## Diffusive interface dynamics in lattices

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Microscopic regularizations of ill-posed forward-backward diffusion equations have many applications and provide dynamical models for diffusive interfaces. Typical examples are the Cahn-Hilliard equation, the viscous approximations studied in [1, 4], and spatially discrete gradient systems such as

$$\dot{u}_j = \Delta p_j, \qquad p_j = \Phi'(u_j), \qquad j \in \mathbb{Z},$$

where  $\Delta$  denotes the discrete Laplacian and  $\Phi'$  is the bistable derivative of a double-well potential. While the Cahn-Hilliard case is well understood, very little is known about the mathematical analysis of viscous or lattice regularizations.

In this talk we choose  $\Phi'$  as a trilinear function and demonstrate that the dynamics of both moving and standing phase interface is intimately related to a family of entropy inequalities and a hysteretic Stefan condition, see the figure and [2] for a simpler case.



**Figure.** Cartoon of the nonlinearity  $\Phi'$  and the macroscopic hysteresis loop.

**Theorem** (macroscopic evolution in the parabolic scaling limit). *Identifying* 

$$u_j(t) = U(\tau, \xi), \qquad p_j(t) = P(\tau, \xi), \qquad \tau = \varepsilon^2 t, \qquad \xi = \varepsilon j,$$

the limit  $\varepsilon \to 0$  for well-prepared single-interface initial data is governed by

$$\partial_{\tau} U = \partial_{\xi}^2 P, \qquad U = P + \mathcal{M}[P],$$

where the relay operator  $\mathcal{M}$  acts pointwise in  $\xi$  on temporal functions. In particular, there exists a unique interface curve  $\xi = \xi_*(\tau)$  which separates space-time regions with either  $U \leq -u_*$  or  $U \geq +u_*$ .

The key arguments in the proof are derived by a careful analysis of the lattice ODE and can informally be summarized as follows (see [3] for the details):

- (1) Macroscopic interfaces exist according to a *persistence lemma*.
- (2) The *waiting lemma* provides an upper bound for the interface speed.
- (3) Microscopic phase transitions can be characterized by a *slow-fast splitting*.
- (4) Several *fluctuation estimates* control the impact of the mesoscopic oscillations.

## References

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