilon} \frac{r_s}{2} |\dot{\phi}|^2 dx$$

In this setting

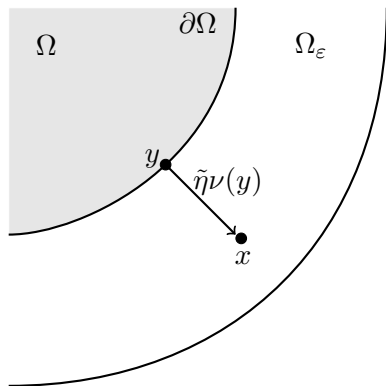
$$r_b \partial_t u - a_b \Delta u + F'_b(u) = 0, \quad \frac{r_s}{\varepsilon} \partial_t \phi - a_s \varepsilon^\alpha \Delta \phi + \frac{1}{\varepsilon} F'_s(\phi) = 0$$

is formally equivalent to

$$0 = D\tilde{\mathcal{R}}_\varepsilon(\dot{u}, \dot{\phi}) + D\tilde{\mathcal{E}}_\varepsilon(u, \phi) \quad (\text{Force balance})$$

Basic existence theory yields solutions

$$(u_\varepsilon, \phi_\varepsilon) \in L^\infty(0, T; \tilde{V}_\varepsilon) \cap H^1(0, T; \tilde{H}_\varepsilon)$$



- Change of coordinates

$$x = y + \tilde{\eta}\nu(y)$$

$$x \mapsto (y, \tilde{\eta}) \in \partial\Omega \times (0, \varepsilon)$$

- Scaling

$$\tilde{\eta} = \varepsilon\eta, \quad \eta \in (0, 1)$$

- Mapping

$$L^2(\Omega_\varepsilon) \longrightarrow L^2(\partial\Omega \times (0, 1))$$

- Spaces

$$\tilde{V}_\varepsilon \longrightarrow V = \{(u, \phi) \in H^1(\Omega) \times H^1(\partial\Omega \times (0, 1)) : u|_{\partial\Omega} = \phi|_{\eta=0}\}$$

$$\tilde{H}_\varepsilon \longrightarrow H = L^2(\Omega) \times L^2(\partial\Omega \times (0, 1))$$

- Energy functionals

$$\begin{aligned}\mathcal{E}_\varepsilon(u, \phi) &= \int_{\Omega} \frac{a_b}{2} |\nabla u|^2 + F_b(u) dx \\ &+ \int_{\partial\Omega} \int_0^1 \left[\frac{a_s \varepsilon^{\alpha+1}}{2} (|B_\varepsilon \nabla_{\parallel} \phi|^2 + \frac{1}{\varepsilon^2} |\partial_\eta \phi|^2) + F_s(\phi) \right] J_\varepsilon d\eta da\end{aligned}$$

- Dissipation functionals

$$\mathcal{R}_\varepsilon(\dot{u}, \dot{\phi}) = \int_{\Omega} \frac{r_b}{2} |\dot{u}|^2 dx + \int_{\partial\Omega} \int_0^1 \frac{r_s}{2} |\dot{\phi}|^2 J_\varepsilon d\eta da$$

- Transformed solutions solve

$$D\mathcal{E}_\varepsilon(u_\varepsilon, \phi_\varepsilon) + D\mathcal{R}_\varepsilon(\dot{u}_\varepsilon, \dot{\phi}_\varepsilon) = 0$$

- Energy functionals

$$\begin{aligned}\mathcal{E}_\varepsilon(u, \phi) = & \mathcal{E}_b(u) \\ & + \int_{\partial\Omega} \int_0^1 \left[\frac{a_s \varepsilon^{\alpha+1}}{2} (|B_\varepsilon \nabla_{\parallel} \phi|^2 + \frac{1}{\varepsilon^2} |\partial_\eta \phi|^2) + F_s(\phi) \right] J_\varepsilon d\eta da\end{aligned}$$

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- Transformed solutions solve

$$D\mathcal{E}_\varepsilon(u_\varepsilon, \phi_\varepsilon) + D\mathcal{R}_\varepsilon(\dot{u}_\varepsilon, \dot{\phi}_\varepsilon) = 0$$

$$\mathcal{E}_\varepsilon(u_\varepsilon(t), \phi_\varepsilon(t)) + \int_0^t \mathcal{R}_\varepsilon(\dot{u}_\varepsilon, \dot{\phi}_\varepsilon) d\tau \leq \mathcal{E}_\varepsilon(u_\varepsilon(0), \phi_\varepsilon(0))$$

- Case $\alpha \leq -1$:

$$\|(u_\varepsilon, \phi_\varepsilon)\|_{L^\infty(0,T;V)} \leq C$$

- Case $\alpha \in (-1, 1)$:

$$\|(u_\varepsilon, \phi_\varepsilon)\|_{L^\infty(0,T;\bar{V})} \leq C$$

with

$$\begin{aligned} \bar{V} := \{ & (u, \phi) \in H^1(\Omega) \times L^2(\partial\Omega \times (0, 1)) : \\ & \partial_\eta \phi \in L^2(\partial\Omega \times (0, 1)), u|_{\partial\Omega} = \phi|_{\eta=0} \} \end{aligned}$$

- Bounded dissipation functional

$$\|(\dot{u}_\varepsilon, \dot{\phi}_\varepsilon)\|_{L^2(0,T;H)} \leq C$$

We have

$$\begin{aligned}(u_\varepsilon, \phi_\varepsilon) &\overset{*}{\rightharpoonup} (u, \phi) \text{ in } L^\infty(0, T; V) \text{ (resp. in } L^\infty(0, T; \bar{V})) \\ (u_\varepsilon, \phi_\varepsilon) &\rightharpoonup (u, \phi) \text{ in } H^1(0, T; H)\end{aligned}$$

In particular

$$\implies \forall t \in [0, T] : z_\varepsilon(t) \rightharpoonup z(t) \text{ in } H$$

Moreover, $z_\varepsilon \in C_w([0, T], \hat{V})$

$$\implies \forall t \in [0, T] : z_\varepsilon(t) \rightharpoonup z(t) \text{ in } \hat{V}$$

with $\hat{V} = V$ for $\alpha \leq -1$ or $\tilde{V} = \bar{V}$ for $\alpha \in (-1, 1)$

F_b and F_s are λ -convex, i.e., $\exists \lambda \in \mathbb{R}$

$$u \mapsto F_b(u) + \lambda|u|^2, \quad \phi \mapsto F_s(\phi) + \lambda|\phi|^2 \quad \text{are convex,}$$

e.g. $C^{1,1}$ -perturbation of a convex function, double well potential

$$\mathbb{W}(u) = \frac{1}{4}(1 - u^2)^2$$

$$F_b(u - v) \geq F_b(u) - F_b'(u)v - \lambda|v|^2, \quad \forall u, v \in \mathbb{R}$$

Consequence: For $z_\varepsilon = (u_\varepsilon, \phi_\varepsilon)$

$$\mathcal{E}_\varepsilon(z_\varepsilon - \widehat{z}_\varepsilon) + \lambda \|\widehat{z}_\varepsilon\|_H^2 \geq \mathcal{E}_\varepsilon(z_\varepsilon) - \langle D\mathcal{E}_\varepsilon(z_\varepsilon), \widehat{z}_\varepsilon \rangle, \quad \forall \widehat{z}_\varepsilon \in V$$

Hence force balance equivalent to

$$\mathcal{E}_\varepsilon(z_\varepsilon(t) - \widehat{z}_\varepsilon) + \lambda \|\widehat{z}_\varepsilon\|_H^2 \geq \mathcal{E}_\varepsilon(z_\varepsilon(t)) + \langle G_\varepsilon \dot{z}_\varepsilon(t), \widehat{z}_\varepsilon \rangle, \quad \forall \widehat{z}_\varepsilon \in V$$

with $G_\varepsilon = D\mathcal{R}_\varepsilon$

Definition: (Attouch [1984], Braides [2002], Dal Maso [1993])
 $\mathcal{I}_0 : \mathcal{Y} \rightarrow \mathbb{R}_\infty$ **Γ -limit** of sequence $(\mathcal{I}_\varepsilon)_{\varepsilon>0}$ w.r.t. topology τ iff

(i) **liminf estimate:**

$$y_\varepsilon \xrightarrow{\tau} y \implies \mathcal{I}_0(y) \leq \liminf_{\varepsilon \rightarrow 0} \mathcal{I}_\varepsilon(y_\varepsilon)$$

(ii) **limsup estimate** (existence of recovery sequence):

$$\forall y \in \mathcal{Y} \exists (\hat{y}_\varepsilon)_{\varepsilon>0} : y_\varepsilon \xrightarrow{\tau} y \text{ and } \mathcal{I}_0(y) \geq \limsup_{\varepsilon \rightarrow 0} \mathcal{I}_\varepsilon(\hat{y}_\varepsilon).$$

Theorem: $(\mathcal{I}_\varepsilon)_{\varepsilon>0}$ equi-coercive $(\bigcup_{\varepsilon>0} \{\mathcal{I}_\varepsilon \leq E\})$ precompact) and $\mathcal{I}_\varepsilon \xrightarrow{\Gamma} \mathcal{I}_0$; then

$$\min_{\mathcal{Y}} \mathcal{I}_0 = \lim_{\varepsilon \rightarrow 0} \min_{\mathcal{Y}} \mathcal{I}_\varepsilon.$$

$(y_\varepsilon)_{\varepsilon>0}$ minimizers of $\mathcal{I}_\varepsilon \implies$ cluster points y are minimizers of \mathcal{I}_0

Minimization problems:

- Homogenization, Two-scale convergence
- Singular limits (Cahn–Hilliard \longrightarrow Sharp interface)
- Young measure relaxation, penalization

Rate independent systems (Mielke, Roubíček, Stefanelli [2008])

- Two-scale Homogenization (Mielke & Timofte [2007], Hanke [2009])
- Damage to Delamination (Mielke, Thomas, Roubíček *in preparation*)
- Dimension reduction in linearized elastoplasticity (Mielke & L. *in preparation*)

Hamiltonian systems (Mielke [2008])

Gradient flows (Sandier & Serfaty [2004])

$$\mathcal{E}_\varepsilon(u, \phi) = \mathcal{E}_b(u) + \int \int_0^1 \left[\frac{a_s \varepsilon^{\alpha+1}}{2} (|B_\varepsilon \nabla_{\parallel} \phi|^2 + \frac{1}{\varepsilon^2} |\partial_\eta \phi|^2) + F_s(\phi) \right] J_\varepsilon d\eta da$$

(i) Case $\alpha = -1$

$$\mathcal{E}_0 : \begin{cases} V & \rightarrow \mathbb{R}_\infty, \\ (u, \phi) & \mapsto \begin{cases} \mathcal{E}_b(u) + \int_{\partial\Omega} \frac{a_s}{2} |\nabla_{\parallel} \phi|^2 + F_s(\phi) da & \partial_\eta \phi = 0, \\ \infty & \text{otherwise} \end{cases} \end{cases}$$

(ii) Case $\alpha < -1$

$$\mathcal{E}_0 : \begin{cases} V & \rightarrow \mathbb{R}_\infty, \\ (u, \phi) & \mapsto \begin{cases} \mathcal{E}_b(u) + \mathcal{H}^{d-1}(\partial\Omega) F_s(\phi) da & \phi \text{ constant,} \\ \infty & \text{otherwise} \end{cases} \end{cases}$$

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(iii) Case $\alpha \in (-1, 1)$

$$\bar{V} := \{(u, \phi) \in H^1(\Omega) \times L^2(\partial\Omega \times (0, 1)) : \\ \partial_\eta \phi \in L^2(\partial\Omega \times (0, 1)), u|_{\partial\Omega} = \phi|_{\eta=0}\}$$

$$\mathcal{E}_0 : \begin{cases} \bar{V} & \rightarrow \mathbb{R}_\infty, \\ (u, \phi) & \mapsto \begin{cases} \mathcal{E}_b(u) + \int_{\partial\Omega} F_s^{**}(\phi) da & \partial_\eta \phi = 0, \\ \infty & \text{otherwise} \end{cases} \end{cases}$$

In the following: F_s convex

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$$\bar{V} := \{(u, \phi) \in H^1(\Omega) \times L^2(\partial\Omega \times (0, 1)) : \\ \partial_\eta \phi \in L^2(\partial\Omega \times (0, 1)), u|_{\partial\Omega} = \phi|_{\eta=0}\}$$

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In the following: F_s convex

Theorem

Let $z_\varepsilon = (u_\varepsilon, \phi_\varepsilon)$ be solutions w.r.t. $(\mathcal{E}_\varepsilon, \mathcal{R}_\varepsilon)$ with

$$(u_\varepsilon, \phi_\varepsilon) \xrightarrow{*} (u, \phi) \text{ in } L^\infty(0, T; V) \text{ (resp. in } L^\infty(0, T; \bar{V}))$$

$$(u_\varepsilon, \phi_\varepsilon) \rightharpoonup (u, \phi) \text{ in } H^1(0, T; H)$$

then $z = (u, \phi)$ is a solution w.r.t. $(\mathcal{E}_0, \mathcal{R}_0)$, i.e.

$$\mathcal{E}_0(z(t) - \hat{z}) + \lambda \|\hat{z}\|_H^2 \geq \mathcal{E}_0(z(t)) + \langle G_0 \dot{z}(t), \hat{z} \rangle, \quad \forall \hat{z} \in V \text{ (resp. } \bar{V})$$

with $G_0 = D\mathcal{R}_0$. In particular, z solves

$$G_0(\dot{z}) + D\mathcal{E}_0(u, \phi) = 0$$

Formally equivalent to

$$r_b \partial_t u - a_b \Delta u + F'_b(u) = 0, \quad \text{in } \Omega \times [0, T]$$

- Case $\alpha = -1$

$$r_s \partial_t \phi - a_b \Delta_{\parallel} \phi + F'_s(\phi) + a_b \partial_{\nu} u = 0, \quad \text{on } \partial\Omega \times [0, T]$$

- Case $\alpha < -1$

$$\begin{aligned} r_s \partial_t \phi + a_b [\partial_{\nu} u] + F'_s(\phi) &= 0, & \text{in } [0, T] \\ u = \phi = \text{const}, & & \text{on } \partial\Omega \times [0, T] \end{aligned}$$

- $\alpha \in (-1, 1)$

$$r_s \partial_t \phi + a_b \partial_{\nu} u + F'_s(\phi) = 0, \quad \text{on } \partial\Omega \times [0, T]$$

Set $\widehat{z}_\varepsilon = z_\varepsilon - \bar{z}_\varepsilon$, with $\bar{z}_\varepsilon(t)$ recovery sequence for $z(t) - \bar{z}(t)$.

$$\int_0^T \mathcal{E}_\varepsilon(\bar{z}_\varepsilon) + \lambda \|z_\varepsilon - \bar{z}_\varepsilon\|_H^2 dt \geq \int_0^T \mathcal{E}_\varepsilon(z_\varepsilon) + \langle G_\varepsilon \dot{z}_\varepsilon, z_\varepsilon - \bar{z}_\varepsilon \rangle dt$$

2. Integration by parts

$$\int_0^T \langle G_\varepsilon \dot{z}_\varepsilon, z_\varepsilon - \bar{z}_\varepsilon \rangle dt = - \int_0^T \langle G_\varepsilon z_\varepsilon, \dot{\bar{z}}_\varepsilon \rangle dt + \mathcal{R}_\varepsilon(z_\varepsilon(T)) - \mathcal{R}_\varepsilon(z_\varepsilon(0))$$

3. Passing to the limit + Integration by parts

$$\int_0^T \mathcal{E}_0(\bar{z}) + \lambda \|z - \bar{z}\|_H^2 dt \geq \int_0^T \mathcal{E}_0(z) - \langle G_0 \dot{z}, z - \bar{z} \rangle dt$$

4. Sequence $\bar{z}^\delta \rightarrow z - \widehat{z}$ in $L^2(0, T; V_0)$. By arbitrariness of \widehat{z}

$$\mathcal{E}_0(z(t) - \widehat{z}) + \lambda \|\widehat{z}\|_H^2 \geq \mathcal{E}_0(z(t)) + \langle G_0 \dot{z}(t), \widehat{z} \rangle, \quad \forall \widehat{z} \in V/\bar{V}$$

- Different scalings in dissipation potential
- More complex dissipation potentials (Heat equation with entropy as driving potential)
- Interfaces

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Thank You