



A thermodynamical theory for nonequilibrium systems

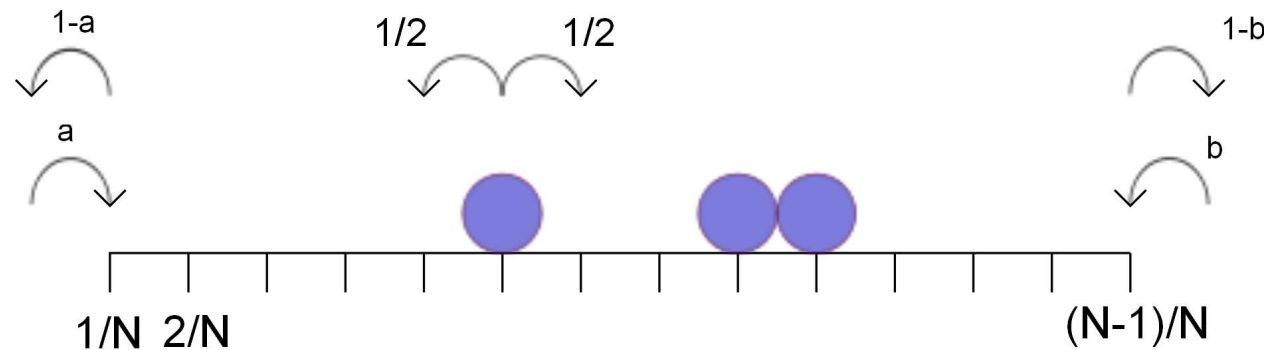
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From particle systems to differential equations

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EP@

- $\Lambda \subset \mathbb{R}^d \quad N \geq 1 \quad \{x/N \in \Lambda : x \in \mathbb{Z}^d\}$
- $\lambda : \partial\Lambda \rightarrow \mathbb{R} \quad \text{chemical potential}$
- Creation: $\frac{e^{\lambda(x/N)}}{1 + e^{\lambda(x/N)}}$ annihilation $\frac{1}{1 + e^{\lambda(x/N)}}$
- $\eta_t = \{\eta_t(x/N) : x/N \in \Lambda\}$
- ν_λ^N stationary state



Nonequilibrium free energy

- $\pi^N = \frac{1}{N^d} \sum_{x/N \in \Lambda} \eta(x/N) \delta_{x/N}$

- $\nu_\lambda^N \{ \pi^N \approx \bar{\rho} \} \sim 1$

$$\begin{cases} \nabla \cdot D(\rho) \nabla \rho = 0 \\ f'(\rho(x)) = \lambda(x) \quad x \in \partial\Lambda \end{cases}$$

- $D(\rho)$ diffusion coefficient $D(\rho) = (1/2)I$

- f equilibrium specific free energy $f(a) = a \log a + (1 - a) \log(1 - a)$

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- $\gamma \neq \bar{\rho}$

- $\nu_\lambda^N \{ \pi^N \approx \gamma \} \sim e^{-N^d S_\lambda(\gamma)}$

- $S_\lambda(\cdot)$ nonequilibrium free energy

Equilibrium \times Nonequilibrium

- λ constant ν_λ^N no correlations

$$S_\lambda(\gamma) = \int_\Lambda \left\{ \gamma \log \frac{\gamma}{\bar{\rho}} + [1 - \gamma] \log \frac{1 - \gamma}{1 - \bar{\rho}} \right\} dx$$

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- $d = 1$ $\Lambda = (0, 1)$ $\lambda(0) \neq \lambda(1)$ $\nu_\lambda^N[\eta(x/N); \eta(y/N)] = O(N^{-1})$

- Derrida, Lebowitz, Speer (2002) Bertini, De Sole, Gabrielli, Jona-Lasinio, L. (2002)

$$S_\lambda(\gamma) = \int_0^1 \left\{ \gamma \log \frac{\gamma}{F} + [1 - \gamma] \log \frac{1 - \gamma}{1 - F} + \log \frac{F_x}{\beta - \alpha} \right\} dx$$

$$\begin{cases} \frac{F_{xx}}{(F_x)^2} = \frac{\gamma - F}{F(1 - F)}, \\ F(x) = e^{\lambda(x)} / [1 + e^{\lambda(x)}], \quad x \in \partial\Lambda \end{cases}$$

(DLS)

Hydrodynamic limit

- $\pi_t^N = \frac{1}{N^d} \sum_{x/N \in \Lambda} \eta_{tN^2}(x/N) \delta_{x/N}$

- $\pi_0^N \rightarrow \gamma(x) dx$

- De Masi, Presutti et al. 80's Guo, Papanicolaou, Varadhan 86

- $\mathbb{P}_\gamma [\pi^N \approx w, [0, T]] \sim 1$

$$\begin{cases} \partial_t w = \nabla \cdot D(w) \nabla w \\ f'(w(t, x)) = \lambda(x) & x \in \partial\Lambda \\ w(0, \cdot) = \gamma(\cdot) \end{cases}$$

- Stationary solution $\bar{\rho}$ globally attractive

Action Functional

• Kipnis, Olla, Varadhan (1989), Donsker, Varadhan (1989), Quastel, Rezakhanlou, Varadhan (1999), Bertini, L., Mourragui (2010)

• $u(t, \cdot) \quad t \in [-T, 0]$

• $\mathbb{P}_{u(-T)} [\pi^N \approx u, [-T, 0]] \sim e^{-N^d I_{[-T, 0]}(u)}$

• $K(\rho)H = -\nabla \cdot (\chi(\rho)\nabla H) \quad H(x) = 0 \quad x \in \partial\Lambda$

• mobility: $\chi(a) = a(1 - a)I$

$$I_{[-T, 0]}(u) := \int_{-T}^0 dt \int_{\Lambda} dx [\partial_t u - \nabla \cdot D(u)\nabla u] K(u)^{-1} [\partial_t u - \nabla \cdot D(u)\nabla u]$$

Quasi-potential

- $I_{[-T,0]}(u)$

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- **Fix** γ $\inf_{\substack{u(-T)=\bar{\rho} \\ u(0)=\gamma}} I_{[-T,0]}(u)$

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- Bodineau and Giacomin (2004), Farfan (2010)

- V_λ is the nonequilibrium free energy:

- $\nu_\lambda^N \{\pi^N \approx \gamma\} \sim e^{-NV_\lambda(\gamma)}$

- Holds for large class of systems in any dimension

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- V_λ is the nonequilibrium free energy:
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- Holds for large class of systems in any dimension
- DLS, BDGJL $d = 1$, few models, Explicit formula for $V_\lambda(\gamma)$



1. Transformations



Basic assumptions

1. Local density $u(t, x)$ and current $j = j(t, x)$

$$\partial_t u + \nabla \cdot j = 0$$

2. Constitutive equation $j = J(t, u(t))$:

$$J(t, \rho) = -D(\rho) \nabla \rho + \chi(\rho) E(t) ,$$

- $D(\rho)$ diffusion coefficient $\chi(\rho)$ mobility

3. Boundary condition

$$f'(u(t, x)) = \lambda(t, x) , \quad x \in \partial\Lambda .$$

- f is the equilibrium specific free energy

4. Einstein relation $D(\rho) = \chi(\rho) f''(\rho)$

Energy balance

$$W_{[0,T]} = \int_0^T dt \left\{ - \int_{\partial\Omega} d\sigma(x) \lambda(t, x) j(t, x) \cdot \hat{n}(x) + \int_{\Omega} dx j(t, x) \cdot E(t, x) \right\} ,$$

- Second law of thermodynamics (Clausius inequality)

$$W_{[0,T]}[\lambda, E, \rho] \geq F(u(T)) - F(\rho) ,$$

- F equilibrium free energy:

$$F(\rho) = \int_{\Omega} dx f(\rho(x))$$

Energy balance

$$W_{[0,T]} = \int_0^T dt \left\{ - \int_{\partial\Lambda} d\sigma(x) \lambda(t, x) j(t, x) \cdot \hat{n}(x) + \int_{\Lambda} dx j(t, x) \cdot E(t, x) \right\} ,$$

• $\lambda = f'(u(t)) \quad D(\rho) = \chi(\rho) f''(\rho) \quad J(t, \rho) = -D(\rho) \nabla \rho + \chi(\rho) E(t)$

$$\begin{aligned} W_{[0,T]} &= \int_0^T dt \left\{ - \int_{\partial\Lambda} d\sigma f'(u(t)) j(t) \cdot \hat{n} + \int_{\Lambda} dx j(t) \cdot E(t) \right\} , \\ &= \int_0^T dt \int_{\Lambda} dx \left\{ - \nabla \cdot [f'(u(t)) j(t)] + j(t) \cdot E(t) \right\} \\ &= \int_0^T dt \int_{\Lambda} dx \left[- f'(u(t)) \nabla \cdot j(t) - f''(u(t)) \nabla u(t) \cdot j(t) + j(t) \cdot E(t) \right] \\ &= \int_0^T dt \frac{d}{dt} \int_{\Lambda} dx f(u(t)) + \int_0^T dt \int_{\Lambda} dx j(t) \cdot \chi(u(t))^{-1} j(t) , \end{aligned}$$



2. Equilibrium states: $J(u) = 0$



Reversible and quasi static transformations

- $E = 0$
- Spatially homogenous equilibrium $\lambda = \text{cte}$
- $\bar{\rho}_\lambda = \text{cte}$ $f'(\bar{\rho}_\lambda) = \lambda$

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- Spatially homogenous equilibrium $\lambda = \text{cte}$
- $\bar{\rho}_\lambda = \text{cte} \quad f'(\bar{\rho}_\lambda) = \lambda$
- $\lambda_0 \quad \lambda_1$
- $\bar{\rho}_0 = \bar{\rho}_{\lambda_0} \longrightarrow \bar{\rho}_1 = \bar{\rho}_{\lambda_1}$
- $\lambda(t) = \lambda_0, t \leq 0 \quad \lambda(t) = \lambda_1, t \geq T$

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- Spatially homogenous equilibrium $\lambda = \text{cte}$
- $\bar{\rho}_\lambda = \text{cte}$ $f'(\bar{\rho}_\lambda) = \lambda$
- λ_0 λ_1
- $\bar{\rho}_0 = \bar{\rho}_{\lambda_0} \longrightarrow \bar{\rho}_1 = \bar{\rho}_{\lambda_1}$
- $\lambda(t) = \lambda_0, t \leq 0$ $\lambda(t) = \lambda_1, t \geq T$
- *reversible transformation*: energy exchanged is minimal
- *quasi static transformation*: variation of the chemical potential is very slow

Reversible and quasi static transformations

$$\begin{aligned} W &= \int_0^\infty dt \frac{d}{dt} \int_\Lambda dx f(u(t)) + \int_0^\infty dt \int_\Lambda dx j(t) \cdot \chi(u(t))^{-1} j(t) \\ &= F(\bar{\rho}_1) - F(\bar{\rho}_0) + \int_0^\infty dt \int_\Lambda dx j(t) \cdot \chi(u(t))^{-1} j(t) \end{aligned}$$

- No regularity assumption of the chemical potential in time

Reversible and quasi static transformations

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• No regularity assumption of the chemical potential in time

• Smooth $\lambda(t)$ $\lambda_\delta(t) = \lambda(\delta t)$ $u_\delta(t)$

• $J(t, \rho) = J(\rho) = -D(\rho) \nabla \rho$ $D(\rho) = \chi(\rho) f''(\rho)$

$$W = F(\bar{\rho}_1) - F(\bar{\rho}_0) + \int_0^\infty dt \int_\Lambda dx \nabla f'(u_\delta(t)) \cdot \chi(u_\delta(t)) \nabla f'(u_\delta(t)) ,$$

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$$W = F(\bar{\rho}_1) - F(\bar{\rho}_0) + \int_0^\infty dt \int_\Lambda dx \nabla f'(u_\delta(t)) \cdot \chi(u_\delta(t)) \nabla f'(u_\delta(t))$$

- $\bar{\rho}_{\lambda_\delta(t)} \quad \lambda_\delta(t)$

- $\nabla f'(\bar{\rho}_{\lambda_\delta(t)}) = 0$

$$\int_0^\infty dt \int_\Lambda dx \nabla [f'(u_\delta(t)) - f'(\bar{\rho}_{\lambda_\delta(t)})] \cdot \chi(u_\delta(t)) \nabla [f'(u_\delta(t)) - f'(\bar{\rho}_{\lambda_\delta(t)})]$$

- $u_\delta(t) - \bar{\rho}_{\lambda_\delta(t)} = O(\delta)$

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- $u_\delta(t) - \bar{\rho}_{\lambda_\delta(t)} = O(\delta)$

- Reversible transformation $W = \Delta F$

- No special property of $\lambda(t)$

Excess work

- $\lambda(0) = \lambda_0 \quad \bar{\rho}_0$

- Transformation $\lambda(t) \quad \lambda(t) \longrightarrow \lambda_1 \quad \bar{\rho}_1$

- *excess work:*

$$W_{\text{ex}} = W[\lambda, E, \rho] - \min W = \int_0^\infty dt \int_\Lambda dx j(t) \cdot \chi(u(t))^{-1} j(t)$$

- $j(t) = J(u(t)) = -D(u)\nabla u \quad D(\rho) = \chi(\rho) f''(\rho)$

$$W_{\text{ex}} = - \int_0^\infty dt \int_\Lambda dx \nabla f'(u(t)) \cdot J(u(t))$$

Excess work and quasi-potential

• Relaxation path: $(\lambda_0, \bar{\rho}_0) \rightarrow \lambda_1$ for $t > 0$ $u(t) \rightarrow \bar{\rho}_1$

$$W_{\text{ex}}[\lambda_1, \bar{\rho}_0] = - \int_0^\infty dt \int_\Lambda dx \nabla f'(u(t)) \cdot J(u(t))$$

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- $\nabla f'(\bar{\rho}_1) = 0$

$$\begin{aligned} W_{\text{ex}}[\lambda_1, \bar{\rho}_0] &= \int_0^\infty dt \int_\Lambda dx [f'(u(t)) - f'(\bar{\rho}_1)] \nabla \cdot J(u(t)) \\ &= - \int_0^\infty dt \int_\Lambda dx [f'(u(t)) - f'(\bar{\rho}_1)] \partial_t u(t) \\ &= \int_\Lambda dx [f(\bar{\rho}_0) - f(\bar{\rho}_1) - f'(\bar{\rho}_1)(\bar{\rho}_0 - \bar{\rho}_1)] = V_{\lambda_1}(\bar{\rho}_0) \end{aligned}$$

- W_{ex} is not the difference of a thermodynamic potential



3. Nonequilibrium states



Adjoint hydrodynamics

- $\partial_t u + \nabla \cdot J(u(t)) = 0$
- Adjoint hydrodynamics chemical potential λ fixed
- $\partial_t u + \nabla \cdot J^*(u(t)) = 0$
- $\frac{1}{2}\{J(\rho) + J^*(\rho)\} = -\chi(\rho) \nabla \frac{\delta V_\lambda(\rho)}{\delta \rho}$

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- $\frac{1}{2}\{J(\rho) + J^*(\rho)\} = -\chi(\rho) \nabla \frac{\delta V_\lambda(\rho)}{\delta \rho}$

- $J_S^\lambda(\rho) = \frac{1}{2}\{J(\rho) + J^*(\rho)\} = -\chi(\rho) \nabla \frac{\delta V_\lambda(\rho)}{\delta \rho}$

- $J_A^\lambda(\rho) = \frac{1}{2}\{J(\rho) - J^*(\rho)\} = J(\rho) - J_S^\lambda(\rho)$

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- $J_A^\lambda(\rho) = \frac{1}{2}\{J(\rho) - J^*(\rho)\} = J(\rho) - J_S^\lambda(\rho)$

- $\int_\Lambda dx J_S^\lambda(\rho) \cdot \chi(\rho)^{-1} J_A^\lambda(\rho) = 0$

Work to maintain stationary state

• $\lambda \quad u(t) = \bar{\rho}_\lambda \quad 0 \leq t \leq T$

$$\begin{aligned} W_{[0,T]} &= \int_0^T dt \frac{d}{dt} \int_\Lambda dx f(u(t)) + \int_0^T dt \int_\Lambda dx j(t) \cdot \chi(u(t))^{-1} j(t) \\ &= \int_0^T dt \int_\Lambda dx J(u(t)) \cdot \chi(u(t))^{-1} J(u(t)) \end{aligned}$$

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• $J(\rho) = J_S^\lambda(\rho) + J_A^\lambda(\rho)$

• $J_S^\lambda(\rho) = -\chi(\rho) \nabla \frac{\delta V_\lambda(\rho)}{\delta \rho}$

• $\frac{\delta V_\lambda(\bar{\rho}_\lambda)}{\delta \rho} = 0 \quad J_S^\lambda(\bar{\rho}_\lambda) = 0$

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• $\frac{\delta V_\lambda(\bar{\rho}_\lambda)}{\delta \rho} = 0 \quad J_S^\lambda(\bar{\rho}_\lambda) = 0$

$$W_{[0,T]} = T \int_\Lambda dx J_A^\lambda(\bar{\rho}_\lambda) \cdot \chi(\bar{\rho}_\lambda)^{-1} J_A^\lambda(\bar{\rho}_\lambda)$$

Renormalized work

• Fix $T > 0$ profile ρ ch. pot. $\lambda(t)$

• $u(t) \quad j(t) \quad t \geq 0$

$$W_{[0,T]}^{\text{ren}}[\lambda, \rho] = W_{[0,T]}[\lambda, \rho] - \int_0^T dt \int_{\Lambda} dx J_{\Lambda}^{\lambda(t)}(u(t)) \cdot \chi(u(t))^{-1} J_{\Lambda}^{\lambda(t)}(u(t))$$

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• Orthogonality $J_{\text{S}}^{\lambda(t)}(u(t)), J_{\text{A}}^{\lambda(t)}(u(t))$

$$W_{[0,T]}^{\text{ren}}[\lambda, \rho] = F(u(T)) - F(\rho) + \int_0^T dt \int_{\Lambda} dx J_{\text{S}}^{\lambda(t)}(u(t)) \cdot \chi(u(t))^{-1} J_{\text{S}}^{\lambda(t)}(u(t))$$

Clausius inequality

$$W_{[0,T]}^{\text{ren}}[\lambda, \rho] = F(u(T)) - F(\rho) + \int_0^T dt \int_{\Lambda} dx J_S^{\lambda(t)}(u(t)) \cdot \chi(u(t))^{-1} J_S^{\lambda(t)}(u(t))$$

• $\lambda(t) \rightarrow \lambda_1$

• $J_S^{\lambda(t)}(u(t)) \rightarrow J_S^{\lambda_1}(\bar{\rho}_1) = 0$

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Quasi static transformations

- $\lambda(0) = \lambda_0 \quad \rho(0) = \bar{\rho}_{\lambda_0}$
- $\lambda(t) = \lambda_1 \quad t \geq T$
- $\delta > 0 \quad \lambda_\delta(t) = \lambda(\delta t) \quad (u_\delta(t), j_\delta(t))$

$$\int_0^\infty dt \int_\Lambda dx J_S^{\lambda_\delta(t)}(u_\delta(t)) \cdot \chi(u_\delta(t))^{-1} J_S^{\lambda_\delta(t)}(u_\delta(t))$$

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- $\bar{\rho}_\delta(t) = \bar{\rho}_{\lambda_\delta(t)}$

- $J_S^{\lambda_\delta(t)}(\bar{\rho}_\delta(t)) = 0$

$$\int_0^\infty dt \int_\Lambda dx [J_S^{\lambda_\delta(t)}(u_\delta(t)) - J_S^{\lambda_\delta(t)}(\bar{\rho}_\delta(t))] \cdot \chi(u_\delta(t))^{-1} [J_S^{\lambda_\delta(t)}(u_\delta(t)) - J_S^{\lambda_\delta(t)}(\bar{\rho}_\delta(t))]$$

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$$W^{\text{ren}} = \Delta F = F(\bar{\rho}_1) - F(\bar{\rho}_0)$$

Excess work

- $\lambda(t) \longrightarrow \lambda_1$
- initial density profile ρ .

$$\begin{aligned} W_{\text{ex}}[\lambda, \rho] &= W^{\text{ren}}[\lambda, \rho] - \min W^{\text{ren}}[\lambda', \rho] \\ &= \int_0^\infty dt \int_\Lambda dx J_S^{\lambda(t)}(u(t)) \cdot \chi(u(t))^{-1} J_S^{\lambda(t)}(u(t)) \end{aligned}$$

Relaxation path: excess work and quasi potential

• $(\lambda_0, \bar{\rho}_0) \rightarrow \lambda_1 \quad t \geq 0$

• $u(t) \rightarrow \bar{\rho}_1$

$$\begin{aligned} W_{\text{ex}}[\lambda_1, \bar{\rho}_0] &= \int_0^\infty dt \int_\Lambda dx J_S^{\lambda_1}(u(t)) \cdot \chi(u(t))^{-1} J_S^{\lambda_1}(u(t)) \\ &= \int_0^\infty dt \int_\Lambda dx J^{\lambda_1}(u(t)) \cdot \chi(u(t))^{-1} J_S^{\lambda_1}(u(t)) \end{aligned}$$

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• $(\lambda_0, \bar{\rho}_0) \rightarrow \lambda_1 \quad t \geq 0$

• $u(t) \rightarrow \bar{\rho}_1$

$$\begin{aligned} W_{\text{ex}}[\lambda_1, \bar{\rho}_0] &= \int_0^\infty dt \int_\Lambda dx J_S^{\lambda_1}(u(t)) \cdot \chi(u(t))^{-1} J_S^{\lambda_1}(u(t)) \\ &= \int_0^\infty dt \int_\Lambda dx J^{\lambda_1}(u(t)) \cdot \chi(u(t))^{-1} J_S^{\lambda_1}(u(t)) \end{aligned}$$

$$\int_0^\infty dt \int_\Lambda dx \nabla \cdot J^{\lambda_1}(u(t)) \frac{\delta V_{\lambda_1}(u(t))}{\delta \rho} = - \int_0^\infty dt \int_\Lambda dx \partial_t u(t) \frac{\delta V_{\lambda_1}(u(t))}{\delta \rho}$$

Relaxation path: excess work and quasi potential

• $(\lambda_0, \bar{\rho}_0) \rightarrow \lambda_1 \quad t \geq 0$

• $u(t) \rightarrow \bar{\rho}_1$

$$\begin{aligned} W_{\text{ex}}[\lambda_1, \bar{\rho}_0] &= \int_0^\infty dt \int_\Lambda dx J_S^{\lambda_1}(u(t)) \cdot \chi(u(t))^{-1} J_S^{\lambda_1}(u(t)) \\ &= \int_0^\infty dt \int_\Lambda dx J^{\lambda_1}(u(t)) \cdot \chi(u(t))^{-1} J_S^{\lambda_1}(u(t)) \end{aligned}$$

$$\int_0^\infty dt \int_\Lambda dx \nabla \cdot J^{\lambda_1}(u(t)) \frac{\delta V_{\lambda_1}(u(t))}{\delta \rho} = - \int_0^\infty dt \int_\Lambda dx \partial_t u(t) \frac{\delta V_{\lambda_1}(u(t))}{\delta \rho}$$

$$W_{\text{ex}}[\lambda_1, \bar{\rho}_0] = V_{\lambda_1}(\bar{\rho}_0) - V_{\lambda_1}(\bar{\rho}_1) = V_{\lambda_1}(\bar{\rho}_0)$$