

g Infinite-Volume Gibbs states

- 1) Pressure provides important macroscopic / thermodynamic properties of the Ising model in infinite volume, e.g. magnetization; or high derivatives of Ψ
- 2) We like to also evaluate the infinite-volume Ising model w.r.t. local observable; e.g. expected spin at origin.

Def (local observable)

A function $f: \Omega \rightarrow \mathbb{R}$ is local if $\exists \Lambda \subset \mathbb{Z}^d$, s.t. that $f(\omega) = f(\omega^\Lambda) \forall \omega, \omega' \in \Omega$ s.t. $\omega_i = \omega'_i \forall i \in \Lambda$.
The smallest s.t. set Λ is called the support of f , $\text{supp}(f)$.

Ex: $\omega \mapsto m_n(\omega)$ is local for $\Lambda \subset \mathbb{Z}^d$
 $\omega \mapsto \omega_i$ is local for $\{i\} \subset \mathbb{Z}^d$

Def (infinite-volume state)

A mapping $f \mapsto \langle f \rangle$ is a state if

1) $\langle 1 \rangle = 1$ normalized

2) $\langle f \rangle \geq 0 \forall f \geq 0$ positive

3) $\langle f + \lambda g \rangle = \langle f \rangle + \lambda \langle g \rangle \forall g, f, \lambda \geq 0$ linear

$\langle f \rangle$ is called the average of f under the state $\langle \cdot \rangle$

Def (Gibbs state)

Let $\Lambda \uparrow \mathbb{Z}^d$ and $(\gamma_n)_{n \geq 1}$ a sequence of boundary conditions. The Gibbs measure $(\mu_{\Lambda, \gamma_n}^z)_{n \geq 1}$ converges to the $\langle \cdot \rangle$: $\Leftarrow \Rightarrow$

$$\lim_{n \uparrow \infty} \mu_{\Lambda, \gamma_n}^z(f) = \langle f \rangle \quad \forall \text{ local } f.$$

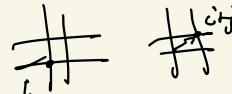
$\langle \cdot \rangle$ is called a Gibbs state at (Λ, h) .

Remark: 1) The topology of local convergence label measures close if they are close label against local observables

2) In fact $\forall \langle \cdot \rangle$, $\exists \mu$ s.t. $\langle f \rangle = \mu(f)$
(Riesz theorem)

3) So convergence tested on local functions can be seen as weak-convergence of probability measures, where the continuity of the test-function is via the product topology

Important class of infinite-volume states are the translation-invariant states:

Let $\Theta_j: \mathbb{Z}^d \mapsto \mathbb{Z}^d \quad i \mapsto \Theta_j \cdot i = i + j$ 

$\Theta_j: \Omega \mapsto \Omega \quad \omega \mapsto \Theta_j \omega$ with $(\Theta_j \omega)_i = \omega_{i-j}$

Def (Translation-invariant states)

A state $\langle \cdot \rangle$ is called translation invariant if

$$\langle \varphi \circ \theta_j \rangle = \langle \varphi \rangle \quad \forall \text{ local } \varphi, j \in \mathbb{Z}^d$$

Do Gibbs states exist for the Ising model?

Thm (Extremal Ising states)

Let $\beta \geq 0, h \in \mathbb{R}$, then there exist translation-invariant states $\langle \cdot \rangle_{\beta, h}^+$ and $\langle \cdot \rangle_{\beta, h}^-$ such that

$$\lim_{n \rightarrow \infty} \mu_{\Lambda_n; \beta, h}^{2^{\pm}} = \langle \cdot \rangle_{\beta, h}^{\pm}$$

independently from the converging sequence $\Lambda_n \uparrow \mathbb{Z}^d$.

⑩ Existence of extremal Ising models

Strategy: 1) Identify limits as monotone limits.
 2) Use well-chosen class of local functions.

$$\forall A \subset \mathbb{Z}^d$$

$$\sigma_A := \prod_{j \in A} \prod_{\substack{\sigma_j \\ \in \pm 1}} \sigma_j \quad ; \quad \tau_A := \prod_{j \in A} \tau_j$$

$\frac{1}{2}(1 + \sigma_j) \in \{0, 1\}$
occupation variables

Lemma (σ_A and τ_A basis for local functions)

$\forall f$ local $\exists (\hat{\varphi}_A)_{A \subset \text{supp}(f)}, (\tilde{\varphi}_A)_{A \subset \text{supp}(f)} \in \mathbb{R}$
 s.t. that

$$f = \sum_{A \subset \text{supp}(f)} \hat{\varphi}_A \sigma_A = \sum_{A \subset \text{supp}(f)} \tilde{\varphi}_A \tau_A$$

Proof: 1) Orthogonality: $\forall B \subset \mathbb{Z}^d, \omega, \tilde{\omega} \in \Omega$

$$2^{-|B|} \sum_{A \subset B} \sigma_A(\tilde{\omega}) \sigma_A(\omega) = \mathbb{1} \{ \omega_i = \tilde{\omega}_i, \forall i \in B \}$$

Indeed: 1) if $\omega_i = \tilde{\omega}_i \forall i \in B$, then $\forall A \subset B \quad \sigma_A(\tilde{\omega}) \sigma_A(\omega) = \prod_{i \in A} \underbrace{\omega_i \tilde{\omega}_i}_{=1} = 1$

2) $\exists i \in B \quad \omega_i \neq \tilde{\omega}_i$, then

$$\sum_{A \subset B} \sigma_A(\tilde{\omega}) \sigma_A(\omega) = \sum_{A \subset B, i \in A} (\sigma_A(\tilde{\omega}) \sigma_A(\tilde{\omega}) + \sigma_{A \cup \{i\}}(\tilde{\omega}) \sigma_{A \cup \{i\}}(\omega))$$

$$-\overbrace{\tau_A(\omega)}^{\tau_A(\omega)} \tau_A(\omega)$$

$$= 0$$

2) Let $B = \text{supp}(f)$

$$\begin{aligned} \psi(\omega) &= \sum_{\omega' \in \Omega_B} \psi(\omega') \underbrace{\mu\{\omega'_i = \omega_i \forall i \in B\}}_{e^{-|B|} \sum_{A \subset B} \tau_A(\omega) \tau_A(\omega')} \\ &= \sum_{A \subset B} \tau_A(\omega) \left(e^{-|B|} \sum_{\omega' \in \Omega_B} \psi(\omega') \tau_A(\omega') \right) \\ &\quad \underbrace{\uparrow}_{\psi_A} \end{aligned}$$

Sketch for τ_A

□

Using the Lemma, we can check convergence w.r.t. τ_A and τ_A .

Key tool to derive monotone limits is via correlation inequalities

- 1) GKS inequality (Griffiths, Kelly, Sherman)
- 2) FKG inequality (Fortuin, Kasteleyn, Ginibre)

1) GKS inequalities: Idea: $\mu_{\beta, h}^{\tau^+}(\tau_i) \geq 0$
 $h \geq 0 \quad i \in \Lambda$

+boundary & non-negative external field should create bias for the expected spin at any site $i \in \Lambda$, $\forall \beta \in \mathbb{R}^d$.

Even $\mu_{1; \beta, h}^{2^+}(\sigma_i=1; \sigma_j=1) \geq \mu_{1; \beta, h}^{2^+}(\sigma_i=1) \mu_{1; \beta, h}^{2^+}(\sigma_j=1)$

\uparrow \uparrow \uparrow
 support each other positive correlation!

Holds more generally:

$$\bar{J} = (J_{ij})_{i,j \in E} \quad ; \quad \bar{h} = (h_i)_{i \in \mathbb{Z}^d}$$

$$\text{let } H_{1; \bar{J}, \bar{h}}(\omega) = - \sum_{i,j \in E} J_{ij} \omega_i \omega_j - \sum_{i \in \mathbb{Z}^d} h_i \omega_i$$

and let $\mu_{1; \bar{J}, \bar{h}}^{2^+}$ the associated Ising model.

Thm (GKS)

Consider \bar{J}, \bar{h} and $\Lambda \subset \mathbb{Z}^d$ with $h_i \geq 0 \quad \forall i$.

Then $\forall A, B \subset \Lambda$

$$\mu_{\Lambda; \bar{J}, \bar{h}}^{2^+}(\sigma_A) \geq 0$$

$$\mu_{\Lambda; \bar{J}, \bar{h}}^{2^+}(\sigma_A \sigma_B) \geq \mu_{\Lambda; \bar{J}, \bar{h}}^{2^+}(\sigma_A) \mu_{\Lambda; \bar{J}, \bar{h}}^{2^+}(\sigma_B)$$

These inequalities also hold for ϕ , by replacing 2^+ .

No proof.

2) FKG inequality:

$$\omega \leq \omega' \quad : \Leftrightarrow \quad \omega_i \leq \omega'_i \quad \forall i \in \mathbb{Z}^d$$

$$E \subset \Omega \text{ increasing} : \Leftrightarrow \text{if } \omega \in E \Rightarrow \omega' \in E \quad \forall \omega \leq \omega'$$

$$\Psi \text{ non-decreasing} : \Leftrightarrow \Psi(\omega) \leq \Psi(\omega') \quad \forall \omega \leq \omega'$$

Idea: ① $\mu_{\lambda, \rho, h}^{2+} (E \cap F) \geq \mu_{\lambda, \rho, h}^{2+} (E) \mu_{\lambda, \rho, h}^{2+} (F)$

② $\mu(fg) - \mu(f)\mu(g) = \frac{1}{2} \int (f(x) - f(y))(g(x) - g(y)) \mu(dx) \mu(dy)$

\uparrow
 both increasing or
 both decreasing

≥ 0 ≥ 0
 ≤ 0 ≤ 0

≥ 0

Note that:

$w \mapsto w_i = v_i(w)$	}	are non-decreasing
$w \mapsto z_i(w)$		
$w \mapsto z_A(w) \quad \forall A \subset \mathbb{Z}^d$		

Thm (FKG)

(consider \mathcal{J} with $\mathcal{J}_{ij} \geq 0$, \bar{h} , $\Lambda \subset \mathbb{Z}^d$; $\mathcal{Z} \in \Omega$.)

Then for two non-decreasing functions f and g

$$\mu_{\lambda, \bar{\mathcal{J}}, \bar{h}}^2 (fg) \geq \mu_{\lambda, \bar{\mathcal{J}}, \bar{h}}^2 (f) \mu_{\lambda, \bar{\mathcal{J}}, \bar{h}}^2 (g)$$

No proof:

Lemma (Monotonicity for non-decreasing functions)

Let f be non-decreasing and $\Lambda_1 \subset \Lambda_2 \subset \mathbb{Z}^d$,

then for $\lambda \geq 0, h \in \mathbb{R}$

$$\mu_{\lambda_1, \rho, h}^{2+} (f) \geq \mu_{\lambda_2, \rho, h}^{2+} (f)$$

this also holds for $\mathcal{Z}^+ \mapsto \mathcal{Z}^-$ and non-decreasing \mapsto non-increasing.

Proof: Spatial Markov property

$$\begin{aligned} \mu_{\Lambda_1; \Delta, h}^{\mathbb{Z}^+}(f) &= \mu_{\Lambda_2; \Delta, h}^{\mathbb{Z}^+}(f \mid \underbrace{w_i = 1 \ \forall i \in \Lambda_2 \setminus \Lambda_1}_{\text{is non-decreasing}}) \\ &\geq \mu_{\Lambda_2; \Delta, h}^{\mathbb{Z}^+}(f) \end{aligned} \quad \square$$

Lemma (Extremality)

let f be non-decreasing; $\beta \geq 0$ and $h \in \mathbb{R}$

$$\mu_{\Lambda; \Delta, h}^{\mathbb{Z}}(f) \leq \mu_{\Lambda; \Delta, h}^{\mathbb{Z}}(f) \leq \mu_{\Lambda; \Delta, h}^{\mathbb{Z}^+}(f)$$

$\forall \Lambda \subset \subset \mathbb{Z}^d$ and $\gamma \in \Omega$ and even $\gamma \mapsto \phi$, so if $\text{supp}(f) \subset \Lambda$

Proof: let $\mathbb{I}(w) = e^{\beta \sum_{i \in \Lambda, j \notin \Lambda} w_i (1 - w_j)}$ for $w \in \Omega_{\Lambda}^{\mathbb{Z}^+}$
non-decreasing.

$$\sum_{w \in \Omega_{\Lambda}^{\mathbb{Z}^+}} e^{-\beta \mathcal{H}_{\Lambda; \Delta, h}(w)} = \sum_{w \in \Omega_{\Lambda}^{\mathbb{Z}^+}} e^{-\beta \mathcal{H}_{\Lambda; \Delta, h}(w)} \mathbb{I}(w)$$

$$\text{further: } \sum_{w \in \Omega_{\Lambda}^{\mathbb{Z}^+}} e^{-\beta \mathcal{H}_{\Lambda; \Delta, h}(w)} f(w) \geq \sum_{w \in \Omega_{\Lambda}^{\mathbb{Z}^+}} e^{-\beta \mathcal{H}_{\Lambda; \Delta, h}(w)} f(w) \mathbb{I}(w)$$

\uparrow
 f might depend on the boundary.

$$\begin{aligned} \Rightarrow \mu_{\Lambda; \Delta, h}^{\mathbb{Z}^+}(f) &= \frac{\sum_{w \in \Omega_{\Lambda}^{\mathbb{Z}^+}} e^{-\beta \mathcal{H}_{\Lambda; \Delta, h}(w)} f(w)}{\sum_{w \in \Omega_{\Lambda}^{\mathbb{Z}^+}} e^{-\beta \mathcal{H}_{\Lambda; \Delta, h}(w)}} \\ &\geq \frac{\sum_{w \in \Omega_{\Lambda}^{\mathbb{Z}^+}} e^{-\beta \mathcal{H}_{\Lambda; \Delta, h}(w)} f(w) \mathbb{I}(w)}{\sum_{w \in \Omega_{\Lambda}^{\mathbb{Z}^+}} e^{-\beta \mathcal{H}_{\Lambda; \Delta, h}(w)} \mathbb{I}(w)} \end{aligned}$$

$$= \frac{\mu_{\Lambda, h}^2(\Gamma \dagger)}{\mu_{\Lambda, h}^2(\Gamma)} \stackrel{\text{alg}}{\geq} \mu_{\Lambda, h}^2(f)$$

Similar for $\bar{\Gamma}$; ϕ ; $\rho_0 \rightarrow$ as exercise. □

Proof for existence of extremal Gibbs states

Let f be local, then

$$\begin{aligned} \mu_{\Lambda, h}^{2+}(f) &= \sum_{\text{Acc supp}(f)} \tilde{\varphi}_f \underbrace{\mu_{\Lambda, h}^{2+}(\tau_A)}_{\substack{\uparrow \text{non-decreasing } \forall A \\ \geq \mu_{\Lambda, h}^{2+}(\tau_A) \geq 0 \\ \forall n \geq 1}} \end{aligned}$$

\Rightarrow convergence:

$$\lim_{n \rightarrow \infty} \mu_{\Lambda, h}^{2+}(f) =: \langle f \rangle_{\Lambda, h}^+ \text{ a Gibbs state.}$$

Independence of sequence via alternating and increasing sequence

$$\Delta_{2k-1} \subset \{ \Lambda_n^{(1)} : n \geq 1 \}; \Delta_{2k} \subset \{ \Lambda_n^{(2)} : n \geq 1 \}; \Delta_k \subset \Delta_{k+1}$$

For the translation invariance first note that $f \circ \theta_j$ is local

$$\begin{aligned} \mu_{\Lambda, h}^{2+}(f) &= \mu_{\Lambda, h}^{2+}(f \circ \theta_j) \\ \downarrow & \qquad \qquad \downarrow \\ \langle f \rangle_{\Lambda, h}^+ &= \langle f \circ \theta_j \rangle_{\Lambda, h}^+ \end{aligned}$$

□