



# Existence and approximation for a 3D model of thermally-induced phase transformations in shape-memory alloys

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joint work with Alexander Mielke and Laetitia Paoli

DFG Research Center MATHEON  
*Mathematics for key technologies*

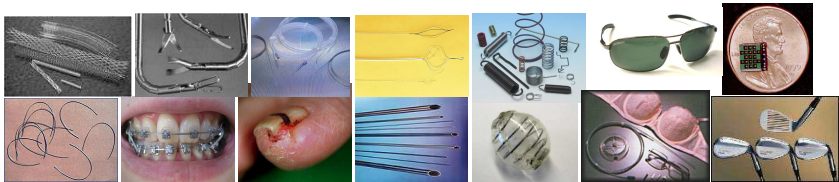
**W I A S**  
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**Shape-memory alloys** are used because of their

- ▷ shape memory under heating and cooling,
- ▷ superelastic properties under mechanical loading,
- ▷ hysteretic behavior for damping of vibrations.

**Applications:** medical, space applications, MEMS...





**AIM:** Model that describes the evolution of **phase mixtures**

Pure phases can be measured experimentally: energy functionals  
 $z \in \underbrace{\{e_1, \dots, e_k, \dots, e_N\}}_{\text{mart1}} \subset \mathbb{R}^N$   $W(\mathbf{E}, e_j), j = 1, \dots, N$   
 $\underbrace{\hspace{10em}}_{\text{martk}} \underbrace{\hspace{10em}}_{\text{aust}}$

Mixtures  $z \in Z = \text{conv}\{e_1, \dots, e_N\}$  is the Gibbs simplex

$W(\mathbf{E}, z) : \mathbb{R}_{\text{sym}}^{d \times d} \times Z \rightarrow \mathbb{R}$  is the **mixture function** called the **free-energy of mixing** by Govindjee, Hackl & Heinen'07

**State variables**

$u : \Omega \rightarrow \mathbb{R}^d$  displacement

$z : \Omega \rightarrow Z$  phase mixture

$u_{\text{Dir}}$ : the time-dependent Dirichlet boundary data

**Applied fields**

$\ell_{\text{appl}} : [0, T] \rightarrow \mathcal{F}^*$  mechanical loading

$\theta_{\text{appl}} : [0, T] \times \Omega \rightarrow \mathbb{R}$  temperature given



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$u : \Omega \rightarrow \mathbb{R}^d$  displacement

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$u_{\text{Dir}} \in C^1([0, T]; H^1(\Omega; \mathbb{R}^d))$

**Applied fields**

$\ell_{\text{appl}} \in C^1([0, T]; H^1(\Omega; \mathbb{R}^d)')$

$\theta_{\text{appl}} \in C^1([0, T]; L^\infty(\Omega; [\theta_{\text{min}}, \theta_{\text{max}}]))$



**Energy:**  $\mathcal{E}(t, u, z) = \int_{\Omega} (W(x, e(u + u_{\text{Dir}}(t)), z, \theta(t)) + \frac{\sigma}{2} |\nabla z|^2) dx - \langle \ell(t), u \rangle$

where  $e(u) = \frac{1}{2}(\nabla u + \nabla u^T)$  is the infinitesimal strain

**Dissipation distance:**  $\mathcal{D}(z_1, z_2) = \int_{\Omega} \psi(x, z_2 - z_1) dx$

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Mielke & Theil'04.

$(u, z) : [0, T] \rightarrow \mathcal{F} \times \mathcal{Z}$  is called **energetic solution**, if

**(S)**  $\mathcal{E}(t, u(t), z(t)) \leq \mathcal{E}(t, \tilde{u}, \tilde{z}) + \mathcal{D}(z(t), \tilde{z})$  for all  $(\tilde{u}, \tilde{z}) \in \mathcal{F} \times \mathcal{Z}$

**(E)**  $\mathcal{E}(t, u(t), z(t)) + \text{Var}_{\mathcal{D}}(z; [0, t]) = \mathcal{E}(0, u_0, z_0) + \int_0^t \partial_s \mathcal{E}(\cdot, u, z) ds$



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**Previous works:**

- ▷ Govindjee, Mielke & Hall'03
- ▷ Souza, Mamiya & Zouain'98, Auricchio & Petrin'04, Mielke & P.'07



**Assumptions on  $W$ .** Let  $\omega : [0, \infty) \rightarrow [0, \infty)$  be a nondecreasing function with  $\lim_{\tau \rightarrow 0^+} \omega(\tau) = 0$  such that

$W(\cdot, z, \theta)$  is strictly convex,

$$W, \partial_\theta W \in C^0(\mathbb{R}^{d \times d} \times Z \times [\theta_{\min}, \theta_{\max}]; \mathbb{R})$$

$$\partial_e W \in C^0(\mathbb{R}^{d \times d} \times Z \times [\theta_{\min}, \theta_{\max}]; \mathbb{R}^{d \times d})$$

$$c(|e|^2 + |z|^2) - C \leq W(e, z, \theta) \leq c(|e|^2 + |z|^2) + C$$

$$|\partial_e W(e, z, \theta)|^2 + |\partial_\theta W(e, z, \theta)| \leq C_1^W (W(e, z, \theta) + C_0^W)$$

$$|\partial_\theta W(e, z, \theta_1) - \partial_\theta W(e, z, \theta_2)| \leq C_1^\theta (W(e, z, \theta_1) + C_0^\theta) \omega(|\theta_1 - \theta_2|)$$

$$|\partial_e W(e, z, \theta_1) - \partial_e W(e, z, \theta_2)|^2 \leq C_1^e (W(e, z, \theta_1) + C_0^e) \omega(|\theta_1 - \theta_2|)$$

$$|\partial_\theta W(e_1, z_1, \theta) - \partial_\theta W(e_2, z_2, \theta)| \leq C^\theta (|e_1 - e_2| + |z_1 - z_2|)(1 + |e_1 + e_2| + |z_1 + z_2|)$$

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**Assumption on  $\mathcal{D}$ .**  $c_1|z_1 - z_2| \leq D(z_1, z_2) \leq c_2|z_1 - z_2|$



## Theorem (The Existence result)

*Under the assumptions on  $W$  and  $\psi$  given above. Let  $(u(0), z(0))$  satisfies **(S)**. Then there exists  $(u, z) : [0, T] \rightarrow \mathcal{F} \times \mathcal{Z}$  satisfying **(S)** & **(E)** such that*

$$u \in L^\infty([0, T]; H^1(\Omega; \mathbb{R}^d))$$

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If  $W$  is  $\alpha$ -uniformly convex jointly in the first two arguments then  $(u, z)$  is Lipschitz continuous



Finite-element spaces  $\mathcal{F}_h \subset \mathcal{F}$  and  $\mathcal{Z}_h \subset \mathcal{Z}$ , time step  $\tau > 0$

Density assumption:  $\underbrace{(u_h, z_h)}_{\in \mathcal{F}_h \times \mathcal{Z}_h} \rightarrow (u, z)$  strongly in  $\mathcal{F} \times \mathcal{Z}$

## Space-Time Discretization

$$\text{(IP)}^{h,\tau} \quad q_k^{h,\tau} \in \underset{\hat{q}^h \in \mathcal{F}_h \times \mathcal{Z}_h}{\text{Argmin}} (\mathcal{E}(t_k^\tau, \hat{q}^h) + \mathcal{D}(z_{k-1}^{h,\tau}, \hat{z}^h))$$

where  $q_k^{h,\tau} = (u_k^{h,\tau}, z_k^{h,\tau})$  and  $\hat{q}^h = (\hat{u}^h, \hat{z}^h)$



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Like Backward Euler for the heat equation:

$$u_t + \Delta u = 0 \quad u_k^{h,\tau} \in \underset{\hat{u}}{\text{Argmin}} \left( \int_{\Omega} \frac{1}{2} |\nabla \hat{u}|^2 dx + \frac{1}{2(t_k - t_{k-1})} \|\hat{u} - u_{k-1}^{h,\tau}\|^2 \right)$$

Piecewise constant interpolants  $\bar{q}^{h,\tau}(t) = (\bar{u}^{h,\tau}, \bar{z}^{h,\tau}) : [0, T] \rightarrow \mathcal{F}_h \times \mathcal{Z}_h$



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**AIM:** investigate the asymptotics as  $h \rightarrow 0$  and  $\tau \rightarrow 0$



Convergence of the space-time discretization follows from

- ▷ uniform a priori estimates  $\rightsquigarrow$  *numerical stability*
- ▷ accumulation points are solutions  $\rightsquigarrow$  *consistency*

**Previous works:** Mielke'05, Roubíček & Mielke'06

## Theorem (Convergence of the space-time discretization)

*Under the assumptions given above. Then there exists a subsequence  $\bar{q}^{h,\tau} = (\bar{u}^{h,\tau}, \bar{z}^{h,\tau})$  which converges to a solution  $(u, z)$  of **(S)&(E)** and*

$$u \in L^\infty([0, T]; H^1(\Omega; \mathbb{R}^d))$$

$$z \in L^\infty([0, T]; H^1(\Omega; Z))$$

$$z \in \text{BV}([0, T]; L^1(\Omega; Z))$$

*such that*

$$\bar{q}^{h_n, \tau_n}(t) \rightarrow q(t) \text{ strongly in } \mathcal{F} \times \mathcal{Z}$$

$$\mathcal{E}(t, \bar{q}^{h_n, \tau_n}(t)) \rightarrow \mathcal{E}(t, q(t))$$

$$\text{Var}_{\mathcal{D}}(\bar{q}^{h_n, \tau_n}; [0, t]) \rightarrow \text{Var}_{\mathcal{D}}(q; [0, t])$$





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- ▷ understand the limit when  $\sigma \rightarrow 0$  (formation of microstructure)
- ▷ include **rate-dependent** effects like a heat equation
- ▷ develop the theory to include other **multifunctional materials** (ferroelectric materials, magnetostrictive materials)



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**Thank you for your attention !**

Papers on line: <http://www.wias-berlin.de/people/petrov>