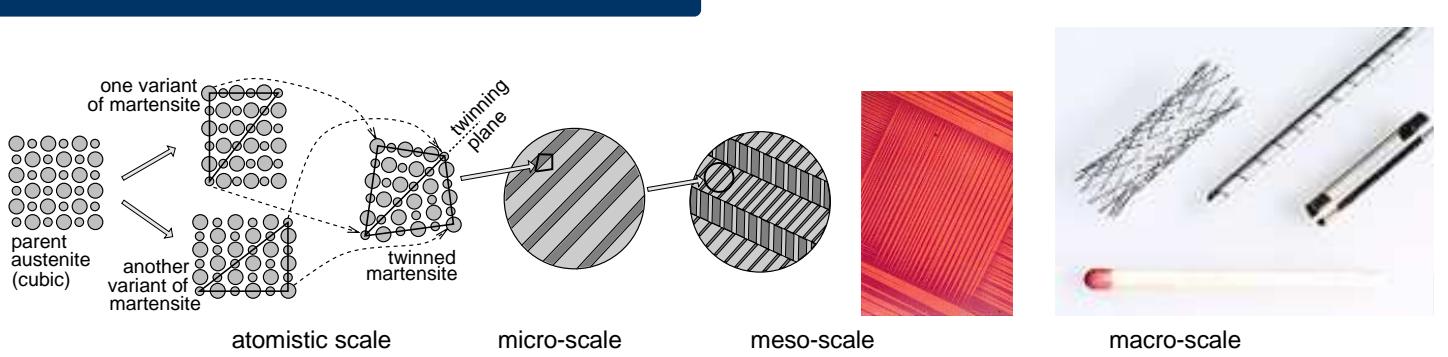




# Analysis and numerics of multidimensional models for elastic phase transformations in shape-memory alloys

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## Multiscale modelling in shape-memory materials



## Mechanical modelling

needs

- energetics and elastic properties of  $N$  phases **austenite** and all the variants of **martensite** (e.g.: tetragonal 1+3; orthorhombic 1+6; 1+24)
- (temperature dependent) criteria for stress and strain induced transformations



## Mathematical model

- **stored energy potential  $E$ :**

$$E(t, \nabla \varphi, z) = \int_{\Omega} W(\varphi, z, \theta) + \frac{\kappa}{r} |\nabla z|^r dx - \langle \ell(t), \varphi \rangle$$

$\varphi$  deformation,  $z \in \mathbb{R}^N$  phase fractions

- **dissipation potential  $R$**

$$R(z, \dot{z}) = \int_{\Omega} R(z(x), \dot{z}(x)) dx$$

up to now:  $\theta = \theta_{\text{applied}}(t, x)$  is prescribed as data

## Energetic formulation of rate-independent evolution via (S) and (E) (see [Mie07])

**(S) static stability**  $E(t, \varphi(t), z(t)) \leq E(t, \tilde{\varphi}, \tilde{z}) + D(z(t), \tilde{z})$  for all  $\tilde{\varphi}, \tilde{z}$

**(E) energy balance**  $E(t, \varphi(t), z(t)) + \int_0^t R(z(s), \dot{z}(s)) ds = E(0, \varphi(0), z(0)) + \int_0^t \partial_s E(s, \varphi(s), z(s)) ds$

## Advantages

- derivative-free formulation ( $\leftarrow$  jumps, nonsmoothness)
- usage of microscopic constitutive laws for each phase  $W(\cdot, e_j)$  for  $j \in \{0, 1, \dots, N\}$  possible
- methods from calculus of variations are available
- time-incremental problem via minimization

$$(IP) \quad (\varphi_k, z_k) \in \text{Arg min} E(\varphi_k, z_k) + D(z_k, \tilde{z})$$

## Challenges

- establish existence and uniqueness for (IP) and for (S)&(E)
- study the influence of interfacial energy and the  $\Gamma$ -limit for vanishing interfacial energy
- numerical approximation for (IP) and convergence to (S)&(E) with optimal rates
- include proper heat equation

## Souza-Auricchio model

- **State variables and applied fields**  
 $\varphi : \Omega \rightarrow \mathbb{R}^d$  displacement ( $e(\varphi) = \frac{1}{2}(\nabla \varphi + \nabla \varphi^T)$ )  
 $z : \Omega \rightarrow \mathbb{R}_{0,\text{sym}}^{d \times d}$  mesoscopic transformation strain  
 $l \in C^1([0, T]; H^{-1}(\Omega))$  exterior loading  
 $\theta = \theta_{\text{appl}} \in C^1([0, T]; L^\infty(\Omega))$  temperature
- **Energy and dissipation potentials**  
 $W(\nabla \varphi, z, \theta) = \frac{1}{2}(e(\varphi) - z) : \mathbb{C}(\theta) : (e(\varphi) - z) + h_\delta(z, \theta)$   
 $R(z, \dot{z}) = \rho |\dot{z}|_1$

Hardening function with  $\delta \in [0, 1]$

$$h_\delta(z, \theta) = c_1(\theta) \sqrt{\delta^2 + |z|^2} + \frac{c_2(\theta)}{2} |z|^2 + \frac{1}{8} (\max\{0, |z| - c_3(\theta)\})^3$$

**Theorem** (see [MP07])

For all  $\delta \geq 0$ , (S)&(E) has a solution  
 $(\varphi, z) \in C^{\text{Lip}}([0, T]; H^1(\Omega; \mathbb{R}^d) \times H^1(\Omega; \mathbb{R}_{0,\text{sym}}^{d \times d}))$ .

For  $\delta > 0$  we have  $E \in C^{2,\text{Lip}}([0, T]; H^1(\Omega))$   
and the solutions are unique.

