



On a 3D model for shape-memory alloys
Dedicated to the memory of M. Schatzman

Adrien Petrov

Weierstraß-Institute für Angewandte Analysis und Stochastik

joint work with A. Mielke, L. Paoli and U. Stefanelli

DFG Research Center MATHEON
Mathematics for key technologies

W I A S
Weierstraß-Institut für Angewandte Analysis und Stochastik

Banff, 2nd September 2010



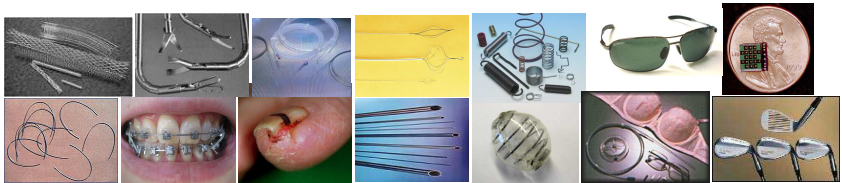
- 1 Background of SMAs
- 2 Background on the Energetic Formulation
- 3 The Souza-Auricchio model for SMAs
- 4 Error estimates for space-time discretizations
- 5 Numerics
- 6 Conclusion



- 1 Background of SMAs
- 2 Background on the Energetic Formulation
- 3 The Souza-Auricchio model for SMAs
- 4 Error estimates for space-time discretizations
- 5 Numerics
- 6 Conclusion



Shape-Memory Alloys are used today in real life:

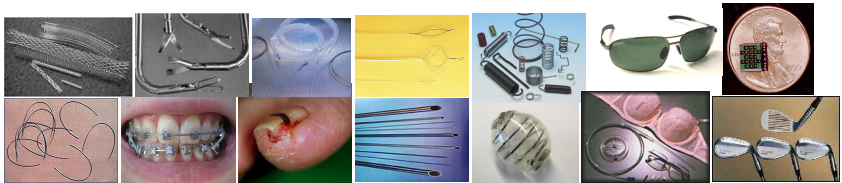


WHY?

¹Souza, Mamiya, Zouain. *Europ. J. Mech., A/Solids*, 1998.



Shape-Memory Alloys are used today in real life:



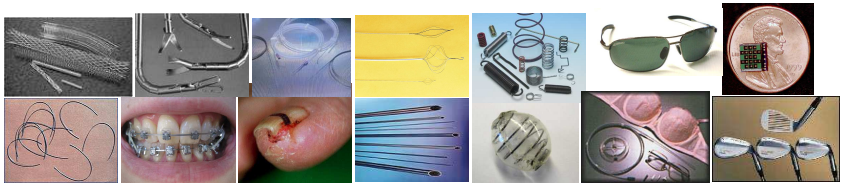
WHY? Because they have the following properties:

- ▷ shape memory under heating and cooling
- ▷ superelastic properties under mechanical loading
- ▷ hysteretic behavior for damping of vibrations

¹Souza, Mamiya, Zouain. *Europ. J. Mech., A/Solids*, 1998.



Shape-Memory Alloys are used today in real life:



WHY? Because they have the following properties:

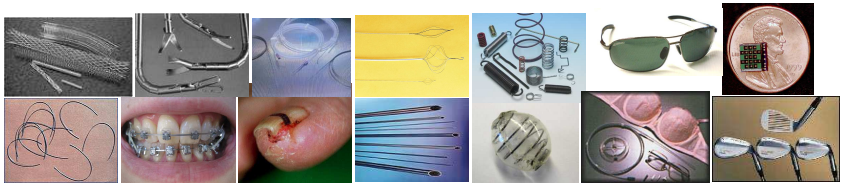
- ▷ shape memory under heating and cooling
- ▷ superelastic properties under mechanical loading
- ▷ hysteretic behavior for damping of vibrations

AIM: Find good mathematical models (analysis and numerics)

¹Souza, Mamiya, Zouain. *Europ. J. Mech., A/Solids*, 1998.



Shape-Memory Alloys are used today in real life:



WHY? Because they have the following properties:

- ▷ shape memory under heating and cooling
- ▷ superelastic properties under mechanical loading
- ▷ hysteretic behavior for damping of vibrations

AIM: Find good mathematical models (analysis and numerics)

⇒ Souza-Auricchio model for shape-memory alloys¹

¹Souza, Mamiya, Zouain. *Europ. J. Mech., A/Solids*, 1998.



- 1 Background of SMAs
- 2 Background on the Energetic Formulation**
- 3 The Souza-Auricchio model for SMAs
- 4 Error estimates for space-time discretizations
- 5 Numerics
- 6 Conclusion



State variables

$u : \Omega \rightarrow \mathbb{R}^d$ displacement

$z : \Omega \rightarrow Z$ phase indicator

Applied field

$\ell_{\text{appl}} : [0, T] \rightarrow \mathcal{X}'$ mechanical loading

Energy functional: $\mathcal{E}(t, u, z) = \int_{\Omega} W(\nabla u, z, \nabla z) dx - \langle \ell_{\text{appl}}(t), u \rangle$

Dissipation functional: $\mathcal{D}(z_1, z_2) = \int_{\Omega} D(x, z_1(x), z_2(x)) dx$

²Mielke, Theil. *Nonl. Diff. Eqns. Appl.*, 2004.



State variables

$u : \Omega \rightarrow \mathbb{R}^d$ displacement
 $z : \Omega \rightarrow Z$ phase indicator

Applied field

$\ell_{\text{appl}} : [0, T] \rightarrow \mathcal{X}'$ mechanical loading

Energy functional: $\mathcal{E}(t, u, z) = \int_{\Omega} W(\nabla u, z, \nabla z) dx - \langle \ell_{\text{appl}}(t), u \rangle$

Dissipation functional: $\mathcal{D}(z_1, z_2) = \int_{\Omega} D(x, z_1(x), z_2(x)) dx$

$(u, z) : [0, T] \rightarrow \mathcal{U} \times \mathcal{Z}$ is called **energetic solution**², if

(S) $\mathcal{E}(t, u(t), z(t)) \leq \mathcal{E}(t, \tilde{u}, \tilde{z}) + \mathcal{D}(z(t), \tilde{z})$ for all $(\tilde{u}, \tilde{z}) \in \mathcal{U} \times \mathcal{Z}$

(E) $\mathcal{E}(t, u(t), z(t)) + \text{Var}_{\mathcal{D}}(z; [0, t]) = \mathcal{E}(0, u_0, z_0) + \int_0^t \partial_s \mathcal{E}(\cdot, u, z) ds$

²Mielke, Theil. *Nonl. Diff. Eqns. Appl.*, 2004.

**State variables**

$u : \Omega \rightarrow \mathbb{R}^d$ displacement
 $z : \Omega \rightarrow \mathcal{Z}$ phase indicator

Applied field

$\ell_{\text{appl}} : [0, T] \rightarrow \mathcal{X}'$ mechanical loading

Energy functional: $\mathcal{E}(t, u, z) = \int_{\Omega} W(\nabla u, z, \nabla z) dx - \langle \ell_{\text{appl}}(t), u \rangle$

Dissipation functional: $\mathcal{D}(z_1, z_2) = \int_{\Omega} D(x, z_1(x), z_2(x)) dx$

$(u, z) : [0, T] \rightarrow \mathcal{U} \times \mathcal{Z}$ is called **energetic solution**², if

(S) $\mathcal{E}(t, u(t), z(t)) \leq \mathcal{E}(t, \tilde{u}, \tilde{z}) + \mathcal{D}(z(t), \tilde{z})$ for all $(\tilde{u}, \tilde{z}) \in \mathcal{U} \times \mathcal{Z}$

(E) $\mathcal{E}(t, u(t), z(t)) + \text{Var}_{\mathcal{D}}(z; [0, t]) = \mathcal{E}(0, u_0, z_0) + \int_0^t \partial_s \mathcal{E}(\cdot, u, z) ds$

If $\mathcal{D}(z_1, z_2) = \Psi(z_2 - z_1)$ and $\mathcal{E}(t, \cdot) : \mathcal{U} \times \mathcal{Z} \rightarrow \mathbb{R}_{\infty}$ convex, then

(S)&(E) $\iff \begin{cases} 0 \in \partial_u \mathcal{E}(t, u, z) & \text{elastic equilibrium} \\ 0 \in \partial \Psi(\dot{z}) + \partial_z \mathcal{E}(t, u, z) & \text{flow rule} \end{cases}$

²Mielke, Theil. *Nonl. Diff. Eqns. Appl.*, 2004.



- 1 Background of SMAs
- 2 Background on the Energetic Formulation
- 3 The Souza-Auricchio model for SMAs**
- 4 Error estimates for space-time discretizations
- 5 Numerics
- 6 Conclusion



State variables

$u : \Omega \rightarrow \mathbb{R}^d$ displacement

$z : \Omega \rightarrow \mathbb{R}_{0,\text{sym}}^{d \times d}$ mesoscopic transformation strain

Applied field

$\ell_{\text{appl}} \in C^3([0, T], \mathcal{X}')$ mechanical load

Energy functional: $\mathcal{E}(t, u, z) = \int_{\Omega} W(\mathbf{e}(u), z, \nabla z) dx - \langle \ell_{\text{appl}}(t), u \rangle$

where $W = \frac{1}{2} \mathbb{C}(\mathbf{e}(u) - z) : (\mathbf{e}(u) - z) + H_{\text{SoAu}}(z) + \sigma |\nabla z|^2$

Dissipation functional: $\mathcal{D}(z_1, z_2) = \Psi(z_2 - z_1) = \int_{\Omega} \rho |z_2 - z_1| dx$



State variables

$u : \Omega \rightarrow \mathbb{R}^d$ displacement
 $z : \Omega \rightarrow \mathbb{R}_{0,\text{sym}}^{d \times d}$ mesoscopic transformation strain

Applied field

$\ell_{\text{appl}} \in C^3([0, T], \mathcal{X}')$ mechanical load

Energy functional: $\mathcal{E}(t, u, z) = \int_{\Omega} W(\mathbf{e}(u), z, \nabla z) dx - \langle \ell_{\text{appl}}(t), u \rangle$

where $W = \frac{1}{2} \mathbb{C}(\mathbf{e}(u) - z) : (\mathbf{e}(u) - z) + H_{\text{SoAu}}(z) + \sigma |\nabla z|^2$

Dissipation functional: $\mathcal{D}(z_1, z_2) = \Psi(z_2 - z_1) = \int_{\Omega} \rho |z_2 - z_1| dx$

$\mathcal{U} \stackrel{\text{def}}{=} \{ u \in H^1(\Omega; \mathbb{R}^d) : u = 0 \text{ on } \Gamma_{\text{Dir}} \}$, $\mathcal{Z} \stackrel{\text{def}}{=} H^1(\Omega; \mathbb{R}_{0,\text{sym}}^{d \times d})$, $\mathcal{Q} \stackrel{\text{def}}{=} \mathcal{U} \times \mathcal{Z}$

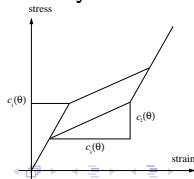
\mathcal{X} : Banach space s.t. $\mathcal{Q} \subset \mathcal{X} \subset \mathcal{Q}'$

$\mathbf{e}(u) = \frac{1}{2}(\nabla u + \nabla u^T)$: infinitesimal strain

$H_{\text{SoAu}}(z) = c_1 |z| + \frac{c_2}{2} |z|^2 + \chi_{\{|z| \leq c_3\}}(z)$

- ▶ c_1 : activation threshold
- ▶ c_2 : hardening in the martensitic regime
- ▶ c_3 : maximal transformation strain

\mathbb{C} : elasticity tensor



**State variables** $u : \Omega \rightarrow \mathbb{R}^d$ displacement $z : \Omega \rightarrow \mathbb{R}_{0,\text{sym}}^{d \times d}$ mesoscopic transformation strain**Applied field** $\ell_{\text{appl}} \in C^3([0, T], \mathcal{X}')$ mechanical load**Energy functional:** $\mathcal{E}(t, u, z) = \int_{\Omega} W(\mathbf{e}(u), z, \nabla z) dx - \langle \ell_{\text{appl}}(t), u \rangle$ where $W = \frac{1}{2} \mathbb{C}(\mathbf{e}(u) - z) : (\mathbf{e}(u) - z) + H_{\text{SoAu}}(z) + \sigma |\nabla z|^2$ **Dissipation functional:** $\mathcal{D}(z_1, z_2) = \Psi(z_2 - z_1) = \int_{\Omega} \rho |z_2 - z_1| dx$ For the regularized version of H_{SoAu} :

$$H_{\delta}(z, \theta) = c_1 \sqrt{\delta^2 + |z|^2} + \frac{c_2}{2} |z|^2 + \frac{1}{\delta} ((|z| - c_3)_+)^3$$

Theorem (Existence and uniqueness^a)*For all $\delta \geq 0$ there exists a solution of **(S)&(E)**.**For $\delta > 0$ the solutions are unique since $\mathcal{E} \in C^3([0, T]; H^1(\Omega))$.*^aMielke, P. *Adv. Math. Sci. Appl.*, 2007.



Since $\mathcal{E}(t, \cdot) : \mathcal{U} \times \mathcal{Z} \rightarrow \mathbb{R}_\infty$ convex, then

$$\mathbf{(S)\&(E)} \iff \begin{cases} 0 \in \partial_u \mathcal{E}(t, u, z) & \text{elastic equilibrium} \\ 0 \in \partial \Psi(\dot{z}) + \partial_z \mathcal{E}(t, u, z) & \text{flow rule} \end{cases}$$

$$\triangleright \partial_u \mathcal{E}(t, u, z) = -\operatorname{div}(\mathbb{C}(\mathbf{e}(u) - z)) - \ell_{\text{appl}}(t)$$

$$\triangleright \partial_z \mathcal{E}(t, u, z) = -\mathbb{C}(\mathbf{e}(u) - z) + \partial_z H_{\text{SoAu}}(z) - \sigma \Delta z$$



Since $\mathcal{E}(t, \cdot) : \mathcal{U} \times \mathcal{Z} \rightarrow \mathbb{R}_\infty$ convex, then

$$\mathbf{(S)\&(E)} \iff \begin{cases} 0 \in \partial_u \mathcal{E}(t, u, z) & \text{elastic equilibrium} \\ 0 \in \partial \Psi(\dot{z}) + \partial_z \mathcal{E}(t, u, z) & \text{flow rule} \end{cases}$$

$$\triangleright \partial_u \mathcal{E}(t, u, z) = -\operatorname{div}(\mathbb{C}(\mathbf{e}(u) - z)) - \ell_{\text{appl}}(t)$$

$$\triangleright \partial_z \mathcal{E}(t, u, z) = -\mathbb{C}(\mathbf{e}(u) - z) + \partial_z \mathcal{H}_{\text{SoAu}}(z) - \sigma \Delta z$$

Then our problem can be rewritten as follows

$$\mathbf{(DI)} \quad 0 \in \partial \Psi(\dot{q}) + \mathbf{A}q + D_q \mathcal{H}_{\text{SoAu}}(q) - \mathbf{L}(t) \text{ a.e. in } [0, T]$$

$$\triangleright q \stackrel{\text{def}}{=} (u, z)^T, \quad \mathbf{L}(t) \stackrel{\text{def}}{=} (\ell(t), 0)^T$$

$$\triangleright \partial \Psi(\dot{q}) \stackrel{\text{def}}{=} (0, \partial \Psi(\dot{z}))^T, \quad \mathcal{H}_{\text{SoAu}}(q) \stackrel{\text{def}}{=} (0, H_{\text{SoAu}}(z) - \frac{c_2}{2} |z|^2)^T$$

$$\triangleright \mathbf{A} \stackrel{\text{def}}{=} \begin{pmatrix} -\operatorname{div}(\mathbb{C}\mathbf{e}(\cdot)) & \operatorname{div}(\mathbb{C}(\cdot)) \\ -\mathbb{C}(\mathbf{e}(\cdot)) & \mathbb{C}(\cdot) - \sigma \Delta(\cdot) + \frac{c_2}{2} \mathbf{I} \end{pmatrix} \in \operatorname{Lin}(\mathcal{Q}, \mathcal{Q}')$$



Since $\mathcal{E}(t, \cdot) : \mathcal{U} \times \mathcal{Z} \rightarrow \mathbb{R}_\infty$ convex, then

$$\text{(S)\&(E)} \iff \begin{cases} 0 \in \partial_u \mathcal{E}(t, u, z) & \text{elastic equilibrium} \\ 0 \in \partial \Psi(\dot{z}) + \partial_z \mathcal{E}(t, u, z) & \text{flow rule} \end{cases}$$

$$\triangleright \partial_u \mathcal{E}(t, u, z) = -\operatorname{div}(\mathbb{C}(\mathbf{e}(u) - z)) - \ell_{\text{appl}}(t)$$

$$\triangleright \partial_z \mathcal{E}(t, u, z) = -\mathbb{C}(\mathbf{e}(u) - z) + \partial_z H_{\text{SoAu}}(z) - \sigma \Delta z$$

Then our problem can be rewritten as follows

$$\text{(DI)} \quad 0 \in \partial \Psi(\dot{q}(t)) + D_q \mathcal{E}(t, q(t)) \text{ a.e. in } [0, T]$$

Here (DI) is equivalent to the *variational inequality*

$$\text{(VI)} \quad \forall v \in \mathcal{Q} : \langle D_q \mathcal{E}(t, q(t)), v - \dot{q}(t) \rangle_{\mathcal{Q}} + \Psi(v) - \Psi(\dot{q}(t)) \geq 0$$



- 1 Background of SMAs
- 2 Background on the Energetic Formulation
- 3 The Souza-Auricchio model for SMAs
- 4 Error estimates for space-time discretizations**
- 5 Numerics
- 6 Conclusion



Finite-element spaces: $\mathcal{U}_h \subset \mathcal{U}$ and $\mathcal{Z}_h \subset \mathcal{Z}$ with $\mathcal{Q}_h = \mathcal{U}_h \times \mathcal{Z}_h$

Time step: $\tau > 0$

Space-Time Discretization: $q_{\tau,h}^k = (u_{\tau,h}^k, z_{\tau,h}^k)$ and $\hat{q}_h = (\hat{u}_h, \hat{z}_h)$

$$(\text{IP})_{\tau,h} \quad q_{\tau,h}^k \in \arg \min_{\hat{q}_h \in \mathcal{Q}_h} (\mathcal{E}(t_{\tau}^k, \hat{q}_h) + \psi(\hat{z}_h - z_{\tau,h}^{k-1}))$$



Finite-element spaces: $\mathcal{U}_h \subset \mathcal{U}$ and $\mathcal{Z}_h \subset \mathcal{Z}$ with $\mathcal{Q}_h = \mathcal{U}_h \times \mathcal{Z}_h$

Time step: $\tau > 0$

Space-Time Discretization: $q_{\tau,h}^k = (u_{\tau,h}^k, z_{\tau,h}^k)$ and $\hat{q}_h = (\hat{u}_h, \hat{z}_h)$

$$(IP)_{\tau,h} \quad q_{\tau,h}^k \in \arg \min_{\hat{q}_h \in \mathcal{Q}_h} (\mathcal{E}(t_\tau^k, \hat{q}_h) + \psi(\hat{z}_h - z_{\tau,h}^{k-1}))$$

▷ $\mathcal{E} : [0, T] \times \mathcal{Q} \rightarrow \mathbb{R}$: the energy functional

(E1) $\mathcal{E} \in C^3([0, T]; \mathcal{Q})$

(E2) $\mathcal{E}(t, \cdot)$ is κ -uniformly convex

(E3) $\exists c, C > 0 : \forall \hat{q} \in \mathcal{Q} \quad c \|\hat{q}\|_{\mathcal{Q}}^2 - C \leq \mathcal{E}(t, \hat{q}) \leq C \|\hat{q}\|_{\mathcal{Q}}^2 + C$

▷ $\Psi : \mathcal{Q} \rightarrow [0, \infty)$: the dissipation functional

(D1) $\forall \gamma \geq 0, q \in \mathcal{Q} : \Psi(\gamma q) = \gamma \Psi(q)$

(D2) $\forall q_1, q_2 \in \mathcal{Q} : \Psi(q_1 + q_2) \leq \Psi(q_1) + \Psi(q_2)$

(D3) $\exists C_\Psi > 0 \forall \hat{q} \in \mathcal{Q} : \Psi(\hat{q}) \leq C_\Psi \|\hat{q}\|_{\mathcal{Q}}$



Finite-element spaces: $\mathcal{U}_h \subset \mathcal{U}$ and $\mathcal{Z}_h \subset \mathcal{Z}$ with $\mathcal{Q}_h = \mathcal{U}_h \times \mathcal{Z}_h$

Time step: $\tau > 0$

Space-Time Discretization: $q_{\tau,h}^k = (u_{\tau,h}^k, z_{\tau,h}^k)$ and $\hat{q}_h = (\hat{u}_h, \hat{z}_h)$

$$\text{(IP)}_{\tau,h} \quad q_{\tau,h}^k \in \arg \min_{\hat{q}_h \in \mathcal{Q}_h} (\mathcal{E}(t_{\tau}^k, \hat{q}_h) + \psi(\hat{z}_h - z_{\tau,h}^{k-1}))$$

Theorem

For all $h \geq 0$, there exists a unique solution $q_h : [0, T] \rightarrow \mathcal{Q}_h$ to (DI) and there exists $R > 0$ such that $\|q_h(t)\|_{\mathcal{Q}} \leq R$ for all $t \in [0, T]$ and $\|\dot{q}_h(t)\|_{\mathcal{Q}} \leq C(R, \kappa)$ for a.e. $t \in [0, T]$.

Idea of the proof: Use Mielke&Theil'04.

AIM: evaluate $\|q_{\tau,h}(t) - q(t)\|_{\mathcal{Q}}$ by some polynomial function of τ and h



Let $\mathbf{P}_h : \mathcal{Q} \rightarrow \mathcal{Q}_h$ s.t. $\forall p_h \in \mathcal{Q}_h: \langle \mathbf{A}\mathbf{P}_h q, p_h \rangle_{\mathcal{Q}} = \langle \mathbf{A}q, p_h \rangle_{\mathcal{Q}}$

There exist $\alpha \in (0, 1]$ $C_P > 0$ s.t. $\forall h > 0 \forall q \in \mathcal{Q}: \|(\mathbf{P}_h - \mathbf{I})q\|_{\mathcal{X}} \leq C_P h^\alpha \|q\|_{\mathcal{Q}}$

(CC) $\exists C > 0 \forall h > 0 \forall (t, \bar{q}_h) \in [0, T] \times \mathcal{Q}_h \forall w \in \mathcal{Q} \exists v_h \in \mathcal{Q}_h :$
 $\langle \mathbf{D}_q \mathcal{E}(t, \bar{q}_h), v_h - w \rangle_{\mathcal{Q}} + \Psi(v_h - w) \leq C h^\alpha \|w\|_{\mathcal{Q}}$

Theorem

There exist $C > 0$ and $\alpha \in (0, 1]$ s.t. for all $h \in [0, h_0]$ and all $q_h(0) \in \mathcal{S}_h(0)$, $\forall t \in [0, T] \forall \tau \in (0, T] : \|q_{\tau, h}(t) - q(t)\|_{\mathcal{Q}} \leq C (h^{\alpha/2} + \sqrt{\tau} + \|q_h(0) - q(0)\|_{\mathcal{Q}})$

$\mathcal{S}_h(t) \stackrel{\text{def}}{=} \{q_h \in \mathcal{Q}_h \mid \forall \hat{q}_h \in \mathcal{Q}_h : \mathcal{E}(t, q_h) \leq \mathcal{E}(t, \hat{q}_h) + \Psi(\hat{q}_h - q_h)\}$: the stable states

Sketch of proof: Let us remark that

$$\|q_{\tau, h}(t) - q(t)\|_{\mathcal{Q}} \leq \underbrace{\|q_{\tau, h}(t) - q_h(t)\|_{\mathcal{Q}}}_{\leq C\sqrt{\tau} \text{ (cf. Mielke\&Theil'04)}} + \underbrace{\|q_h(t) - q(t)\|_{\mathcal{Q}}}_{\leq C(h^{\alpha/2} + \|q_h(0) - q(0)\|_{\mathcal{Q}})}$$



Define now

$$(G1) \quad \gamma(t) \stackrel{\text{def}}{=} \langle D_q \mathcal{E}(t, q_h) - D_q \mathcal{E}(t, q), q_h - q \rangle_{\mathcal{Q}} \geq \kappa \|q - q_h\|_{\mathcal{Q}}^2$$



Define now

$$(G1) \quad \gamma(t) \stackrel{\text{def}}{=} \langle D_q \mathcal{E}(t, q_h) - D_q \mathcal{E}(t, q), q_h - q \rangle_Q \geq \kappa \|q - q_h\|_Q^2$$

$$\mathcal{E}(t, \cdot) \text{ is } \kappa\text{-unif. convex} \xRightarrow{\text{Ciarlet'82}} \mathcal{E}(t, \hat{q}) \geq \mathcal{E}(t, q) + \langle D_q \mathcal{E}(t, q), \hat{q} - q \rangle_Q + \frac{\kappa}{2} \|q - \hat{q}\|_Q^2$$



Define now

$$(G1) \quad \gamma(t) \stackrel{\text{def}}{=} \langle D_q \mathcal{E}(t, q_h) - D_q \mathcal{E}(t, q), q_h - q \rangle_{\mathcal{Q}} \geq \kappa \|q - q_h\|_{\mathcal{Q}}^2$$

By differentiation, we get

$$\begin{aligned} \dot{\gamma}(t) &= \langle \partial_t D_q \mathcal{E}(t, q_h) - \partial_t D_q \mathcal{E}(t, q), q_h - q \rangle_{\mathcal{Q}} \\ &\quad + \langle D_q \mathcal{E}(t, q) - D_q \mathcal{E}(t, q_h) + D_q^2 \mathcal{E}(t, q_h)[q_h - q], \dot{q}_h \rangle_{\mathcal{Q}} \\ &\quad + \langle D_q \mathcal{E}(t, q_h) - D_q \mathcal{E}(t, q) + D_q^2 \mathcal{E}(t, q)[q - q_h], \dot{q} \rangle_{\mathcal{Q}} \\ &\quad + 2 \underbrace{\langle D_q \mathcal{E}(t, q_h) - D_q \mathcal{E}(t, q), \dot{q}_h - \dot{q} \rangle_{\mathcal{Q}}} \\ &\quad \stackrel{(VI)}{\leq} \langle D_q \mathcal{E}(t, q_h), v_h - \dot{q} \rangle_{\mathcal{Q}} + \Psi(v_h - \dot{q}) \stackrel{(CC)}{\leq} Ch^\alpha \|\dot{q}\|_{\mathcal{Q}} \end{aligned}$$



Define now

$$\text{(G1)} \quad \gamma(t) \stackrel{\text{def}}{=} \langle D_q \mathcal{E}(t, q_h) - D_q \mathcal{E}(t, q), q_h - q \rangle_Q \geq \kappa \|q - q_h\|_Q^2$$

By differentiation, we get

$$\dot{\gamma}(t) \stackrel{\text{(E1)-(E3)}}{\leq} Ch^\alpha \|\dot{q}\|_Q + C(1 + \|\dot{q}\|_Q + \|\dot{q}_h\|_Q) \|q - q_h\|_Q^2 \stackrel{\text{Thm. M\&T'04}}{\leq} C \left(h^\alpha + \frac{\gamma(t)}{\kappa} \right)$$

↓ Int. + Grönwall

$$\gamma(t) \leq (e^{Ct/\kappa} - 1) \kappa h^\alpha + e^{Ct/\kappa} \underbrace{\gamma(0)}_{\leq C \|q(0) - q_h(0)\|_Q^2}$$

↓ (G1)

$$\|q(t) - q_h(t)\|_Q^2 \leq (e^{Ct/\kappa} - 1) h^\alpha + \frac{C e^{Ct/\kappa}}{\kappa} \|q(0) - q_h(0)\|_Q^2$$



Lemma

There exists $C > 0$ such that for all $h \in (0, h_0]$, we have $q_h(0) \in \mathcal{S}_h(0)$ and

$$\|q_h(0) - q(0)\|_{\mathcal{Q}} \leq Ch^{\alpha/2}$$

Idea of the proof: Since we have $D_q \mathcal{E}(0, q(0)) \in \text{Lin}(\mathcal{Q}, \mathcal{X}')$ and $\|D_q \mathcal{E}(0, q(0))\|_{\text{Lin}(\mathcal{Q}, \mathcal{X}')} \leq C_\Psi$ which implies the lemma.

³Mielke, Paoli, P., Stefanelli. *SINUM*, 2010.



Lemma

There exists $C > 0$ such that for all $h \in (0, h_0]$, we have $q_h(0) \in \mathcal{S}_h(0)$ and

$$\|q_h(0) - q(0)\|_{\mathcal{Q}} \leq Ch^{\alpha/2}$$

Idea of the proof: Since we have $D_q \mathcal{E}(0, q(0)) \in \text{Lin}(\mathcal{Q}, \mathcal{X}')$ and $\|D_q \mathcal{E}(0, q(0))\|_{\text{Lin}(\mathcal{Q}, \mathcal{X}')} \leq C_\Psi$ which implies the lemma.

Theorem (Error estimates for space-time discretizations³)

There exists $C > 0$ such that $q(0) \in \mathcal{S}(0)$, we have

$$\forall h \in [0, h_0] \exists q_h(0) \in \mathcal{S}_h(0) \forall t \in [0, T] \forall \tau \in (0, T] :$$

$$\|q_{\tau, h}(t) - q(t)\|_{\mathcal{Q}} \leq C(h^{\alpha/2} + \sqrt{\tau})$$

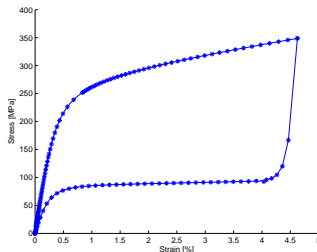
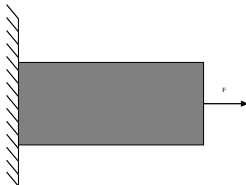
³Mielke, Paoli, P., Stefanelli. *SINUM*, 2010.



- 1 Background of SMAs
- 2 Background on the Energetic Formulation
- 3 The Souza-Auricchio model for SMAs
- 4 Error estimates for space-time discretizations
- 5 Numerics**
- 6 Conclusion



We have developed a 2D FE simulation tools based on some ideas of Carstensen et al.⁴



⁴Carstensen, Klose. *J. Numer. Math.*, 2002.



- 1 Background of SMAs
- 2 Background on the Energetic Formulation
- 3 The Souza-Auricchio model for SMAs
- 4 Error estimates for space-time discretizations
- 5 Numerics
- 6 Conclusion**



- ▷ include **rate-dependent** effects like a heat equation
(work in progress with A. Mielke and L. Paoli)
- ▷ develop the theory to include other **multifunctional materials**
(ferroelectric materials, magnetostrictive materials)



- ▷ include **rate-dependent** effects like a heat equation (work in progress with A. Mielke and L. Paoli)
- ▷ develop the theory to include other **multifunctional materials** (ferroelectric materials, magnetostrictive materials)

Thank you for your attention !

Papers on line: <http://www.wias-berlin.de/people/petrov>