

# Coulomb and spin-orbit interactions in nanorings

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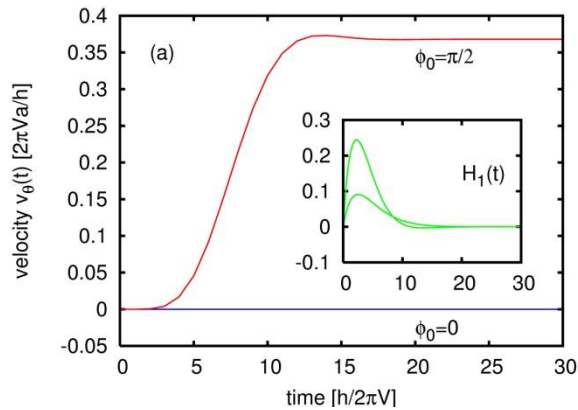
WIAS Workshop 2011  
5 February



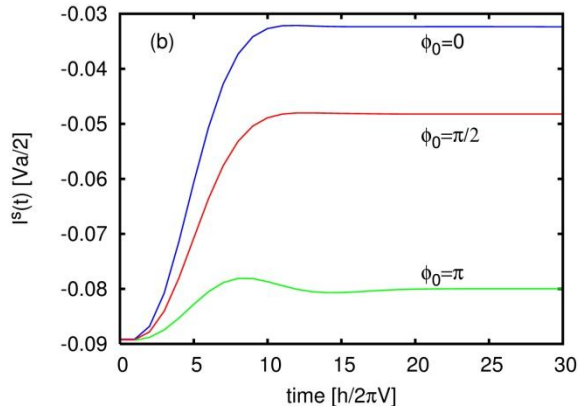
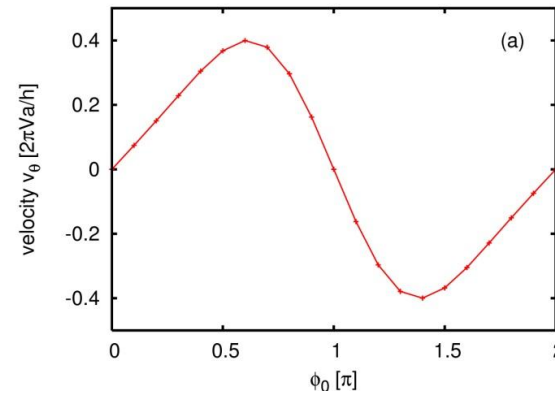
# Persistent charge and spin currents

1D ring with SOI interaction excited by an external EM pulse

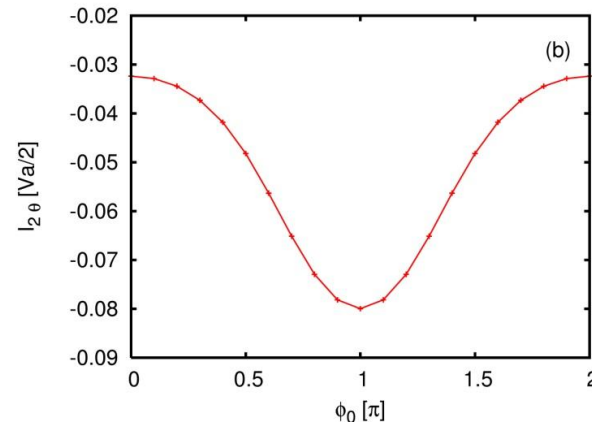
$$V_{pulse}(t) = Ae^{-\Gamma t} \left[ \sin(\omega_1 t) \cos \varphi + \sin(\omega_2 t) \cos(\varphi + \varphi_0) \right]$$



Charge current



Spin current



# SOI models

Rashba - Due to the vertical confinement

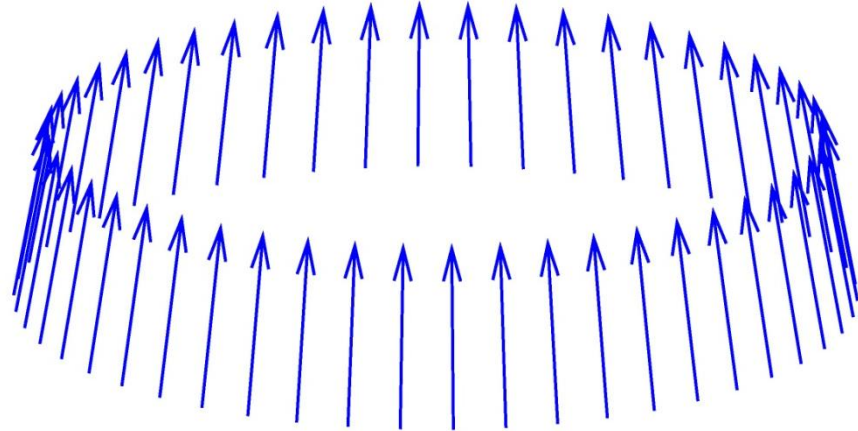
$$H_R = \frac{\alpha_R}{\hbar} (\sigma_x p_y - \sigma_y p_x)$$

Dresselhaus – Due to the electric field in the crystal

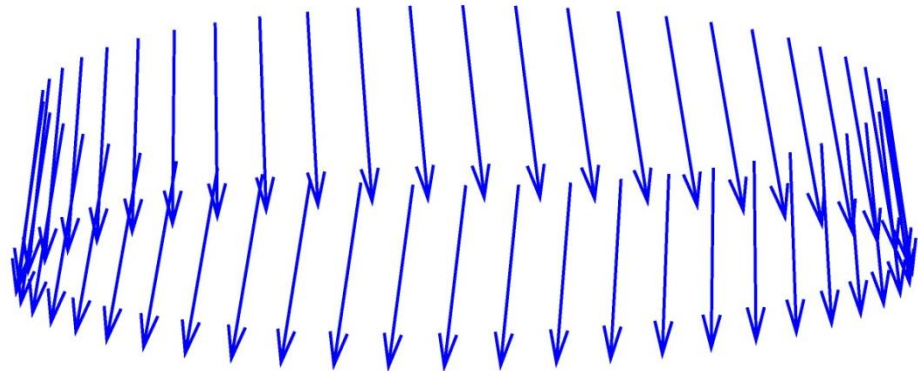
$$H_D = \frac{\beta_D}{\hbar} (\sigma_x p_x - \sigma_y p_y)$$

# Known results on the 1D ring

Rashba



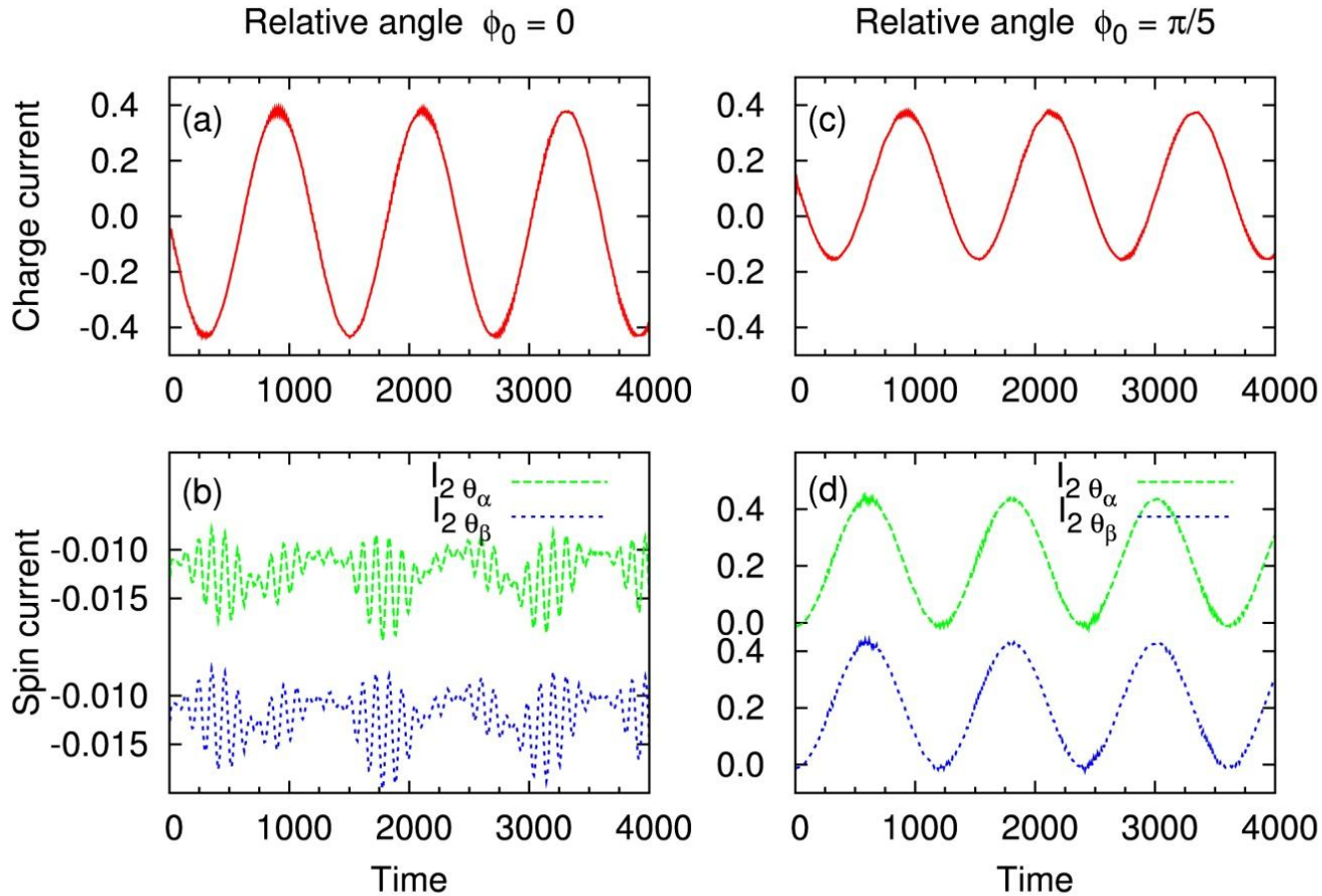
Dresselhaus



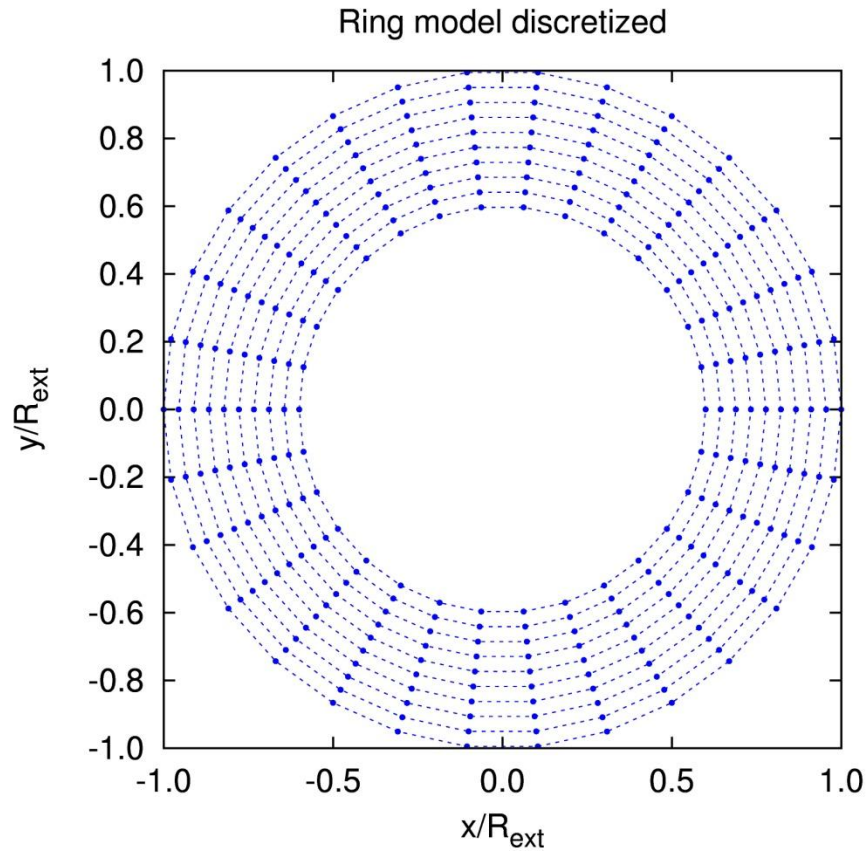
3 electrons, no  
Coulomb interaction

J. S. Cheng and K. Chang PRB 74, 235315 (2006)

# Charge and spin currents for R+D SOI



# The sample



Between 1-10 rings and 20-100 angular sites all connected by hopping matrix elements

# The sample Hamiltonian

$$H = -\frac{\hbar^2}{2m} \Delta = -\frac{\hbar^2}{2m} \left[ \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2}{\partial \varphi^2} \right]$$

Matrix elements for discrete polar coordinates:

$$\begin{aligned} \langle r\varphi | H | r'\varphi' \rangle &= t_r \left( 2\delta_{rr'} - \delta_{r,r'-1} - \delta_{r,r'+1} \right) \delta_{\varphi\varphi'} \\ &\quad + t_\varphi \left( 2\delta_{\varphi\varphi'} - \delta_{\varphi,\varphi'-1} - \delta_{\varphi,\varphi'+1} \right) \delta_{rr'} \end{aligned}$$

$$t_R = \frac{\hbar^2}{2mR^2} \quad t_r = t_R \left( \frac{R}{\delta r} \right)^2 \quad t_\varphi = t_R \left( \frac{R}{r\delta\varphi} \right)^2$$

# The Coulomb interaction

One particle:  $H_1 \psi_a = \varepsilon_a \psi_a$

N particles:  $H_N = \sum_a \varepsilon_a c_a^\dagger c_a + \frac{1}{2} \sum_{a,b;c,d} V_{ab,cd} c_a^\dagger c_b^\dagger c_d c_c$

Coulomb matrix elements:  $V_{ab,cd} = \int d\vec{r} d\vec{r}' \psi_a^*(\vec{r}) \psi_b^*(\vec{r}') \frac{e^2}{|\vec{r} - \vec{r}'|} \psi_c(\vec{r}) \psi_d(\vec{r}')$

Many-body states for independent electrons: ( $V_{ab,cd} = 0$ )

$|\alpha\rangle$ :  $|1, 0, 0, \dots\rangle$   $|0, 1, 0, \dots\rangle$   $|1, 1, 0, \dots\rangle$   $|1, 1, 1, \dots\rangle$

"bit strings" with N particles

$V_{ab,cd} \neq 0 \Rightarrow H_N |\mu\rangle = E_\mu |\mu\rangle$

States for interacting electrons:  $|\mu\rangle = \sum_\alpha A_{\alpha\mu} |\alpha\rangle$  and



# The “exact diagonalization”

$$H_N |\mu\rangle = E_\mu |\mu\rangle$$

We want eigenvalues/vectors of the matrix  $\langle\alpha|H_N|\beta\rangle$

All we need are the matrix elements  $V_{ab,cd}$  and  $\langle\alpha|c_a^\dagger c_b^\dagger c_c c_d|\beta\rangle$

Creation/destruction (annihilation) operators:

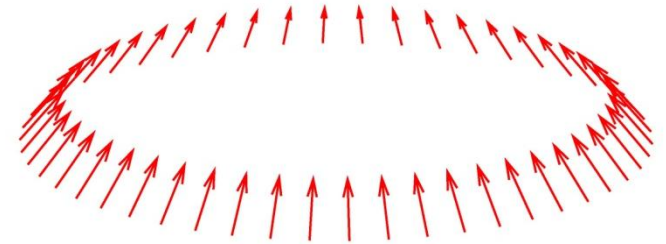
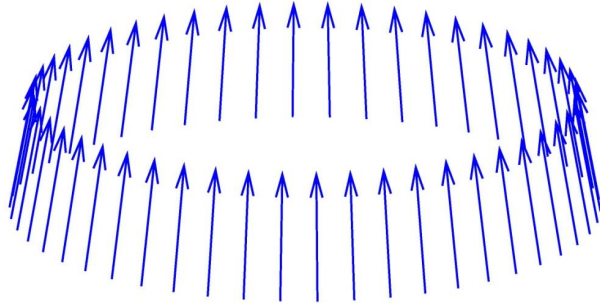
$$c_1^\dagger |\underline{0}, 1, 0, \dots\rangle = |\underline{1}, 1, 0, \dots\rangle \quad c_1 |\underline{1}, 1, 0, \dots\rangle = |\underline{0}, 1, 0, \dots\rangle$$

$$c_1^\dagger |\underline{1}, 1, 0, \dots\rangle = 0 \quad c_1 |\underline{0}, 1, 0, \dots\rangle = 0$$

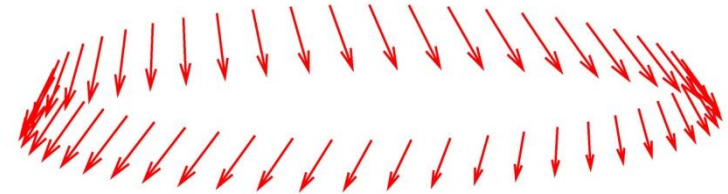
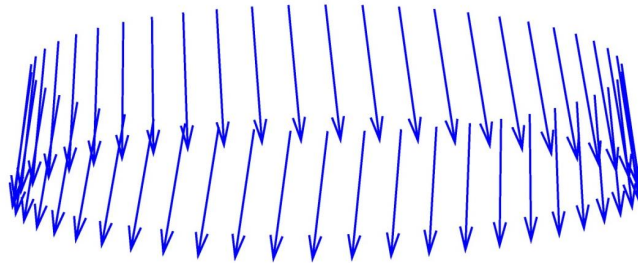
$$\langle\alpha|c_a^\dagger c_b^\dagger c_c c_d|\beta\rangle = \langle c_a c_b \alpha | c_c c_d \beta \rangle = 0, \pm 1$$

# Combined Coulomb and SOI effects (groundstate)

Rashba



Dresselhaus



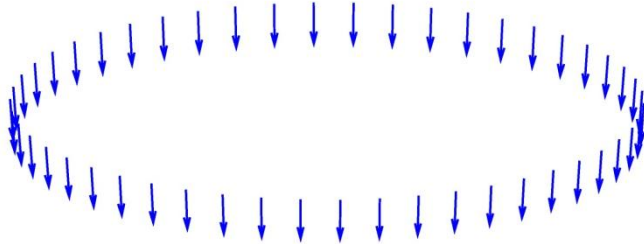
N=3 no interaction

N=3 with interaction

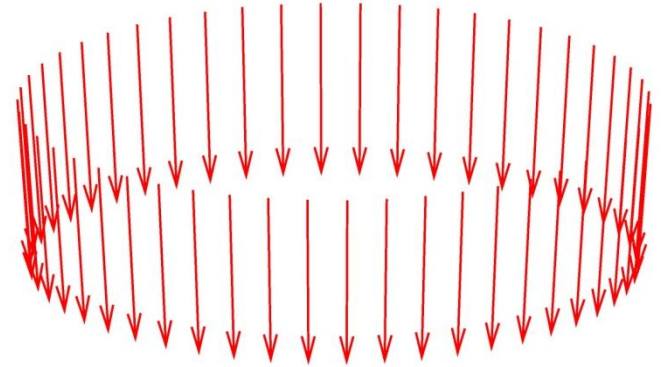
**The spin polarization may be reduced by the Coulomb interaction ...**

# Combined Coulomb and SOI effects

Rashba



N=4 no interaction

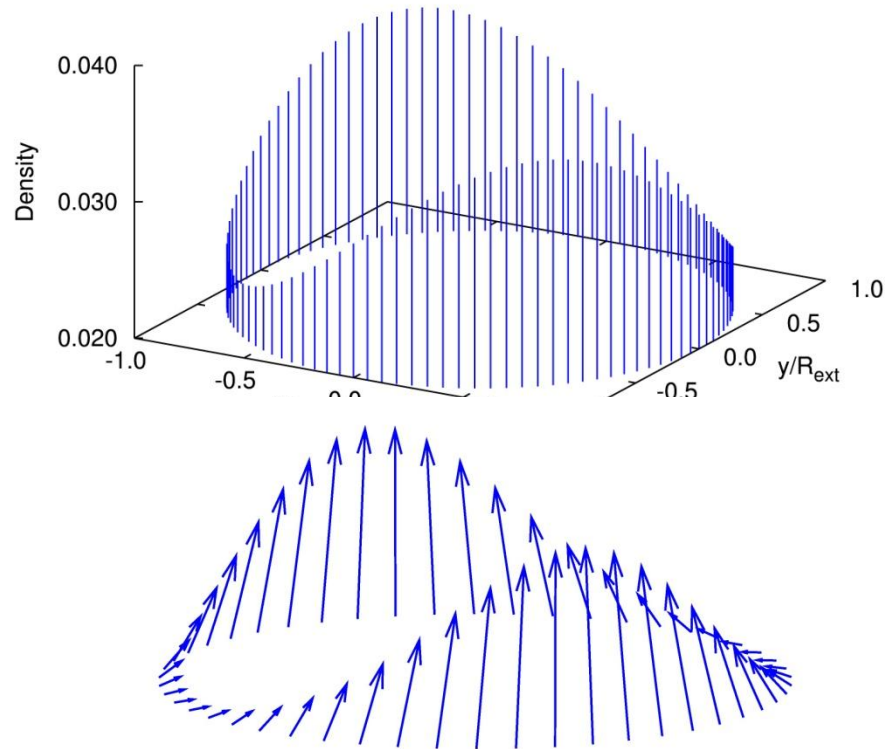


N=4 with interaction

**...but the spin polarization may also be amplified by the interaction ...**

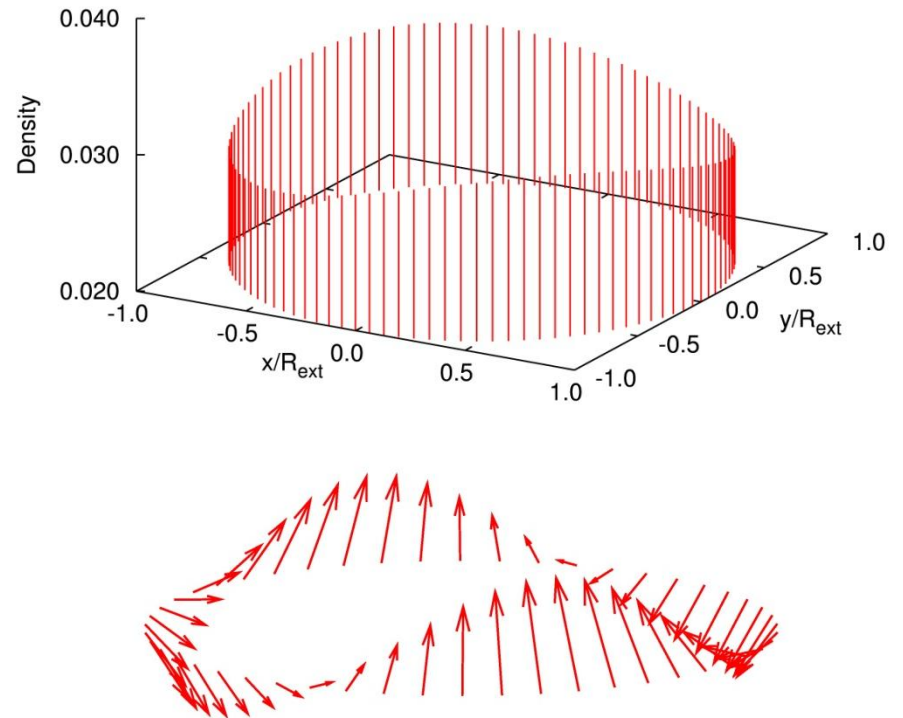
# Combined Rashba and Dresselhaus SOI in 1D

Electron density in the ground state 1D no interaction



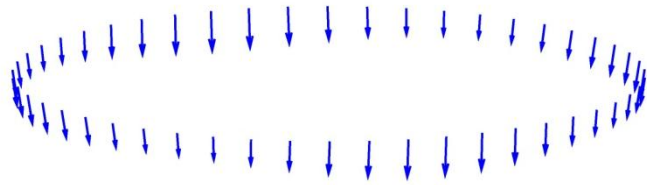
$N=3$ , no interaction,  
inhomogeneous charge density  
and spin density

Electron density in the ground state 1D with interaction

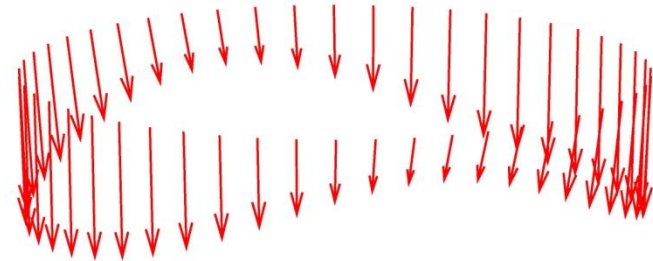


$N=3$ , Coulomb interaction,  
screened charge density,  
unscreened spin density?

# Combined Rashba and Dresselhaus SOI



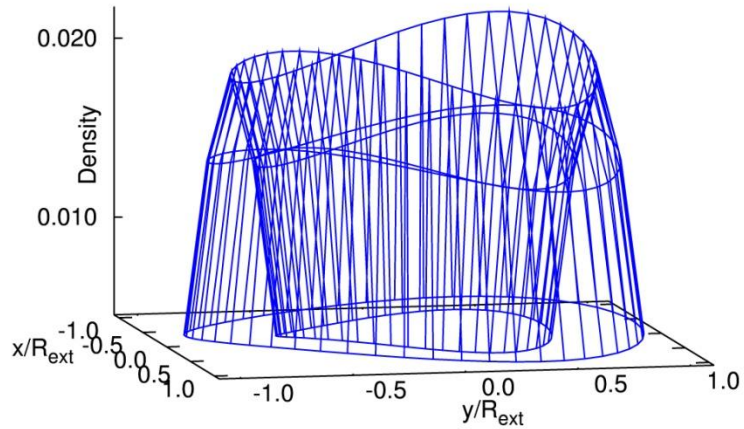
N=4, no interaction,  
inhomogeneous spin density



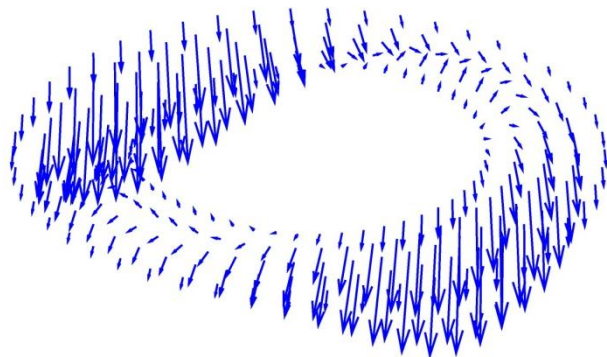
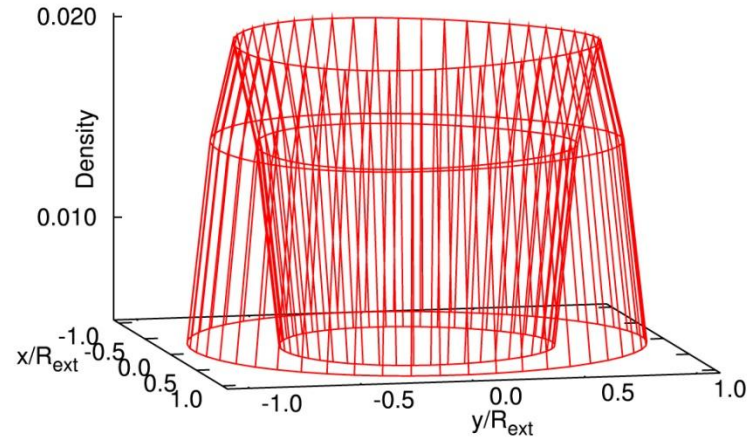
N=4, Coulomb interaction,  
unscreened spin density?

# Combined Coulomb and SOI in 2D

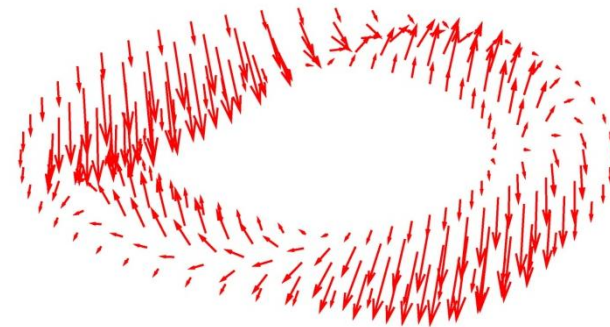
Electron density in the ground state 2D, no interaction



Electron density in the ground state 2D, with interaction



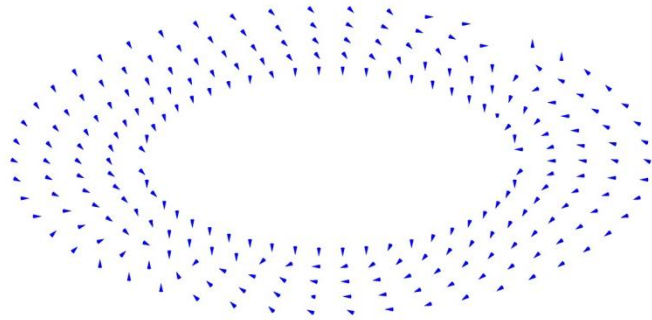
N=3 no interaction



N=3 with interaction

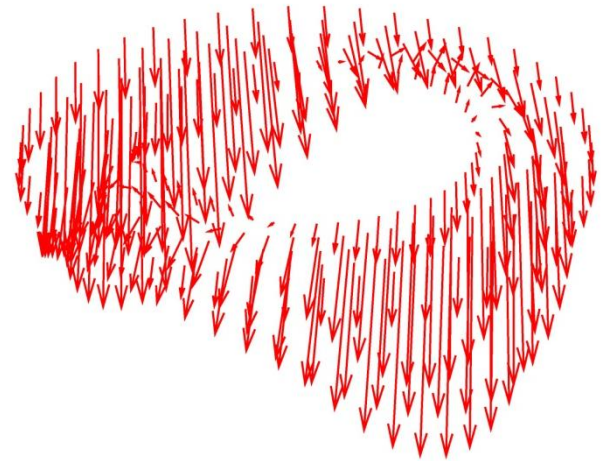
# Combined Coulomb and SOI in 2D

2D ring  $u_c=0$   $\alpha_R=1$   $\beta_D=0.8$



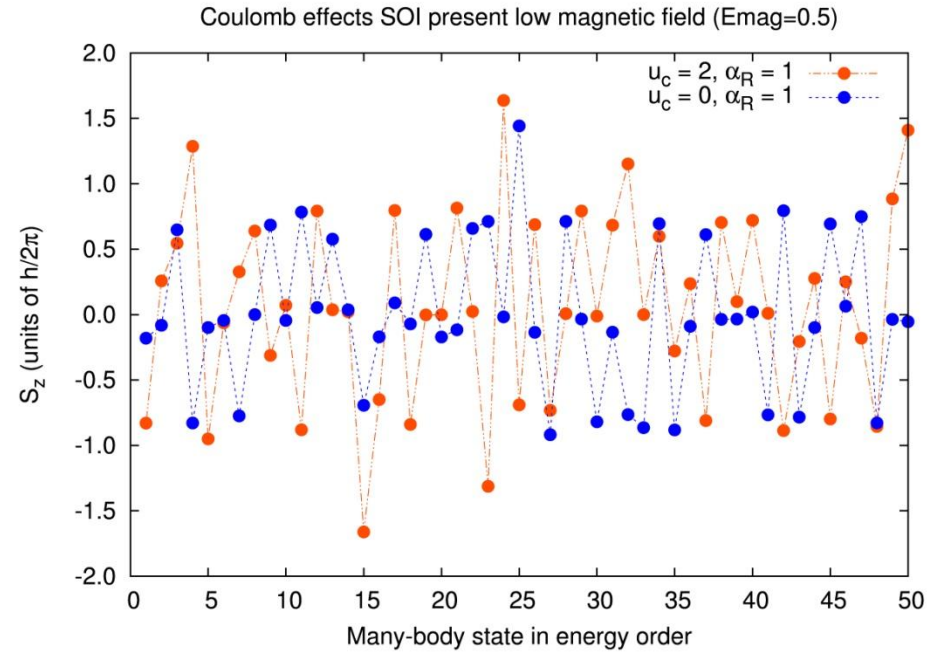
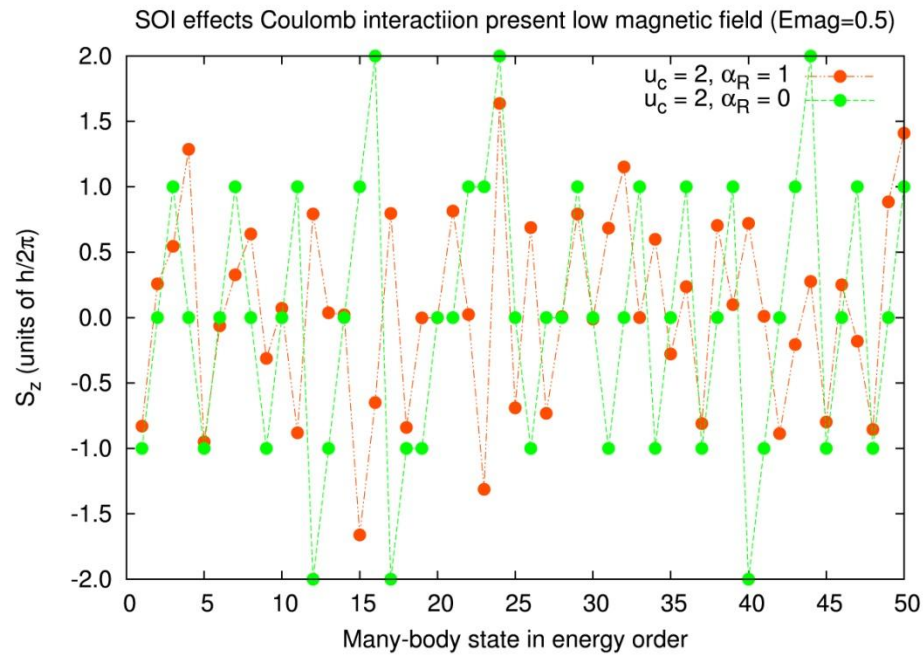
N=4 no interaction

2D ring  $u_c=2$   $\alpha_R=1$   $\beta_D=0.8$



N=4 with interaction

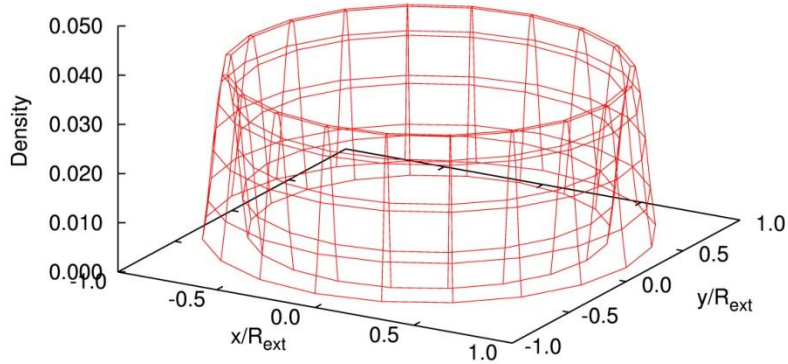
# Total spin and excited states



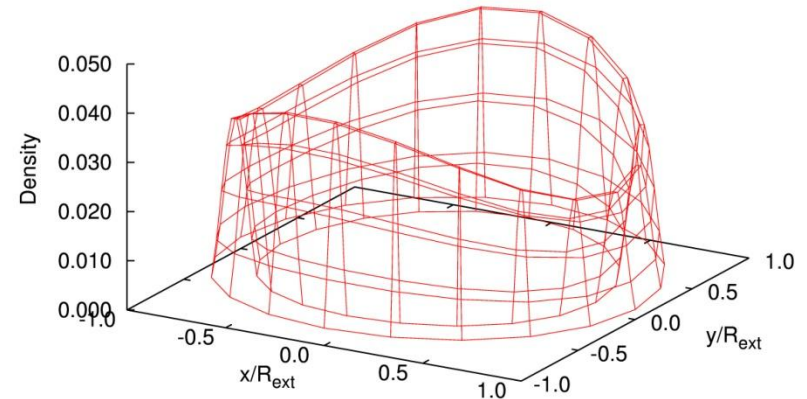


# Excited states in the presence of an external field. “Collective modes” ?

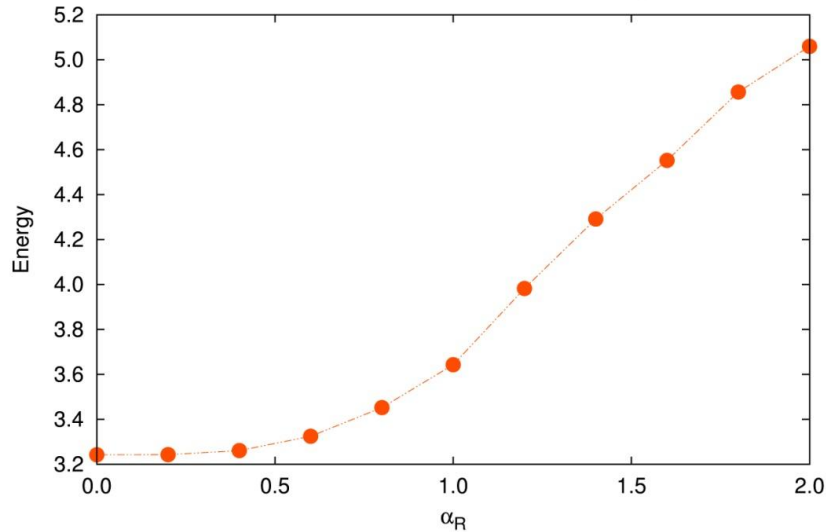
Density of electron in the ground state  $\alpha_R=1.0$



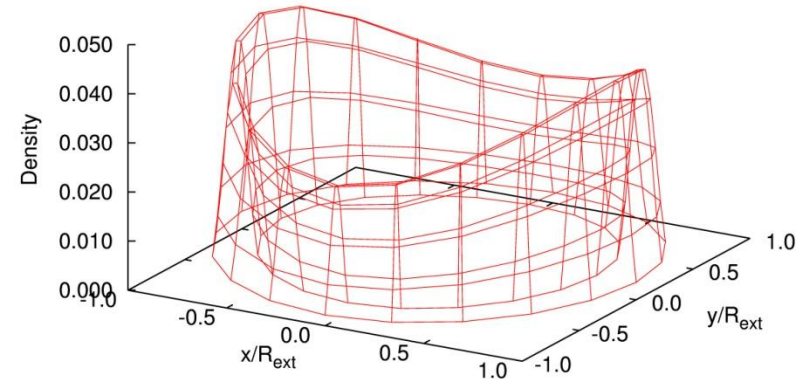
Density of electron in the state 23  $\alpha_R=1.0$



Collective excitations for different SOI strength



Density of electron in the state 24  $\alpha_R=1.0$



Increased energy of the “collective modes” with the SOI parameter.

For plasma oscillations: **D.C. Marinescu and F.Lung, PRB 82, 205322 (2010)**

# Continuous 2D ring model

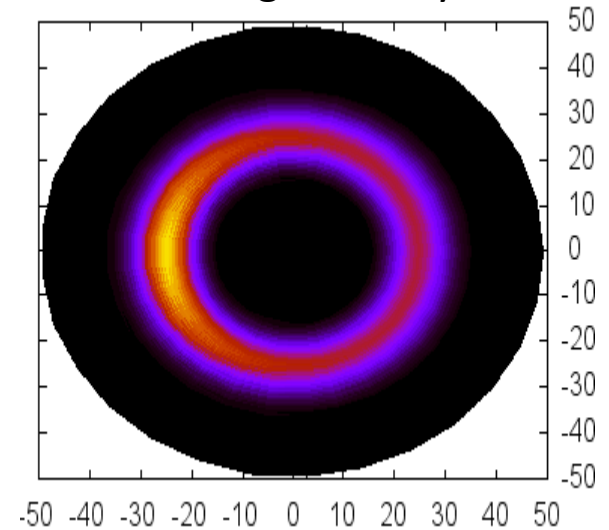
Lateral confinement  $V(r) = V_0 \left( \frac{r}{r_0} - \frac{r_0}{r} \right)$

Analytical single-particle wave functions

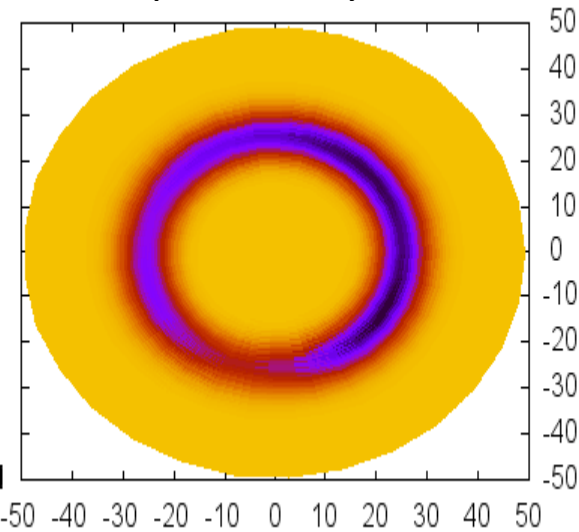
**Tan & Inckson, Semic. Sci. Technol., 11 1635 (1996)**

We include a Coulomb impurity

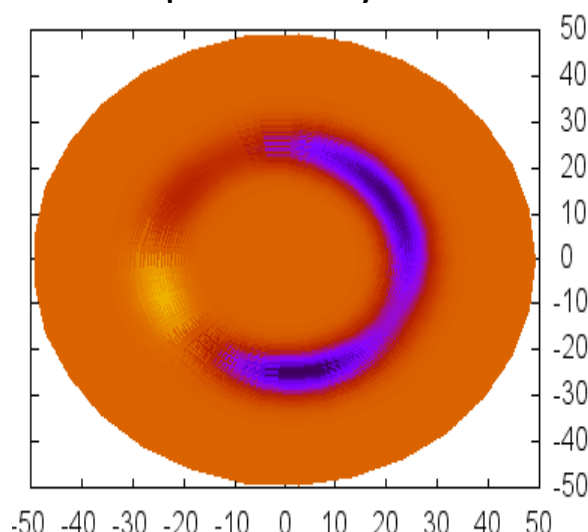
Charge density



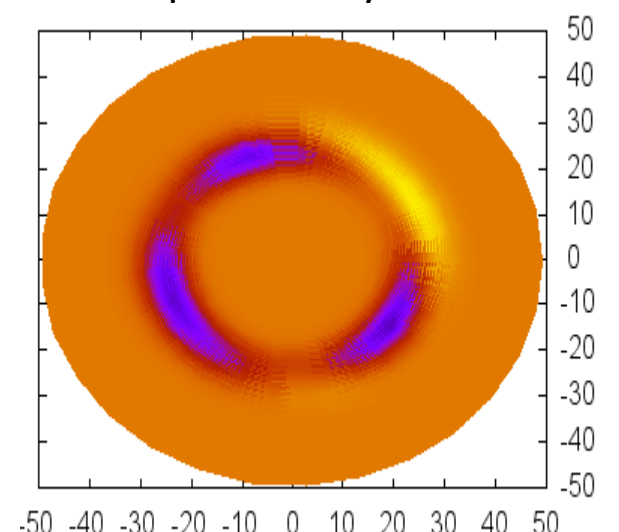
Spin density GS



Spin density ES1



Spin density ES2



# (Preliminary) conclusions

- The Coulomb interaction has strong effects on the spin polarization
- The spin density may propagate on longer distances than the charge density
- The spin density may have a richer structure in space than the charge density
- Collective modes have higher energy due to SOI ?

# Co-workers

- Csaba Daday (Reykjavik)
- Marian Niță (Bucharest)
- Catalina Marinescu (Clemson)
- Viðar Guðmundsson (Reykjavik)