Some notes about quasi-incompressible fluids

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Models for liquids (rather than for gases) where mass density variation can be neglected but propagation of pressure waves (in contrast to ideally incompressible Navier-Stokes equations) with a dispersion are presented. Advantages may be that the pressure can be well determined even pointwise everywhere (including the boundary). One variant is governed by the system

(1)
$$\varrho \frac{\partial v}{\partial t} + \varrho(v \cdot \nabla)v - \operatorname{div} \nu \nabla v + \frac{\varrho}{2}(\operatorname{div} v)v + \nabla \pi = g \text{ and } \frac{1}{K} \frac{\partial \pi}{\partial t} + \operatorname{div} v = \frac{1}{H} \Delta \frac{\partial \pi}{\partial t},$$

where v is the velocity, π pressure, K is the elastic bulk modulus (in the physical units Pa=J/m³) and H is a "hyper" bulk modulus (in the physical units Pa/m²=J/m⁵). The bulk force $\frac{\varrho}{2}(\operatorname{div} v) v$ was invented by R. Temam to fit the expected energetics, cf. [7]. This model is usually considered rather as an artificial regularization (or numerical stabilization) of the incompressible model only, possibly with $H = \infty$ or ignoring the force $\frac{\varrho}{2}(\operatorname{div} v) v$, cf. e.g. [1, 2] or [4]. Yet, it has a good physical relevancy by itself as the modulus K with the mass density ϱ determines the speed of P- (=pressure) waves for low-frequency range, while $H < \infty$ leads to a micro-inertia like [3] and facilitates a normal dispersion for higher frequencies leading to a lower speed of P-waves.

Existence of a weak solution and certain regularity as well as asymptotics towards the incompressible Navier-Stokes model for $K, H \to \infty$ is at disposal in this model.

Various modifications of the basic model (1) will also be presented. For example, $K = K(\pi)$ leading to a pressure-dependent speed of P-waves, or enhancement towards anomalous dispersion, or a coupling with a phase field χ governed by the Cahn-Hilliard equation (like e.g. [5] for incompressible case) giving rise to a Korteweg-type stress and then $K = K(\chi)$ leading to a phase-dependent P-wave speed in a two-phase flow or, e.g., to a salinity-dependent P-wave speed if χ is concentration of salt in seawater.

Other (likely physical relevant) modifications to be discussed are replacement of the Temam's force by a Bernoulli-type pressure, or making also the second equation in (1) parabolic $\frac{1}{K} \frac{\partial \pi}{\partial t} + \text{div } v = \frac{1}{H} \Delta \pi$ as in [6], or possibly also the fully convective variant, i.e.

$$\varrho \frac{\partial v}{\partial t} + \varrho (v \cdot \nabla) v - \operatorname{div} \left(\nu \nabla v + \left(\frac{\varrho}{2} |v|^2 + \frac{\pi^2}{2K} - \pi \right) \mathbb{I} \right) = g \text{ and } \frac{1}{K} \left(\frac{\partial \pi}{\partial t} + v \cdot \nabla \pi \right) + \operatorname{div} v = \frac{\Delta \pi}{H}$$
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