

Mean Field Models for Self-Gravitating Particles

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We consider in this lecture parabolic-elliptic systems of the form

$$n_t = \nabla \cdot (D_*(\nabla p + n\nabla\varphi)), \quad (1)$$

$$\Delta\varphi = n, \quad (2)$$

which appear in statistical mechanics as hydrodynamical *mean field* models for self-interacting particles, cf. e.g. [4] and other papers by P.-H. CHAVANIS.

Here $n = n(x, t) \geq 0$ is the density function defined for $(x, t) \in \Omega \times \mathbb{R}^+$, $\Omega \subset \mathbb{R}^d$, $\varphi = \varphi(x, t)$ is the Newtonian potential generated by the particles of density n , and the pressure $p \geq 0$ is determined by the density-pressure relation with a sufficiently regular function $p = p(n, \vartheta)$. The parameter $\vartheta > 0$ plays the role of the temperature, and $D_* > 0$ is a diffusion coefficient which may depend on n, ϑ, x, \dots

Such systems can be studied either in the *canonical ensemble* (i.e. the *isothermal* setting), when $\vartheta = \text{const}$ is fixed, or in the *microcanonical ensemble* with a variable temperature: $\vartheta = \vartheta(t)$, and the energy

$$E = \frac{d}{2} \int_{\Omega} p \, dx + \frac{1}{2} \int_{\Omega} n\varphi \, dx = \text{const}, \quad (3)$$

which, for a given n , defines $\vartheta = \vartheta(t)$ in an implicit way.

In this work we consider examples of density-pressure relations

$$p(n, \vartheta) = \vartheta^{d/2+1} P\left(\frac{n}{\vartheta^{d/2}}\right)$$

more general than MAXWELL–BOLTZMANN, FERMI–DIRAC and polytropic.

Interesting questions are:

- existence of entropy functionals and entropy production rates,
- existence of steady states with prescribed mass and temperature, or prescribed mass and energy,
- nonexistence of global in time solutions and their blow up,
- continuation of local in time solutions with polytropic density-pressure relations.

These results have been obtained in collaboration with TADEUSZ NADZIEJA (Zielona Góra), PHILIPPE LAURENÇOT (Toulouse), and ROBERT STAŃCZY (Łódź and Wrocław).

References

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