## Abstract for MSH Sept 2012 Using gradient structures for modeling semiconductors Alexander Mielke (joint work with Annegret Glitzky and Matthias Liero) WIAS Berlin

We show that classical energy-drift-diffusion systems used for modeling of semiconductor devices can be written as gradient systems in the form

 $\partial_t \mathbf{u} = -\mathcal{K}(\mathbf{u}) \mathrm{D}\Phi(\mathbf{u}),$ 

where  $\mathbf{u}$  is a vector of densities of the different species and possibly the temperature (or the internal energy), cf. [Mie11]. While the potential  $\Phi$  is the free energy in the isothermal case and the negative entropy in the temperature-dependent case, the Onsager operator  $\mathcal{K}(\mathbf{u}) = G(\mathbf{u})^{-1}$  is the inverse of the metric tensor  $\mathcal{G}$  and has an additive split into the different dissipation mechanics:

$$\mathcal{K}(\mathbf{u})\boldsymbol{\xi} = \mathcal{K}_{ ext{diff}}(\mathbf{u})\boldsymbol{\xi} + \mathcal{K}_{ ext{react}}(\mathbf{u})\boldsymbol{\xi} = - ext{div}(\mathbb{M}(\mathbf{u})
abla \boldsymbol{\xi}) + \mathbb{K}(\mathbf{u})\boldsymbol{\xi}.$$

In the case of Fermi-Dirac statistics this modeling lead to new and more consistent modeling of diffusion enhancement and reaction kinetics. In particular, the modeling of semiconductors via gradient system  $(\Phi, \mathcal{K})$  yields always thermodynamically consistent models and highlights the different physical aspects of the theory in a clear way.

It is shown in [GlM12, Mie12] that gradient structures can also used effectively to model hybrid models with nontrivial bulk-interface interaction and active interfaces, which are need in photovoltaics. Finally, we discuss how the theory of geodesically  $\lambda$ -convex gradient can be applied in this field, cf. [LiM12].

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