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Coupling of Fluid Flow and Solute Transport Using a Divergence-Free Reconstruction of the Crouzeix-Raviart Element



ENERGIESPEICHER Forschungsinitiative der Bundesregierung

Problem Description: Transport in a Fluid Flow

The coupling between a fluid flow with velocity \vec{u} , pressure p and viscosity η and the steady transport of a dissolved species with concentration c and diffusion coefficient D is governed by the incompressible Navier-Stokes equations

 $-\eta \Delta \vec{u} + (\vec{u} \cdot \nabla)\vec{u} + \nabla p = \vec{f} \text{ in } \Omega, \quad \nabla \cdot \vec{u} = 0 \text{ in } \Omega \text{ and } \vec{u} = \vec{u}_D \text{ along } \partial \Omega$

Coupling To Solute Transport by FVM

The transport equation is discretised with a finite volume method. Let \mathcal{P} denote a set of points of a Delaunay mesh and let \mathcal{K} denote the associated set of Voronoi cells with facets \mathcal{F} .

and the transport equation

 $\nabla \cdot (-D\nabla c + c\vec{u}) = s \text{ in } \Omega$ and $c = c_D$ along $\partial \Omega$.

The discretisation of the divergence-constraint is a crucial part of the coupling for numerical stability and physical correctness (e.g. maximum principles).

Flow Solver: Modified Crouzeix-Raviart FEM

The Navier-Stokes equations are discretised with a modified Crouzeix-Raviart finite element method from [2] on a regular triangulation \mathcal{T} . The velocity test functions $CR(\mathcal{T})$ are piecewise affine, vector-valued functions and continuous in the barycenters of the faces, while the pressure test functions $Q(\mathcal{T})$ are piecewise constant with zero integral mean.

The modified method seeks $(\vec{u}_h, p_h) \in CR(\mathcal{T}) \times Q(\mathcal{T})$ with $\int_E \vec{u}_D - \vec{u}_h ds = 0$ for every boundary face *E*, and

$$\int_{\Omega} \eta \nabla \vec{u}_h : \nabla \vec{v}_h \, dx + \int_{\Omega} (\Pi_{\mathrm{RT}} \vec{u}_h \cdot \nabla) \vec{u}_h \cdot \Pi_{\mathrm{RT}} \vec{v}_h \, dx - \int_{\Omega} p_h \nabla \cdot \vec{v}_h \, dx = \int_{\Omega} \vec{f} \cdot \Pi_{\mathrm{RT}} \vec{v}_h \, dx,$$

On every
$$\sigma_{KL} := \partial K \cap \partial L \in \mathcal{F}$$
 define
 $u_{\sigma_{KL}} := \int_{\sigma_{KL}} \prod_{RT} \vec{u}_h \cdot (\vec{x}_L - \vec{x}_K) \, ds / \left| \sigma_{KL} \right|,$
 $\tau_{\sigma_{KL}} := \left| \sigma_{KL} \right| / \left| \vec{x}_L - \vec{x}_K \right|.$



Since $\nabla \cdot \prod_{RT} \vec{u}_h = 0$, the fluxes $u_{\sigma_{KL}}$ are discretely divergence-free in the FV sense

$$\sum_{K \text{ neighbour of } K} \tau_{\sigma_{KL}} u_{\sigma_{KL}} = \int_{\partial K} \Pi_{\mathrm{RT}} \vec{u}_h \cdot \vec{n}_K \, dx = \int_K \nabla \cdot \Pi_{\mathrm{RT}} \vec{u}_h \, dx = 0 \quad \text{ for all } K \in \mathcal{K}.$$

This guarantees the preservation of maximum principles.

The finite volume scheme seeks $c_h \in P_0(\mathcal{K})$ with $c_K := c_h|_K = c_D(\vec{x}_K)$ for all cells $K \in \mathcal{K}_D$ at the Dirichlet boundary and

$$\sum_{K \in K} \tau_{\sigma_{KL}} g(c_K, c_L, u_{\sigma_{KL}}) = |K| \, s_K \quad \text{for all } K \in \mathcal{K}_0 := \mathcal{K} \setminus \mathcal{K}_D$$
 where we have a neighbour of K

where $g(c_K, c_L, u_{\sigma_{KL}}) := D(B(u_{\sigma_{KL}}/D)c_K - B(-u_{\sigma_{KL}}/D)c_L)$ define the exponentially fitted flux approximations with Bernoulli function $B(z) = z/(1 - e^{-z})$.

 $-\int_{\Omega} q_h \nabla \cdot \vec{u}_h \, dx = 0 \quad \text{for all } (\vec{v}_h, q_h) \in \operatorname{CR}_0(\mathcal{T}) \times Q(\mathcal{T}).$

The Fortin interpolation π_{RT} maps discretely divergence-free test functions to globally divergence-free Raviart-Thomas test functions.

Advantages

- Divergence-free reconstruction ensures preservation of maximum principles for the solute transport (see below).
- Significantly smaller number of degrees of freedom compared to previously introduced divergence-free coupling schemes (e.g. Scott-Vogelius FEM [1]).
- Optimal pressure-independent a priori velocity error estimate. BDM interpolation instead of Π_{RT} yields optimal L^2 error convergence rates [3].

Preservation of Maximum Principle: CR-FEM vs. Modified CR-FEM



Application: Flowcell Simulation



Divergence-free coupling methods allow e.g. to simulate the transport of ions in an electrolyte through a flowcell to validate measurements and determine problem parameters in regimes not covered by the literature. The figure on the left depicts isosurfaces of the ion concentration from [1]. The ultimate goal is to find an efficient discretisation scheme with reasonable complexity.

References

- J. Fuhrmann, H. Langmach, and A. Linke. A numerical method for mass conservative coupling between fluid flow and solute transport. *Appl. Numer. Math.*, 61(4):530–553, 2011.
- A. Linke. On the role of the Helmholtz decomposition in mixed methods for incompressible flows and a new variational crime. *Comput. Methods Appl. Mech. Engrg.*, 268:782–800, 2014.

Isolines (on the base) and elevation graph of stationary concentration in the longitudinal section of an U-shaped pipe: standard CR-FEM velocity without re-construction (left) and with modified CR-FEM velocity (right).

C. Brennecke, A. Linke, C. Merdon, and J. Schöberl. Optimal and pressure-independent L2 velocity error estimates for a modified Crouzeix-Raviart Stokes element with BDM reconstructions, 2014. WIAS Preprint 1929.

This research has been partially funded in the framework of the project "Macroscopic Modeling of Transport and Reaction Processes in Magnesium-Air-Batteries" (Grant 03EK3027D) under the research initiative "Energy storage" of the German Federal government.

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