



On the error estimates for space-time discretizations of rate-independent processes

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joint work with A. Mielke, L. Paoli and U. Stefanelli

DFG Research Center MATHEON
Mathematics for key technologies

W I A S
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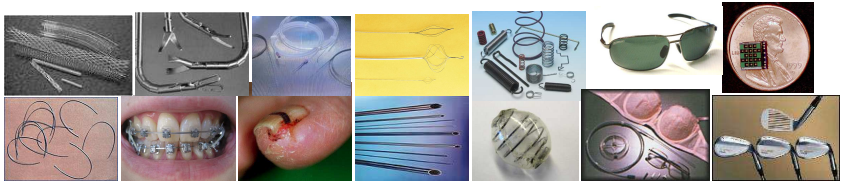
- 1 Background of SMAs
- 2 An abstract space-time discretization
- 3 Application to the isothermal Souza-Auricchio model
- 4 Conclusion



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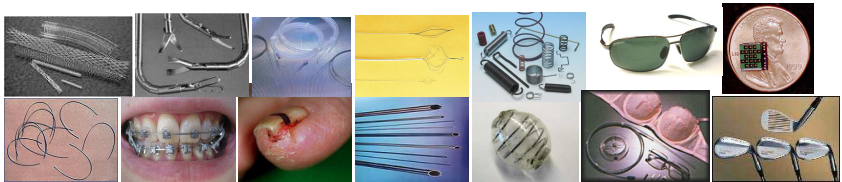
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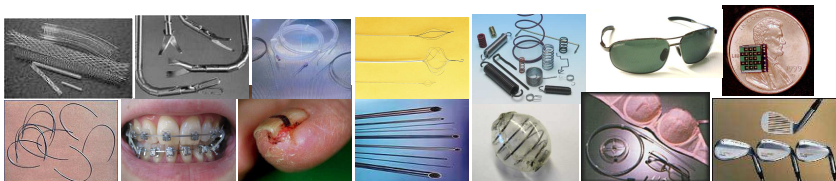


WHY? Because they have the following properties:

- ▷ shape memory under heating and cooling
- ▷ superelastic properties under mechanical loading
- ▷ hysteretic behavior for damping of vibrations



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AIM: Find good mathematical models (analysis and numerics)



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▷ **Energy functional:** $\mathcal{E}(t, q) \stackrel{\text{def}}{=} \frac{1}{2} \langle \mathbf{A}q, q \rangle_{\mathcal{Q}} + \underbrace{\mathcal{H}(q)}_{\text{hardening funct.}} - \underbrace{\langle \mathbf{I}(t), q \rangle}_{\text{load. funct.}}$

(E1) $\mathbf{A} \in \text{Lin}(\mathcal{Q}, \mathcal{Q}')$ is a symmetric operator
and there exists $\kappa > 0$: $\langle \mathbf{A}\hat{q}, \hat{q} \rangle_{\mathcal{Q}} \geq \kappa \|\hat{q}\|_{\mathcal{Q}}$ for all $\hat{q} \in \mathcal{Q}$

(E2) $\mathcal{H} \in C^3(\mathcal{Q}, \mathbb{R})$ is convex and $D_q \mathcal{H} \in C^0(\mathcal{Q}; \mathcal{X}')$

(E3) $\mathbf{I} \in C^3([0, T]; \mathcal{X}')$



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▷ **Dissipation functional:** $\Psi : \mathcal{Q} \rightarrow [0, \infty)$

(D1) $\Psi(\gamma q) = \gamma \Psi(q)$ for all $\gamma \geq 0, q \in \mathcal{Q}$

(D2) $\Psi(q_1 + q_2) \leq \Psi(q_1) + \Psi(q_2)$ for all $q_1, q_2 \in \mathcal{Q}$

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▷ \mathcal{X} : the Banach space such that $\mathcal{Q} \subset \mathcal{X} \subset \mathcal{Q}'$

▷ $\mathcal{S}(t) \stackrel{\text{def}}{=} \{q \in \mathcal{Q} \mid \mathcal{E}(t, q) \leq \mathcal{E}(t, \hat{q}) + \Psi(\hat{q} - q) \text{ for all } \hat{q} \in \mathcal{Q}\}$: the stable set



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$$(DI) \quad 0 \in \partial \Psi(\dot{q}(t)) + D_q \mathcal{E}(t, q(t)) \text{ a.e. in } [0, T]$$



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$$\text{(VI)} \quad \langle D_q \mathcal{E}(t, q(t)), v - \dot{q}(t) \rangle_{\mathcal{Q}} + \Psi(v) - \Psi(\dot{q}(t)) \geq 0 \text{ for all } v \in \mathcal{Q}$$



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(DI) is also equivalent to the *energetic formulation* (cf. Mielke&Theil'04)

$$(S) \quad \mathcal{E}(t, q(t)) \leq \mathcal{E}(t, \tilde{q}) + \Psi(\tilde{q} - q(t)) \text{ for all } \tilde{q} \in \mathcal{Q}$$

$$(E) \quad \mathcal{E}(t, q(t)) + \int_0^t \Psi(\dot{q}(s)) ds = \mathcal{E}(0, q(0)) + \int_0^t \partial_s \mathcal{E}(s, q(s)) ds$$



- ▷ Finite-element spaces $\mathcal{Q}_h \subset \mathcal{Q}$
- ▷ Time step $\tau > 0$

Space-Time Discretization

$$\text{(IP)}^{\tau,h} \quad q_k^{\tau,h} \in \underset{\hat{q}^h \in \mathcal{Q}_h}{\text{Argmin}} (\mathcal{E}(t_k^\tau, \hat{q}^h) + \psi(\hat{q}^h - q_{k-1}^{\tau,h}))$$



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Theorem (Mielke&Theil'04)

Assume that **(E1)**-**(E3)** and **(D1)**-**(D3)** hold. Then for all $h \geq 0$, there exists a unique solution $q_h : [0, T] \rightarrow \mathcal{Q}_h$ to **(DI)** and there exists $R > 0$ s.t. $\|q_h(t)\|_{\mathcal{Q}} \leq R$ for all $t \in [0, T]$ and $\|\dot{q}_h(t)\|_{\mathcal{Q}} \leq C(R, \kappa)$ for a.e. $t \in [0, T]$.



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AIM: evaluate $\|q^{\tau,h}(t) - q(t)\|_{\mathcal{Q}}$ by some polynomial function of τ and h



Let $\mathbf{P}_h : \mathcal{Q} \rightarrow \mathcal{Q}_h$ be a linear operator such that there exist $C_i^{\mathbf{P}} > 0$ and α_i , $i = 1, 2, 3$, such that for all $v \in \mathcal{Q}$ and $v_h \in \mathcal{Q}_h$, we have

$$\text{(P1)} \quad \|(\mathbf{P}_h - \text{Id})v\|_{\mathcal{X}} \leq C_1^{\mathbf{P}} h^{\alpha_1} \|v\|_{\mathcal{Q}}$$

$$\text{(P2)} \quad \|(\mathbf{P}_h^* \mathbf{A} - \mathbf{A} \mathbf{P}_h)v\|_{\mathcal{Q}'} \leq C_2^{\mathbf{P}} h^{\alpha_2} \|v\|_{\mathcal{Q}}$$

$$\text{(P3)} \quad \|(\mathbf{P}_h - \text{Id})v_h\|_{\mathcal{Q}} \leq C_3^{\mathbf{P}} h^{\alpha_3} \|v_h\|_{\mathcal{Q}}$$

Theorem

Assume that **(E1)-(E3)**, **(D1)-(D3)** hold. Then there exists $C > 0$ such that for all $h \in [0, h_0]$ and all $q_h(0) \in \mathcal{S}_h(0)$,

$$\|q_{\tau,h}(t) - q(t)\|_{\mathcal{Q}} \leq C (h^{\alpha/2} + \sqrt{\tau} + \|q_h(0) - q(0)\|_{\mathcal{Q}})$$

for all $t \in [0, T]$ and $\tau \in (0, T]$.

Sketch of proof: Let us remark that

$$\begin{aligned} \|q_{\tau,h}(t) - q(t)\|_{\mathcal{Q}} &\leq \underbrace{\|q_{\tau,h}(t) - q_h(t)\|_{\mathcal{Q}}}_{\leq C\sqrt{\tau} \text{ (cf. Mielke\&Theil'04)}} + \underbrace{\|q_h(t) - q(t)\|_{\mathcal{Q}}}_{\leq C(h^{\alpha/2} + \|q_h(0) - q(0)\|_{\mathcal{Q}})} \end{aligned}$$



Lemma

Assume that (E1)-(E3), (D1)-(D3) hold. Then we have

$$\begin{aligned}
 |\langle \mathbf{B}_h q, q \rangle_{\mathcal{Q}}| &\leq C^{\mathbf{B}} h^\alpha \text{ with } \mathbf{B}_h \stackrel{\text{def}}{=} -(\mathbf{A}(\mathbf{P}_h - \text{Id}) + (\mathbf{A}(\mathbf{P}_h - \text{Id}))^*) \\
 \langle \mathbf{D}_q \mathcal{E}(t, q_h), v_h - w \rangle_{\mathcal{Q}} + \Psi(v_h - w) &\leq C(h^\alpha + \|q_h - q\|_{\mathcal{Q}}^2) \|w\|_{\mathcal{Q}} + \langle \mathbf{B}_h q, w \rangle_{\mathcal{Q}} \\
 \text{with } \alpha &= \min\{\alpha_1, 2\alpha_2, \alpha_3\}.
 \end{aligned}$$

Proof: Notice that

$$\begin{aligned}
 &\langle \mathbf{D}_q \mathcal{E}(t, q_h), v_h - w \rangle_{\mathcal{Q}} + \Psi(v_h - w) \\
 &= \langle \mathbf{A} q_h + \mathbf{D}_q \mathcal{H}(q_h) - \mathbf{l}(t), v_h - w \rangle_{\mathcal{Q}} + \Psi(v_h - w) \\
 &\leq \langle \mathbf{A}(q_h - q), v_h - w \rangle_{\mathcal{Q}} + C C_1^{\mathbf{P}} h^{\alpha_1} \|w\|_{\mathcal{Q}}
 \end{aligned}$$

Taking $v_h = \mathbf{P}_h w$, we have

$$\begin{aligned}
 \langle \mathbf{A}(q_h - q), v_h - w \rangle_{\mathcal{Q}} &= \underbrace{\langle \mathbf{A}(q_h - q), (\mathbf{P}_h - \text{Id})w \rangle_{\mathcal{Q}}}_{=} \\
 &= \langle (\mathbf{P}_h^* \mathbf{A} - \mathbf{A} \mathbf{P}_h)(q_h - q), w \rangle_{\mathcal{Q}} + \langle \mathbf{A}(\mathbf{P}_h - \text{Id})q_h, w \rangle_{\mathcal{Q}} + \langle \mathbf{A}q, (\mathbf{P}_h - \text{Id})w \rangle_{\mathcal{Q}} + \langle \mathbf{B}_h q, w \rangle_{\mathcal{Q}} \\
 &\leq C(C_2^{\mathbf{P}} h^{\alpha_2} \|q_h - q\|_{\mathcal{Q}} + C_3^{\mathbf{P}} h^{\alpha_3} \|q_h\|_{\mathcal{Q}} + C_1^{\mathbf{P}} h^{\alpha_1}) \|w\|_{\mathcal{Q}} + \langle \mathbf{B}_h q, w \rangle_{\mathcal{Q}}
 \end{aligned}$$



Define now

$$(G1) \quad \gamma \stackrel{\text{def}}{=} \langle D_q \mathcal{E}(t, q_h) - D_q \mathcal{E}(t, q), q_h - q \rangle_Q \geq \kappa \|q_h - q\|_Q^2$$

$$\mathcal{E}(t, \cdot) \text{ is } \kappa\text{-unif. convex} \xrightarrow{\text{Ciarlet'82}} \mathcal{E}(t, \hat{q}) \geq \mathcal{E}(t, q) + \langle D_q \mathcal{E}(t, q), \hat{q} - q \rangle_Q + \frac{\kappa}{2} \|q - \hat{q}\|_Q^2$$



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By differentiation, we get

$$\begin{aligned} \dot{\gamma} &= \langle \partial_t D_q \mathcal{E}(t, q_h) - \partial_t D_q \mathcal{E}(t, q), q_h - q \rangle_{\mathcal{Q}} \\ &\quad + \langle D_q \mathcal{E}(t, q) - D_q \mathcal{E}(t, q_h) + D_q^2 \mathcal{E}(t, q_h)[q_h - q], \dot{q}_h \rangle_{\mathcal{Q}} \\ &\quad + \langle D_q \mathcal{E}(t, q_h) - D_q \mathcal{E}(t, q) + D_q^2 \mathcal{E}(t, q)[q - q_h], \dot{q} \rangle_{\mathcal{Q}} \\ &\quad + 2 \underbrace{\langle D_q \mathcal{E}(t, q_h) - D_q \mathcal{E}(t, q), \dot{q}_h - \dot{q} \rangle_{\mathcal{Q}}} \\ &\quad \stackrel{(VI)}{\leq} \langle D_q \mathcal{E}(t, q_h), v_h - \dot{q} \rangle_{\mathcal{Q}} + \underbrace{\Psi(v_h - \dot{q})}_{(CC)} \leq C(h^\alpha + \|q_h - q\|_{\mathcal{Q}}^2) \|\dot{q}\|_{\mathcal{Q}} + \langle \mathbf{B}_h q, \dot{q} \rangle_{\mathcal{Q}} \end{aligned}$$



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By differentiation, we get

$$\dot{\gamma} \stackrel{(E1)-(E3)}{\leq} C(h^\alpha \|\dot{q}\|_Q + (\|\dot{q}\|_Q + \|\dot{q}_h\|_Q) \|q - q_h\|_Q^2) + 2 \langle \mathbf{B}_h q, \dot{q} \rangle_Q$$



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By differentiation, we get

$$\dot{\gamma} \stackrel{\text{Thm. M\&T'04}}{\leq} C \left(h^\alpha + \frac{\gamma}{\kappa} \right) + 2 \langle \mathbf{B}_h q, \dot{q} \rangle_Q$$

\Downarrow Mult. by $e^{-Ct/\kappa}$ + Int.

$$\begin{aligned} \gamma &\leq e^{Ct/\kappa} \left(\underbrace{\gamma(0)}_{\leq C \|q_h(0) - q(0)\|_Q^2} + (\kappa + 2C^{\mathbf{B}}) h^\alpha \right) \\ &\leq C \|q_h(0) - q(0)\|_Q^2 \\ &\Downarrow (G1) \end{aligned}$$

$$\|q_h(t) - q(t)\|_Q^2 \leq e^{Ct/\kappa} \left(\frac{C}{\kappa} \|q_h(0) - q(0)\|_Q^2 + (\kappa + 2C^{\mathbf{B}}) h^\alpha \right)$$



Lemma

Assume that **(E1)-(E3)**, **(D1)-(D3)** hold. Then there exists $C_0 > 0$ such that for all $h \in (0, h_0]$, we have $q_h(0) \in \mathcal{S}_h(0)$ and

$$\|q_h(0) - q(0)\|_{\mathcal{Q}} \leq C_0 h^{\alpha/2} \quad \text{with} \quad \alpha = \min\{\alpha_1, 2\alpha_2, \alpha_3\}.$$

Proof: We deduce from $\mathcal{E}(0, q(0)) + \frac{\kappa}{2} \|\hat{q} - q(0)\|_{\mathcal{Q}}^2 \leq \mathcal{E}(0, \hat{q}) + \Psi(\hat{q} - q(0))$ that

$$\begin{aligned} \frac{\kappa}{2} \|q_h(0) - q(0)\|_{\mathcal{Q}}^2 &\leq \mathcal{E}(0, \mathbf{P}_h q(0)) - \mathcal{E}(0, q(0)) + \Psi((\mathbf{P}_h - \text{Id})q(0)) \\ &\leq \langle \mathbf{D}_q \mathcal{E}(0, q(0)), (\mathbf{P}_h - \text{Id})q(0) \rangle_{\mathcal{Q}} + \mathcal{I} + \Psi((\mathbf{P}_h - \text{Id})q(0)) \end{aligned}$$

where

$$\begin{aligned} \mathcal{I} &\stackrel{\text{def}}{=} \int_0^1 \langle \mathbf{D}_q \mathcal{E}(0, q(0) + s(\mathbf{P}_h - \text{Id})q(0)) - \mathbf{D}_q \mathcal{E}(0, q(0)), (\mathbf{P}_h - \text{Id})q(0) \rangle_{\mathcal{Q}} ds \\ &\leq C^\kappa (\|(\mathbf{P}_h - \text{Id})q(0)\|_{\mathcal{Q}}^2 + \|(\mathbf{P}_h - \text{Id})q(0)\|_{\mathcal{X}}^2) \end{aligned}$$



Lemma

Assume that $q(0) \in \mathcal{S}(0)$, **(D1)** and **(D3)** hold. Then we have $D_q \mathcal{E}(0, q(0)) \in \text{Lin}(\mathcal{Q}, \mathcal{X}')$ and $\|D_q \mathcal{E}(0, q(0))\|_{\text{Lin}(\mathcal{Q}, \mathcal{X}')} \leq C^\Psi$.



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$$\begin{aligned} \frac{\kappa}{2} \|q_h(0) - q(0)\|_{\mathcal{Q}}^2 &\leq \mathcal{E}(0, \mathbf{P}_h q(0)) - \mathcal{E}(0, q(0)) + \Psi((\mathbf{P}_h - \text{Id})q(0)) \\ &\leq C^\kappa (\|(\mathbf{P}_h - \text{Id})q(0)\|_{\mathcal{Q}}^2 + \|(\mathbf{P}_h - \text{Id})q(0)\|_{\mathcal{X}}^2) + 2C^\Psi \|(\mathbf{P}_h - \text{Id})q(0)\|_{\mathcal{X}} \\ \stackrel{\text{(P1)}}{\leq} &C^\kappa \|(\mathbf{P}_h - \text{Id})q(0)\|_{\mathcal{Q}}^2 + C^\kappa (C_1^{\mathbf{P}})^2 h^{2\alpha_1} \|q(0)\|_{\mathcal{Q}}^2 + 2C^\Psi C_1^{\mathbf{P}} h^{\alpha_1} \|q(0)\|_{\mathcal{Q}} \end{aligned}$$



Lemma

Assume that $q(0) \in \mathcal{S}(0)$, **(D1)** and **(D3)** hold. Then we have $D_q \mathcal{E}(0, q(0)) \in \text{Lin}(\mathcal{Q}, \mathcal{X}')$ and $\|D_q \mathcal{E}(0, q(0))\|_{\text{Lin}(\mathcal{Q}, \mathcal{X}')} \leq C^\Psi$.

$$\begin{aligned} \frac{\kappa}{2} \|q_h(0) - q(0)\|_{\mathcal{Q}}^2 &\leq \mathcal{E}(0, \mathbf{P}_h q(0)) - \mathcal{E}(0, q(0)) + \Psi((\mathbf{P}_h - \text{Id})q(0)) \\ &\leq C^\kappa \|(\mathbf{P}_h - \text{Id})q(0)\|_{\mathcal{Q}}^2 + C^\kappa (C_1^{\mathbf{P}})^2 h^{2\alpha_1} \|q(0)\|_{\mathcal{Q}}^2 + 2C^\Psi C_1^{\mathbf{P}} h^{\alpha_1} \|q(0)\|_{\mathcal{Q}} \end{aligned}$$

It remains to evaluate $\|(\mathbf{P}_h - \text{Id})q(0)\|_{\mathcal{Q}}$

$$\begin{aligned} \frac{\kappa}{2} \|(\mathbf{P}_h - \text{Id})q(0)\|_{\mathcal{Q}}^2 &\leq C \underbrace{\langle \mathbf{A}(\mathbf{P}_h - \text{Id})q(0), (\mathbf{P}_h - \text{Id})q(0) \rangle_{\mathcal{Q}}} \\ &= \langle (\mathbf{P}_h^* \mathbf{A} - \mathbf{A} \mathbf{P}_h)(\mathbf{P}_h - \text{Id})q(0), q(0) \rangle_{\mathcal{Q}} + \langle \mathbf{A}(\mathbf{P}_h - \text{Id})\mathbf{P}_h q(0), q(0) \rangle_{\mathcal{Q}} - \langle \mathbf{A}(\mathbf{P}_h - \text{Id})q(0), q(0) \rangle_{\mathcal{Q}} \\ &+ C \|(\mathbf{P}_h - \text{Id})q(0)\|_{\mathcal{X}}^2 + 2C^\Psi \|(\mathbf{P}_h - \text{Id})q(0)\|_{\mathcal{X}} \end{aligned}$$

$$\stackrel{\text{(P1)-(P3)}}{\implies} \|(\mathbf{P}_h - \text{Id})q(0)\|_{\mathcal{Q}}^2 \leq C_0 h^\alpha \text{ with } \alpha = \min\{\alpha_1, 2\alpha_2, \alpha_3\}.$$



Theorem

Assume that **(E1)**-**(E3)**, **(D1)**-**(D3)** hold. Then there exists $C_* > 0$ such that $q(0) \in \mathcal{S}(0)$, we have

$$\|q_{\tau,h}(t) - q(t)\|_{\mathcal{Q}} \leq C_*(h^{\alpha/2} + \sqrt{\tau})$$

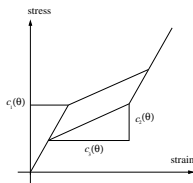
with $\alpha = \min\{\alpha_1, 2\alpha_2, \alpha_3\}$.



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- 4 Conclusion

**State variables** $u : \Omega \rightarrow \mathbb{R}^d$ displacement $z : \Omega \rightarrow \mathbb{R}_{0,\text{sym}}^{d \times d}$ mesoscopic transformation strain**Applied field** $\ell_{\text{appl}} \in C^3([0, T], \mathcal{F}^*)$ mechanical load**Energy functional:** $\mathcal{E}(t, u, z) = \int_{\Omega} W(\mathbf{e}(u), z, \nabla z) dx - \langle \ell_{\text{appl}}(t), u \rangle$ where $W(\mathbf{e}(u), z, \nabla z) = \frac{1}{2} \mathbb{C}(\mathbf{e}(u) - z) : (\mathbf{e}(u) - z) + H_{\text{SoAu}}(z) + \nu |\nabla z|^2$ ▷ $\mathbf{e}(u) = \frac{1}{2}(\nabla u + \nabla u^T)$: infinitesimal strain▷ \mathbb{C} : elasticity tensor▷ $H_{\text{SoAu}}(z) = c_1 |z| + \frac{c_2}{2} |z|^2 + \chi_{\{|z| \leq c_3\}}(z)$

- ▶ c_1 : activation threshold
- ▶ c_2 : hardening in the martensitic regime
- ▶ c_3 : maximal transformation strain

**Dissipation distance:** $\Psi(z_2 - z_1) = \int_{\Omega} \rho |z_2 - z_1| dx$ **Notations:** $\mathcal{Q} = \mathcal{F} \times \mathcal{Z}$ and $q = (u, z)$

**State variables** $u : \Omega \rightarrow \mathbb{R}^d$ displacement $z : \Omega \rightarrow \mathbb{R}_{0,\text{sym}}^{d \times d}$ mesoscopic transformation strain**Applied field** $\ell_{\text{appl}} \in C^3([0, T], \mathcal{F}^*)$ mechanical load**Energy functional:** $\mathcal{E}(t, u, z) = \int_{\Omega} W(\mathbf{e}(u), z, \nabla z) dx - \langle \ell_{\text{appl}}(t), u \rangle$ where $W(\mathbf{e}(u), z, \nabla z) = \frac{1}{2} \mathbb{C}(\mathbf{e}(u) - z) : (\mathbf{e}(u) - z) + H_{\text{SoAu}}(z) + \nu |\nabla z|^2$ Regularized version of H_{SoAu} :

$$H_{\delta}(z) = c_1 \sqrt{\delta^2 + |z|^2} + \frac{c_2}{2} |z|^2 + \frac{1}{\delta} ((|z| - c_3)_+)^3$$

Dissipation distance: $\Psi(z_2 - z_1) = \int_{\Omega} \rho |z_2 - z_1| dx$ **Theorem (Existence and uniqueness)***For all $\delta \geq 0$ there exists a solution of **(S)&(E)**.**For $\delta > 0$ the solutions are unique since $\mathcal{E} \in C^3([0, T] \times H^1(\Omega))$.*



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$$0 \in \partial\Psi(\dot{q}) + \mathbf{A}q + D_q \mathcal{H}_{\text{SoAu}}(q) - \mathbf{l}(t)$$

$$\triangleright \partial\Psi(\dot{q}) \stackrel{\text{def}}{=} (0, \partial\Psi(\dot{z}))^\top$$

$$\triangleright \mathcal{H}_{\text{SoAu}}(q) \stackrel{\text{def}}{=} (0, H_{\text{SoAu}}(z) + \frac{c_2}{2}|z|^2)^\top$$

$$\triangleright \mathbf{l}(t) \stackrel{\text{def}}{=} (\ell(t), 0)^\top$$

$$\triangleright \mathbf{A} \stackrel{\text{def}}{=} \begin{pmatrix} -\text{div}(\mathbb{C}\mathbf{e}(\cdot)) & \text{div}(\mathbb{C}(\cdot)) \\ -\mathbb{C}(\mathbf{e}(\cdot)) & \mathbb{C}(\cdot) - \nu\Delta(\cdot) + \frac{c_2}{2}\text{Id}(\cdot) \end{pmatrix}$$

Theorem

There exists $C_*^{\text{SoAu}} > 0$ such that for all $h \in (0, 1]$, we have

$$\|q_{\tau, h}(t) - q(t)\|_{\mathcal{Q}} \leq C_*^{\text{SoAu}} (\sqrt{h} + \sqrt{\tau}) \text{ for all } t \in [0, T].$$



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- ▷ include **rate-dependent** effects like a heat equation
- ▷ develop the theory to include other **multifunctional materials** (ferroelectric materials, magnetostrictive materials)



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Thank you for your attention !

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