

# On the convergence for kinetic variational inequality to quasi-static variational inequality with application to elastic-plastic systems with hardening

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In this note, a priori estimates for the kinetic problem are obtained that imply, that the kinetic solutions converge to the quasi-static ones, when the size of initial perturbations and the rate of application of the forces tend to 0. An application to three-dimensional elastic-plastic systems with hardening is given.

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## 1 Mathematical formulation

We start with a Hilbert space  $H \stackrel{\text{def}}{=} H_1 \times H_2$  with dual  $H^*$ , the dual pairing and the norm are respectively denoted by  $\langle \cdot, \cdot \rangle : H \times H^* \rightarrow \mathbb{R}$  and  $\|\cdot\|_H$ . Let  $V \stackrel{\text{def}}{=} V_1 \times V_2$  be such that  $V \subset H \subset V^*$  with the dual  $V^*$ . We use below the following norms:  $\|\cdot\|^2 \stackrel{\text{def}}{=} \langle A \cdot, \cdot \rangle$ ,  $\|\cdot\|_*^2 \stackrel{\text{def}}{=} \langle \cdot, A^{-1} \cdot \rangle$  and  $|\cdot|_M \stackrel{\text{def}}{=} \|M^{1/2} \cdot\|_H$ . We denote by  $A : H \rightarrow H^*$  a strictly positive operator. We consider the variational inclusion

$$\varepsilon^2 M \ddot{q}(t) + Aq(t) + \partial \mathcal{R}(\dot{q}(t)) \ni l(t) \text{ where } M \stackrel{\text{def}}{=} \begin{pmatrix} m & 0 \\ 0 & 0 \end{pmatrix}, A \stackrel{\text{def}}{=} \begin{pmatrix} a_{11} & a_{12} \\ a_{21} & a_{22} \end{pmatrix}, \quad (1.1)$$

$q(t) \stackrel{\text{def}}{=} (u(t), z(t))^T$  with  $u \in H_1$  and  $z \in H_2$ ,  $l(t) \stackrel{\text{def}}{=} (f_u(t), f_z(t))^T$ ,  $\varepsilon$  is a parameter which eventually tends to 0 and  $\partial \mathcal{R}$  is the subdifferential of the dissipation functional  $\mathcal{R}$  which is assumed to be convex, lower-semicontinuous, homogeneous of degree 1. Here we assume to be given initial data  $(\dot{u}(0), q(0)) \in H_1 \times V$ . The corresponding quasi-static solution  $\bar{q} \stackrel{\text{def}}{=} (\bar{u}, \bar{z})$  solves (1.1) with  $\varepsilon = 0$  together with  $\bar{q}(0) \in V$ . In what concerns the kinetic problem, the theory of maximal monotone operators is used to prove existence of a unique solution whereas for the quasi-static problem, we refer to [3] or the recently developed theory in [6].

## 2 A priori estimates

The aim of this section is to estimate the distance between  $q$  and  $\bar{q}$  by introducing a special intermediate solution  $\hat{q} = (\hat{u}, \hat{z})$  satisfies (1.1) together with initial conditions  $(\hat{u}(0), \hat{q}(0)) = (0, \bar{q}(0))$ . Notice that the existence of a unique  $\hat{q}$  is obtained using the theory of maximal monotone operators. First, we provide a priori estimates for the problem which allow us to control the term  $M \ddot{\hat{q}}$  in  $H$  instead of the usual estimates in  $V^*$ . The problem occurs through the fact that  $\partial \mathcal{R}$  is nonsmooth and classical techniques for smooth problems do not suffice. One way to handle this is to use Yosida regularization (see [4]). Here we choose a different technique that is based on difference quotients. Our a priori estimates can be derived most easily by using the bilinear form  $B : V \times H \times V^* \rightarrow \mathbb{R}$  defined via  $B[q, \dot{q}, l] \stackrel{\text{def}}{=} |\varepsilon \dot{q}|_M^2 + \|q - A^{-1}l\|^2 + \|l\|_*^2$ . The construction is such that for solutions we have  $B[q, \dot{q}, l] = 2\mathcal{E}(t, q, \dot{q}) + 2\|l\|_*^2$  with  $\mathcal{E}(t, q, \dot{q}) = \frac{1}{2} \langle \varepsilon^2 M \dot{q}, \dot{q} \rangle + \frac{1}{2} \langle Aq, q \rangle - \langle l(t), q \rangle$ . Moreover,  $B$  defines an equivalent norm on  $V \times H \times V^*$ , since  $\frac{1}{g^2} (\|q\|^2 + |\varepsilon \dot{q}|_M^2 + \|l\|_*^2) \leq B[q, \dot{q}, l] \leq g^2 (\|q\|^2 + |\varepsilon \dot{q}|_M^2 + \|l\|_*^2)$  with  $g \stackrel{\text{def}}{=} \frac{1+\sqrt{5}}{2} \approx 1.618$  is the golden ratio. Then, one can deduce the following proposition

**Proposition 2.1** *Let  $l_1, l_2 \in W^{1,1}([0, T]; V^*)$  and  $q_1$  and  $q_2$  be solutions of (1.1) with right-hand sides  $l_1$  and  $l_2$  respectively, then  $w = q_1 - q_2$  satisfies the estimate*

$$\forall t \in [0, T] : B[w(t), \dot{w}(t), l_1(t) - l_2(t)]^{1/2} \leq B[w(0), \dot{w}(0), l_1(0) - l_2(0)]^{1/2} + g\sqrt{2} \int_0^t \|\dot{l}_1(\tau) - \dot{l}_2(\tau)\|_* \, d\tau. \quad (2.1)$$

Since  $\hat{u}(0) = 0$ , we may choose  $\hat{q}(t) = \bar{q}(0)$  for all  $t \in [-\rho, 0]$ ,  $\rho > 0$ , and let  $l(t) = A\hat{q}(0)$ . Hence, we deduce that the stability condition  $l(0) \in \partial \mathcal{R}(0) + A\hat{q}(0)$  holds and that the following limits for  $h \searrow 0$  exist:

$$\frac{1}{h} (\hat{q}(0) - \hat{q}(-h)) \rightarrow 0 \text{ in } V, \frac{1}{h} (\hat{u}(0) - \hat{u}(-h)) \rightarrow 0 \text{ in } H_1, \frac{1}{h} (l(0) - l(-h)) \rightarrow 0 \text{ in } V^*. \quad (2.2)$$

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Using Proposition 2.1 and (2.2), we find

$$|\varepsilon \ddot{q}(t)|_M \leq gB[\hat{q}(t), \ddot{q}(t), \dot{l}(t)]^{1/2} \leq gB[0, 0, 0]^{1/2} + g^2\sqrt{2}\left(\|\dot{l}(0)\|_* + \int_0^t \|\ddot{l}(s)\|_* ds\right). \quad (2.3)$$

On the other hand, we use the standard trick of adding variational inequalities leads to

$$|\varepsilon \dot{q}(t)|_M^2 + \|\hat{q}(t) - \bar{q}(t)\|^2 \leq 2\bar{C} \operatorname{ess\,sup}_{s \in [0, T]} |\varepsilon \ddot{q}(s)|_M \int_0^t \varepsilon \|l(s)\|_* ds. \quad (2.4)$$

The difference between the solution  $q$  and the special solution  $\hat{q}$  comes from Proposition 2.1. Hence, (2.3) and (2.4) imply the following theorem:

**Theorem 2.2** *Let the above assumptions on  $M$ ,  $A$  and  $\mathcal{R}$  hold and assume  $l = (f_u, f_z) \in W^{2,1}([0, T], V^*)$ . For  $\bar{q}(0) \in V$  with  $l(0) \in \partial\mathcal{R}(0) + A\bar{q}(0) \subset \{0\} \times H_2$  let  $\bar{q}$  be the unique solution of (1.1) with  $\varepsilon = 0$ . For arbitrary  $q(0) = (u(0), z(0)) \in V$  and  $\dot{u}(0) \in H_1$  satisfying  $f_z(0) \in a_{21}u(0) + a_{22}z(0) + \partial\mathcal{R}(0)$  let  $q$  be the unique solution of (1.1). Then there exists  $C > 0$  such that*

$$(|\varepsilon \dot{q}(t)|_M^2 + \|q(t) - \bar{q}(t)\|^2)^{1/2} \leq (|\varepsilon \dot{u}(0)|_M^2 + \|q(0) - \bar{q}(0)\|^2)^{1/2} + C\sqrt{\varepsilon}. \quad (2.5)$$

### 3 Application: elastic-plastic systems with hardening

We consider a body with reference configuration  $\Omega \subset \mathbb{R}^d$ . This body may undergo displacements  $u : \Omega \rightarrow \mathbb{R}^d$ . The plastic strain will be characterized by  $e^{\text{pl}} : \Omega \rightarrow \mathbb{S}_0^d$  where  $\mathbb{S}_0^d$  is the space of symmetric  $d \times d$  tensors such that the trace of  $e^{\text{pl}}$  vanishes. The set of admissible displacements  $\mathcal{F}$  is chosen as a suitable subspace of  $W^{1,2}(\Omega; \mathbb{R}^d)$  by describing Dirichlet data at the part  $\Gamma_{\text{Dir}}$  of  $\partial\Omega$  and the plastic variable  $e^{\text{pl}}$  lives in  $\mathcal{Z} \stackrel{\text{def}}{=} L^2(\Omega; \mathbb{S}_0^d)$ . We denote by  $e(u) \stackrel{\text{def}}{=} \frac{1}{2}(\nabla u + \nabla u^T)$  and  $\mathbb{E}$  respectively the linearized strain tensor and elasticity tensor. We define the dissipation potential by  $\mathcal{R}(e^{\text{pl}}) \stackrel{\text{def}}{=} \int_{\Omega} R(x, e^{\text{pl}}(x)) dx$  where  $R(x, \cdot)$  is 1-homogeneous, convex and satisfies  $0 < r_1|v| \leq R(x, v) \leq r_2|v|$ . We consider the governing system

$$\begin{cases} \varepsilon^2 \rho \ddot{u} - \operatorname{div}(\mathbb{E}:(e(u) - e^{\text{pl}})) = l_{\text{ext}}(t), \\ -\dot{e}^{\text{pl}} + \partial R^*(\mathbb{E}:(e(u) - e^{\text{pl}}) - \mathbb{H}e^{\text{pl}}) \ni 0. \end{cases} \quad (3.1)$$

We assume to be given initial data  $(u(0), \dot{u}(0), e^{\text{pl}}(0)) \in \mathcal{F} \times \mathcal{F} \times \mathcal{Z}$ . Here  $\rho > 0$  is the density,  $l_{\text{ext}}(t)$  is the applied mechanical loading and  $R^*$  is the Legendre transform of  $R$ . Let  $(\bar{u}, \bar{e}^{\text{pl}})$  solves (3.1) with  $\varepsilon = 0$  and  $(\bar{u}(0), \bar{e}^{\text{pl}}(0)) \in \mathcal{F} \times \mathcal{Z}$ . Notice that the existence result for the kinetic problem was established in [7, 8] whereas existence and uniqueness of a strong solution to the corresponding quasi-static problem was obtained in [1, 2]. On the other hand, applying Theorem 2.2 and using Korn's inequality, we deduce the following Corollary:

**Corollary 3.1** *Assume that  $(u(0), e^{\text{pl}}(0))$  and  $(\bar{u}(0), \bar{e}^{\text{pl}}(0))$  satisfy  $0 \in \mathbb{E}:(e^{\text{pl}}(0) - e(u(0))) + \mathbb{H}e^{\text{pl}}(0) + \partial\mathcal{R}(0)$  and  $(l_{\text{ext}}(0), 0)^T \in A(\bar{u}(0), \bar{e}^{\text{pl}}(0))^T + \{0\} \times \partial R(0)$ , respectively. Then there exist  $c, C > 0$  such that*

$$\begin{aligned} \forall \varepsilon > 0 : & \left( \|\varepsilon \rho^{1/2} \dot{u}(t)\|_{L^2}^2 + \|u(t) - \bar{u}(t)\|_{W^{1,2}}^2 + \|e^{\text{pl}}(t) - \bar{e}^{\text{pl}}(t)\|_{L^2}^2 \right)^{1/2} \\ & \leq c \left( \|\varepsilon \rho^{1/2} \dot{u}(0)\|_{L^2}^2 + \|u(0) - \bar{u}(0)\|_{W^{1,2}}^2 + \|e^{\text{pl}}(0) - \bar{e}^{\text{pl}}(0)\|_{L^2}^2 \right)^{1/2} + C\sqrt{\varepsilon}. \end{aligned}$$

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