

**Weierstraß-Institut
für Angewandte Analysis und Stochastik
Leibniz-Institut im Forschungsverbund Berlin e. V.**

Preprint

ISSN 2198-5855

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the Oseen–Frank free energy**

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submitted: January 24, 2018

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No. 2474
Berlin 2018



2010 Mathematics Subject Classification. 35Q35, 35K52, 35R06, 76A15.

Key words and phrases. Liquid crystal, Ericksen–Leslie equation, measure-valued solution, weak-strong uniqueness.

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Weak-strong uniqueness for measure-valued solutions to the Ericksen–Leslie model equipped with the Oseen–Frank free energy

Robert Lasarzik

Abstract

We analyze the Ericksen–Leslie system equipped with the Oseen–Frank energy in three space dimensions. Recently, the author introduced the concept of measure-valued solutions to this system and showed the global existence of these generalized solutions. In this paper, we show that suitable measure-valued solutions, which fulfill an associated energy inequality, enjoy the weak-strong uniqueness property, i. e. the measure-valued solution agrees with a strong solution if the latter exists. The weak-strong uniqueness is shown by a relative energy inequality for the associated nonconvex energy functional.

1 Introduction

Nonlinear partial differential equations require generalized solution concepts. Uniqueness of solutions, however often is an open problem because of lack of regularity. A prominent example are the Navier–Stokes equations in three space dimensions with Leray's [30] weak solutions and Serrin's [39] uniqueness result. Already there the concept of weak-strong uniqueness is applied: a generalized solution is compared with a solution exhibiting more regularity. The generalized solution concept of measure-valued solutions was first introduced by Tartar [40] using conventional Young measures. This concept is also well-known in the context of the Euler equations. DiPerna and Majda [11] define measure-valued solutions to the Euler equation by introducing so-called *generalized Young-measures*, capturing oscillation and concentration effects for sequences that are only bounded in L^1 . Another step in the analysis of such sequences is done in Alibert and Bouchitté [2] and we will heavily rely on the techniques introduced there. Brenier, De Lellis, and Székelyhidi [5] prove the first weak-strong uniqueness result for measure-valued solutions, involving generalized Young measures. They consider the incompressible Euler equation and certain hyperbolic conservation laws. Some of the techniques introduced in [5] are also used in this work. Additionally, there are works on the weak-strong uniqueness of measure-valued solutions by Demoulini, Stuart and Tzavaras [9] in the context of elastodynamics and by Gwiazda, Świerczewska-Gwiazda & Wiedemann [24] for a class of compressible Euler equations.

An important argument of this article relies on the formulation of a relative energy, which is often called relative entropy in this context (see Feireisl, Jin & Novotný [19]). The idea of a relative entropy to compare two solutions goes back to Dafermos [7].

For a strictly convex entropy function $\eta : \mathbb{R} \rightarrow \mathbb{R}$, the relative entropy of two solutions u and \tilde{u} is given by (see Ref. [8, Sec. 5.3])

$$\mathcal{E} := \eta(u) - \eta(\tilde{u}) - \langle \eta'(\tilde{u}), (u - \tilde{u}) \rangle. \quad (1.1)$$

The strict convexity of η guarantees that \mathcal{E} is nonnegative.

The approach of the relative energy for convex functionals has been used, e.g., to show the weak-strong uniqueness property of solutions (see Feireisl and Novotný [20]), the stability of an equilibrium state (see Feireisl [17]), the convergence to a singular limit problem (see Breit, Feireisl and Hofmanova [4] as well as Feireisl [18]), or to derive *a posteriori* estimates of numerical solutions (see Fischer [22]). Another possible application is the definition of a generalized solution concept, the so-called dissipative solutions. The formulation of such a concept relies on an inequality instead of an equality (see Lions [36, Sec. 4.4]).

Our aim is to generalize the above concept of weak-strong uniqueness for measure-valued solutions via a relative energy approach to the three dimensional Ericksen–Leslie model equipped with the Oseen–Frank energy describing nematic liquid crystal flow.

Nematic liquid crystals are anisotropic fluids. The rod-like molecules build, or are dispersed in, a fluid and are directionally ordered. This ordering and its direction heavily influences the properties of the material such as light scattering or flow

behaviour. This gives rise to many applications, where *liquid crystal displays* are only the most prominent ones. The *Ericksen–Leslie model* is the most common model to describe nematic liquid crystals. The direction of the aligned molecule is modelled by a unit-vector field and the fluid flow by a velocity field. Since this model was proposed by Ericksen [16] and Leslie [31] in the 60ies, it is extensively studied.

The first mathematical analysis of a simplified Ericksen–Leslie model is done by Lin and Liu [32]. They show global existence of weak solutions and local existence of strong solutions. Additionally, they manage to generalise this results to a more realistic model [34]. They also show partial regularity of weak solutions to the considered system [33]. Following this works, there have been many articles considering slightly more complicated models, see [3], [6], [21], or [41] for example. To the best of the author's knowledge, the only generalisation with respect to the free energy potential is performed by Emmrich and the author in [15].

There are also results on *local existence of solutions* to realistic models (see for instance [26], [42] or [25]). Especially, local strong solutions are known to exists to different simplifications of the system considered in this article. The full (thermodynamically consistent) Ericksen–Leslie system with the one constant approximation of the Oseen–Frank energy is considered in [25]. Whereas, the simplified Ericksen–Leslie system with the full Oseen–Frank energy is studied in [26]. For more on liquid crystals, we refer to Emmrich, Klapp and Lasarzik [14]. Since finite time singularities in nematic liquid crystals have been observed experimentally [1] and analytically [27], it seems appropriate to inquire a weakened solution concept such as measure-valued solutions.

Recently, the existence of measure-valued solutions to the Ericksen–Leslie model equipped with the Oseen–Frank energy has been proven by the author (see [29]). It is the first solution concept for the Ericksen–Leslie model in its full generality. The measure-valued solutions are a very weak solution concept. The Gradient of the director is only represented by a generalized Young measure. This article is dedicated to the proof of the weak-strong uniqueness property of the measure-valued solutions to the Ericksen–Leslie model. The main theorem (see Theorem 2.1) shows that a suitable measure-valued solution coincides with a strong solution emanating from the same initial data, as long as the latter exists. A suitable measure-valued solution fulfills an additional energy inequality. This energy inequality can only be shown to hold under the assumption of Parodi's relation (see (2.2e) below), which follows from Onsager's reciprocal principle. The weak-strong uniqueness can also shown to hold for suitable measure-valued solutions without assuming Parodi's relation, but these solutions are not known to exist.

The novelty of this paper is the generalization of the relative energy approach to a system with a nonconvex energy. Since the Oseen–Frank energy is not convex, the relative energy defined by (1.1) is not necessarily positive anymore. In this paper we use an alternative way to define the relative energy (see (4.1)) for the nonconvex energy Oseen–Frank energy and derive a relative energy inequality resulting in the weak-strong uniqueness property of the solutions. This new approach can hopefully be used to prove other properties, like *a posteriori* estimates or the existence of dissipative solutions, of the system too.

The paper is organised as follows: In Section 1, we collect some notation. Section 2 contains the model, the definition of the suitable generalized solutions, the definition of the strong solutions and the main result. In Section 3, some auxiliary lemmas are collected and the energy equality for the strong solution is shown. While Section 4 introduces the relative energy and some associated estimates, Section 5 collects certain integration-by-parts formulae. The proof of the main result is carried out in Section 6. In Appendix A, some tensor calculations are collected.

Notation

Vectors of \mathbb{R}^3 are denoted by bold small Latin letters. Matrices of $\mathbb{R}^{3 \times 3}$ are denoted by bold capital Latin letters. We also use tensors of higher order, which are denoted by bold capital Greek letters. Moreover, numbers are denoted be small Latin or Greek letters, and capital Latin letters are reserved for potentials. The euclidean scalar product in \mathbb{R}^3 is denoted by a dot $\mathbf{a} \cdot \mathbf{b} := \mathbf{a}^T \mathbf{b} = \sum_{i=1}^3 a_i b_i$, for $\mathbf{a}, \mathbf{b} \in \mathbb{R}^3$ and the Frobenius product in $\mathbb{R}^{3 \times 3}$ by a double point $\mathbf{A} : \mathbf{B} := \text{tr}(\mathbf{A}^T \mathbf{B}) = \sum_{i,j=1}^3 A_{ij} B_{ij}$, for $\mathbf{A}, \mathbf{B} \in \mathbb{R}^{3 \times 3}$. Additionally, the scalar product in the space of Tensors of order three is denoted by three dots

$$\boldsymbol{\Upsilon} : \boldsymbol{\Gamma} := \left[\sum_{j,k,l=1}^3 \boldsymbol{\Upsilon}_{jkl} \boldsymbol{\Gamma}_{jkl} \right], \quad \boldsymbol{\Upsilon} \in \mathbb{R}^{3 \times 3 \times 3}, \boldsymbol{\Gamma} \in \mathbb{R}^{3 \times 3 \times 3}.$$

The associated norms are all denoted by $|\cdot|$, where also the norms of tensors of higher order are denoted in the same way

$$|\boldsymbol{\Lambda}|^2 := \sum_{i,j,k,l=1}^3 \boldsymbol{\Lambda}_{ijkl}^2, \quad \text{for } \boldsymbol{\Lambda} \in \mathbb{R}^{3^4} \quad \text{and} \quad |\boldsymbol{\Theta}|^2 := \sum_{i,j,k,l,m,n=1}^3 \boldsymbol{\Theta}_{ijklmn}^2, \quad \text{for } \boldsymbol{\Theta} \in \mathbb{R}^{3^6}$$

respectively. Similar, we define the products of tensors of different order. The product of a tensor of third order with a matrix is defined by

$$\boldsymbol{\Gamma} : \mathbf{A} := \left[\sum_{j,k=1}^3 \boldsymbol{\Gamma}_{ijk} \mathbf{A}_{jk} \right]_{i=1}^3, \quad \boldsymbol{\Gamma} \cdot \mathbf{A} := \left[\sum_{k=1}^3 \boldsymbol{\Gamma}_{ijk} \mathbf{A}_{kl} \right]_{i,j,l=1}^3, \quad \boldsymbol{\Gamma} \in \mathbb{R}^{3 \times 3 \times 3}, \mathbf{A} \in \mathbb{R}^{3 \times 3}$$

and the product of a tensor of third order with a vector by

$$\boldsymbol{\Gamma} \cdot \mathbf{a} := \left[\sum_{k=1}^3 \boldsymbol{\Gamma}_{ijk} \mathbf{a}_k \right]_{i,j=1}^3, \quad \boldsymbol{\Gamma} \in \mathbb{R}^{3 \times 3 \times 3}, \mathbf{a} \in \mathbb{R}^3.$$

The product of a tensor of fourth order with a matrix and a vector is defined by

$$\boldsymbol{\Lambda} : \mathbf{A} := \left[\sum_{k,l=1}^3 \boldsymbol{\Lambda}_{ijkl} \mathbf{A}_{kl} \right]_{i,j=1}^3, \quad \boldsymbol{\Lambda} : \mathbf{a} := \left[\sum_{l=1}^3 \boldsymbol{\Lambda}_{ijkl} \mathbf{a}_l \right]_{i,j,k=1}^3, \quad \boldsymbol{\Lambda} \in \mathbb{R}^{3^4}, \mathbf{A} \in \mathbb{R}^{3 \times 3}, \mathbf{a} \in \mathbb{R}^3.$$

The product of tensors of fourth and third order is given by

$$\boldsymbol{\Lambda} : \boldsymbol{\Gamma} := \left[\sum_{k,l=1}^3 \boldsymbol{\Lambda}_{ijkl} \boldsymbol{\Gamma}_{klm} \right]_{i,j,m=1}^3, \quad \boldsymbol{\Lambda} : \boldsymbol{\Gamma} := \left[\sum_{j,k,l=1}^3 \boldsymbol{\Lambda}_{ijkl} \boldsymbol{\Gamma}_{jkl} \right]_{i=1}^3, \quad \boldsymbol{\Lambda} \in \mathbb{R}^{3^4}, \boldsymbol{\Gamma} \in \mathbb{R}^{3 \times 3 \times 3}.$$

The product of a tensor of fourth order and a matrix or a tensor of third order is defined via

$$\mathbf{A} : \boldsymbol{\Theta} := \left[\sum_{i,j=1}^3 \mathbf{A}_{ij} \boldsymbol{\Theta}_{ijklmn} \right]_{k,l,m,n=1}^3, \quad \boldsymbol{\Theta} : \boldsymbol{\Gamma} := \left[\sum_{l,m,n=1}^3 \boldsymbol{\Theta}_{ijklmn} \boldsymbol{\Gamma}_{lmn} \right]_{i,j,k=1}^3, \quad \boldsymbol{\Theta} \in \mathbb{R}^{3^6}, \mathbf{A} \in \mathbb{R}^{3 \times 3}, \boldsymbol{\Gamma} \in \mathbb{R}^{3 \times 3 \times 3}.$$

The product of a vector and a tensor of fourth order is defined differently. The definition is adjusted to the cases of this work:

$$\mathbf{a} \cdot \boldsymbol{\Theta} := \left[\sum_{k=1}^3 \mathbf{a}_k \boldsymbol{\Theta}_{ijklmn} \right]_{i,j,l,m,n=1}^3, \quad \boldsymbol{\Theta} \in \mathbb{R}^{3^6}, \mathbf{a} \in \mathbb{R}^3.$$

The standard matrix and matrix-vector multiplication is written without an extra sign for brevity,

$$\mathbf{AB} = \left[\sum_{j=1}^3 \mathbf{A}_{ij} \mathbf{B}_{jk} \right]_{i,k=1}^3, \quad \mathbf{A}\mathbf{a} = \left[\sum_{j=1}^3 \mathbf{A}_{ij} \mathbf{a}_j \right]_{i=1}^3, \quad \mathbf{A} \in \mathbb{R}^{3 \times 3}, \mathbf{B} \in \mathbb{R}^{3 \times 3}, \mathbf{a} \in \mathbb{R}^3.$$

The outer vector product is given by $\mathbf{a} \otimes \mathbf{b} := \mathbf{ab}^T = [\mathbf{a}_i \mathbf{b}_j]_{i,j=1}^3$ for two vectors $\mathbf{a}, \mathbf{b} \in \mathbb{R}^3$ and by $\mathbf{A} \otimes \mathbf{a} := \mathbf{Aa}^T = [\mathbf{A}_{ij} \mathbf{a}_k]_{i,j,k=1}^3$ for a matrix $\mathbf{A} \in \mathbb{R}^{3 \times 3}$ and a vector $\mathbf{a} \in \mathbb{R}^3$. The symmetric and skew-symmetric parts of a matrix are given by $\mathbf{A}_{\text{sym}} := \frac{1}{2}(\mathbf{A} + \mathbf{A}^T)$ and $\mathbf{A}_{\text{skw}} := \frac{1}{2}(\mathbf{A} - \mathbf{A}^T)$, respectively ($\mathbf{A} \in \mathbb{R}^{3 \times 3}$). For the product of two matrices $\mathbf{A}, \mathbf{B} \in \mathbb{R}^{3 \times 3}$, we observe

$$\mathbf{A} : \mathbf{B} = \mathbf{A} : \mathbf{B}_{\text{sym}}, \quad \text{if } \mathbf{A}^T = \mathbf{A} \quad \text{and} \quad \mathbf{A} : \mathbf{B} = \mathbf{A} : \mathbf{B}_{\text{skw}}, \quad \text{if } \mathbf{A}^T = -\mathbf{A}.$$

Furthermore, it holds $\mathbf{A}^T \mathbf{B} : \mathbf{C} = \mathbf{B} : \mathbf{AC}$ for $\mathbf{A}, \mathbf{B}, \mathbf{C} \in \mathbb{R}^{3 \times 3}$ and $\mathbf{a} \otimes \mathbf{b} : \mathbf{A} = \mathbf{a} \cdot \mathbf{Ab}$ for $\mathbf{a}, \mathbf{b} \in \mathbb{R}^3, \mathbf{A} \in \mathbb{R}^{3 \times 3}$ and hence $\mathbf{a} \otimes \mathbf{a} : \mathbf{A} = \mathbf{a} \cdot \mathbf{Aa} = \mathbf{a} \cdot \mathbf{A}_{\text{sym}} \mathbf{a}$.

We use the Nabla symbol ∇ for real-valued functions $f : \mathbb{R}^3 \rightarrow \mathbb{R}$, vector-valued functions $\mathbf{f} : \mathbb{R}^3 \rightarrow \mathbb{R}^3$ as well as matrix-valued functions $\mathbf{A} : \mathbb{R}^3 \rightarrow \mathbb{R}^{3 \times 3}$ denoting

$$\nabla f := \left[\frac{\partial f}{\partial \mathbf{x}_i} \right]_{i=1}^3, \quad \nabla \mathbf{f} := \left[\frac{\partial \mathbf{f}_i}{\partial \mathbf{x}_j} \right]_{i,j=1}^3, \quad \nabla \mathbf{A} := \left[\frac{\partial \mathbf{A}_{ij}}{\partial \mathbf{x}_k} \right]_{i,j,k=1}^3.$$

The divergence of a vector-valued and a matrix-valued function is defined by

$$\nabla \cdot \mathbf{f} := \sum_{i=1}^3 \frac{\partial \mathbf{f}_i}{\partial \mathbf{x}_i} = \text{tr}(\nabla \mathbf{f}), \quad \nabla \cdot \mathbf{A} := \left[\sum_{j=1}^3 \frac{\partial \mathbf{A}_{ij}}{\partial \mathbf{x}_j} \right]_{i=1}^3.$$

Additionally, we abbreviate $\nabla \nabla$ and $\nabla \cdot \nabla \cdot$ by ∇^2 and $\nabla^2 :$ respectively. For a given tensor of fourth order, we abbreviate the associated second order operator by $\Delta_{\Lambda} \mathbf{d} := \nabla \cdot \mathbf{A} : \nabla \mathbf{d}$ acting on functions $\mathbf{d} \in \mathcal{C}^1(\Omega \times [0, T]; \mathbb{R}^3)$.

Throughout this paper, let $\Omega \subset \mathbb{R}^3$ be a bounded domain of class $\mathcal{C}^{3,1}$. We rely on the usual notation for spaces of continuous functions, Lebesgue and Sobolev spaces. Spaces of vector-valued functions are emphasised by bold letters, for example $\mathbf{L}^p(\Omega) := L^p(\Omega; \mathbb{R}^3)$, $\mathbf{W}^{k,p}(\Omega) := W^{k,p}(\Omega; \mathbb{R}^3)$. The standard inner product in $L^2(\Omega; \mathbb{R}^3)$ is just denoted by (\cdot, \cdot) , in $L^2(\Omega; \mathbb{R}^{3 \times 3})$ by $(\cdot; \cdot)$, and in $L^2(\Omega; \mathbb{R}^{3 \times 3 \times 3})$ by $(\cdot ; \cdot)$.

The space of smooth solenoidal functions with compact support is denoted by $\mathcal{C}_{c,\sigma}^\infty(\Omega; \mathbb{R}^3)$. By $\mathbf{L}_\sigma^p(\Omega)$, $\mathbf{H}_{0,\sigma}^1(\Omega)$, and $\mathbf{W}_{0,\sigma}^{1,p}(\Omega)$, we denote the closure of $\mathcal{C}_{c,\sigma}^\infty(\Omega; \mathbb{R}^3)$ with respect to the norm of $\mathbf{L}^p(\Omega)$, $\mathbf{H}^1(\Omega)$, and $\mathbf{W}^{1,p}(\Omega)$ respectively.

The dual space of a Banach space V is always denoted by V^* and equipped with the standard norm; the duality pairing is denoted by $\langle \cdot, \cdot \rangle$. The duality pairing between $\mathbf{L}^p(\Omega)$ and $\mathbf{L}^q(\Omega)$ (with $1/p + 1/q = 1$), however, is denoted by (\cdot, \cdot) , $(\cdot; \cdot)$, or $(\cdot ; \cdot)$.

The unit ball in d dimensions is denoted by $B_d := \{\mathbf{x} \in \mathbb{R}^d; |\mathbf{x}| < 1\}$ and the sphere in d -dimensions by $\mathcal{S}^{d-1} := \{\mathbf{x} \in \mathbb{R}^d; |\mathbf{x}| = 1\}$.

For $Q \subset \mathbb{R}^d$, the Radon measures are denoted by $\mathcal{M}(Q)$, the positive Radon measures by $\mathcal{M}^+(Q)$, and probability measures by $\mathcal{P}(Q)$. We recall that the Radon measures equipped with the total variation are a Banach space and for compact sets Q , it can be characterized by $\mathcal{M}(Q) = (\mathcal{C}(Q))^*$ (see [12, Theorem 4.10.1]). The integration of a function $f \in \mathcal{C}(Q)$ with respect to a measure $\mu \in \mathcal{M}(Q)$ is denoted by $\int_Q f(\mathbf{h}) \mu(d\mathbf{h})$. In case of the Lebesgue measure we just write $\int_Q f(\mathbf{h}) d\mathbf{h}$.

The cross product of two vectors is denoted by \times . We introduce the notation $[\cdot]_{\mathbf{X}}$, which is defined via

$$[\cdot]_{\mathbf{X}} : \mathbb{R}^d \rightarrow \mathbb{R}^{d \times d}, \quad [\mathbf{h}]_{\mathbf{X}} := \begin{pmatrix} 0 & -\mathbf{h}_3 & \mathbf{h}_2 \\ \mathbf{h}_3 & 0 & -\mathbf{h}_1 \\ -\mathbf{h}_2 & \mathbf{h}_1 & 0 \end{pmatrix}. \quad (1.2)$$

The i -th component of the vector $\mathbf{h} \in \mathbb{R}^3$ is denoted by \mathbf{h}_i . The mapping $[\cdot]_{\mathbf{X}}$ has some nice properties, for instance

$$[\mathbf{a}]_{\mathbf{X}} \mathbf{b} = \mathbf{a} \times \mathbf{b}, \quad [\mathbf{a}]_{\mathbf{X}}^T [\mathbf{b}]_{\mathbf{X}} = (\mathbf{a} \cdot \mathbf{b}) I - \mathbf{b} \otimes \mathbf{a},$$

for all $\mathbf{a}, \mathbf{b} \in \mathbb{R}^3$, where I denotes the identity matrix in $\mathbb{R}^{3 \times 3}$ or

$$[\mathbf{a}]_{\mathbf{X}} : \nabla \mathbf{b} = [\mathbf{a}]_{\mathbf{X}} : (\nabla \mathbf{b})_{\text{skw}} = \mathbf{a} \cdot \nabla \times \mathbf{b}, \quad \nabla \cdot [\mathbf{a}]_{\mathbf{X}} = -\nabla \times \mathbf{a}, \quad \frac{1}{2} [\nabla \times \mathbf{a}]_{\mathbf{X}} = (\nabla \mathbf{a})_{\text{skw}},$$

for all $\mathbf{a}, \mathbf{b} \in \mathcal{C}^1(\overline{\Omega})$. Displaying the cross product by this matrix makes the operation associative.

Additionally, we define $[\cdot]_{-\mathbf{X}} : \mathbb{R}^{3 \times 3} \rightarrow \mathbb{R}^3$, it is the left inverse of $[\cdot]_{\mathbf{X}}$ and given by

$$[\mathbf{A}]_{-\mathbf{X}} := \begin{pmatrix} \mathbf{A}_{3,2} \\ \mathbf{A}_{1,3} \\ \mathbf{A}_{2,1} \end{pmatrix}, \quad \text{for all } \mathbf{A} \in \mathbb{R}^{3 \times 3}.$$

For this mapping holds $[[\mathbf{a}]_{\mathbf{X}}]_{-\mathbf{X}} = \mathbf{a}$ and, thus $2[(\nabla \mathbf{a})_{\text{skw}}]_{-\mathbf{X}} = \nabla \times \mathbf{a}$, for all $\mathbf{a} \in \mathcal{C}^1(\overline{\Omega}; \mathbb{R}^3)$.

We also use the Levi–Civita tensor $\Upsilon \in \mathbb{R}^{3^3}$. Let \mathfrak{S}_3 be the symmetric group of all permutations of $(1, 2, 3)$. The sign of a given permutation $\sigma \in \mathfrak{S}_3$ is denoted by $\text{sgn } \sigma$. The Tensor Υ is defined via

$$\Upsilon_{ijk} := \begin{cases} \text{sgn } \sigma, & (i, j, k) = \sigma(1, 2, 3) \text{ with } \sigma \in \mathfrak{S}_3, \\ 0, & \text{else.} \end{cases}$$

This tensor allows it two write the cross product as

$$\mathbf{a} \times \mathbf{b} = \Upsilon : (\mathbf{a} \otimes \mathbf{b}) = \Upsilon_{ijk} \mathbf{a}_i \mathbf{b}_k, \quad \text{for all } \mathbf{a}, \mathbf{b} \in \mathbb{R}^d$$

and the curl via

$$\nabla \times \mathbf{d} = \Upsilon_{ijk} \partial_j \mathbf{d}_k, \quad \text{for all } \mathbf{d} \in \mathcal{C}^1(\Omega).$$

For a given Banach space V , Bochner–Lebesgue spaces are denoted by $L^p(0, T; V)$. Moreover, $W^{1,p}(0, T; V)$ denotes the Banach space of abstract functions in $L^p(0, T; V)$ whose weak time derivative exists and is again in $L^p(0, T; V)$ (see also Diestel and Uhl [10, Section II.2] or Roubíček [38, Section 1.5] for more details). By $\mathcal{C}([0, T]; V)$, and $\mathcal{C}_w([0, T]; V)$, we denote the spaces of abstract functions mapping $[0, T]$ into V that are absolutely continuous, continuous, and continuous with respect to the weak topology in V , respectively. We often omit the time interval $(0, T)$ and the domain Ω and just write, e.g., $L^p(\mathbf{W}^{k,p})$ for brevity.

Finally, by $c > 0$, we denote a generic positive constant and by C_δ a constant depending on δ .

2 Model and main result

Governing equations

Let Ω be of class $\mathcal{C}^{3,1}$. We consider the Ericksen–Leslie model as introduced in [29]. The governing equations read as

$$\partial_t \mathbf{v} + (\mathbf{v} \cdot \nabla) \mathbf{v} + \nabla p + \nabla \cdot \mathbf{T}^E - \nabla \cdot \mathbf{T}^L = \mathbf{g}, \quad (2.1a)$$

$$\mathbf{d} \times (\partial_t \mathbf{d} + (\mathbf{v} \cdot \nabla) \mathbf{d} - (\nabla \mathbf{v})_{\text{skw}} \mathbf{d} + \lambda (\nabla \mathbf{v})_{\text{sym}} \mathbf{d} + \mathbf{q}) = 0, \quad (2.1b)$$

$$\nabla \cdot \mathbf{v} = 0, \quad (2.1c)$$

$$|\mathbf{d}| = 1. \quad (2.1d)$$

We recall that $\mathbf{v} : \bar{\Omega} \times [0, T] \rightarrow \mathbb{R}^3$ denotes the velocity of the fluid, $\mathbf{d} : \bar{\Omega} \times [0, T] \rightarrow \mathbb{R}^3$ represents the orientation of the rod-like molecules, and $p : \bar{\Omega} \times [0, T] \rightarrow \mathbb{R}$ denotes the pressure. The Helmholtz free energy potential F , which is described rigorously in the next section, is assumed to depend only on the director and its gradient, $F = F(\mathbf{d}, \nabla \mathbf{d})$. The free energy functional \mathcal{F} is defined by

$$\mathcal{F} : \mathbf{H}^1 \rightarrow \mathbb{R}, \quad \mathcal{F}(\mathbf{d}) := \int_{\Omega} F(\mathbf{d}, \nabla \mathbf{d}) \, dx,$$

and \mathbf{q} is its variational derivative (see Furihata and Matsuo [23, Section 2.1]),

$$\mathbf{q} := \frac{\delta \mathcal{F}}{\delta \mathbf{d}}(\mathbf{d}) = \frac{\partial F}{\partial \mathbf{d}}(\mathbf{d}, \nabla \mathbf{d}) - \nabla \cdot \frac{\partial F}{\partial \nabla \mathbf{d}}(\mathbf{d}, \nabla \mathbf{d}). \quad (2.2a)$$

The Ericksen stress tensor \mathbf{T}^E is given by

$$\mathbf{T}^E = \nabla \mathbf{d}^T \frac{\partial F}{\partial \nabla \mathbf{d}}(\mathbf{d}, \nabla \mathbf{d}). \quad (2.2b)$$

The Leslie tensor is given by

$$\begin{aligned} \mathbf{T}^L &= \mu_1 (\mathbf{d} \cdot (\nabla \mathbf{v})_{\text{sym}} \mathbf{d}) \mathbf{d} \otimes \mathbf{d} + \mu_4 (\nabla \mathbf{v})_{\text{sym}} + (\mu_5 + \mu_6) (\mathbf{d} \otimes (\nabla \mathbf{v})_{\text{sym}} \mathbf{d})_{\text{sym}} \\ &\quad + (\mu_2 + \mu_3) (\mathbf{d} \otimes \mathbf{e})_{\text{sym}} + \lambda (\mathbf{d} \otimes (\nabla \mathbf{v})_{\text{sym}} \mathbf{d})_{\text{skw}} + (\mathbf{d} \otimes \mathbf{e})_{\text{skw}}, \end{aligned} \quad (2.2c)$$

where

$$\mathbf{e} := \partial_t \mathbf{d} + (\mathbf{v} \cdot \nabla) \mathbf{d} - (\nabla \mathbf{v})_{\text{skw}} \mathbf{d}. \quad (2.2d)$$

We emphasise that Parodi's law,

$$\lambda = \mu_2 + \mu_3, \quad (2.2e)$$

is neither essential for the reformulation nor the existence of measure-valued solutions, but is essential to prove the existence of *suitable* measure-valued solutions (see Remark 4) for which the weak-strong uniqueness property holds.

To ensure the dissipative character of the system, we assume that

$$\begin{aligned} \mu_1 &> 0, \quad \mu_4 > 0, \quad (\mu_5 + \mu_6) - \lambda(\mu_2 + \mu_3) > 0, \quad \mu_1 + \lambda(\mu_2 + \mu_3) > 0, \\ 4((\mu_5 + \mu_6) - \lambda(\mu_2 + \mu_3)) &> ((\mu_2 + \mu_3) - \lambda)^2. \end{aligned} \quad (2.2f)$$

Finally, we impose boundary and initial conditions as follows:

$$\mathbf{v}(\mathbf{x}, 0) = \mathbf{v}_0(\mathbf{x}) \quad \text{for } \mathbf{x} \in \Omega, \quad \mathbf{v}(\mathbf{x}, t) = \mathbf{0} \quad \text{for } (t, \mathbf{x}) \in [0, T] \times \partial\Omega, \quad (2.3a)$$

$$\mathbf{d}(\mathbf{x}, 0) = \mathbf{d}_0(\mathbf{x}) \quad \text{for } \mathbf{x} \in \Omega, \quad \mathbf{d}(\mathbf{x}, t) = \mathbf{d}_1(\mathbf{x}) \quad \text{for } (t, \mathbf{x}) \in [0, T] \times \partial\Omega. \quad (2.3b)$$

We always assume that $\mathbf{d}_1 = \mathbf{d}_0$ on $\partial\Omega$, which is a compatibility condition providing regularity.

The general Oseen–Frank energy

The *Oseen–Frank* energy is given by (see Leslie [31])

$$F(\mathbf{d}, \nabla \mathbf{d}) := \frac{K_1}{2} (\nabla \cdot \mathbf{d})^2 + \frac{K_2}{2} (\mathbf{d} \cdot \nabla \times \mathbf{d})^2 + \frac{K_3}{2} |\mathbf{d} \times \nabla \times \mathbf{d}|^2,$$

where $K_1, K_2, K_3 > 0$. This energy can be reformulated using the norm one restriction, to

$$2F(\mathbf{d}, \nabla \mathbf{d}) := k_1 (\nabla \cdot \mathbf{d})^2 + k_2 |\nabla \times \mathbf{d}|^2 + k_3 |\mathbf{d}|^2 (\nabla \cdot \mathbf{d})^2 + k_4 (\mathbf{d} \cdot \nabla \times \mathbf{d})^2 + k_5 |\mathbf{d} \times \nabla \times \mathbf{d}|^2. \quad (2.4)$$

where $k_1 = k_3 = K_1/2$, $k_2 = \min\{K_2, K_3\}/2$, $k_4 = K_2 - k_2$, and $k_5 = K_3 - k_2$ are again positive constants. We remark that $|\mathbf{d}|^2 |\nabla \times \mathbf{d}|^2 = (\mathbf{d} \cdot \nabla \times \mathbf{d})^2 + |\mathbf{d} \times \nabla \times \mathbf{d}|^2$.

We introduce short notations for the derivatives of the free energy (2.4) with respect to $\nabla \mathbf{d}$ and \mathbf{d} . The free energy (2.4) can be seen as a function $F : \mathbb{R}^3 \times \mathbb{R}^{3 \times 3} \rightarrow \mathbb{R}$, where we replace \mathbf{d} in definition (2.4) by $\mathbf{h} \in \mathbb{R}^3$ and $\nabla \mathbf{d}$ by $\mathbf{S} \in \mathbb{R}^{3 \times 3}$. By an easy vector calculation, we find

$$2F(\mathbf{h}, \mathbf{S}) = k_1 \operatorname{tr}(\mathbf{S})^2 + k_2 |(\mathbf{S})_{\text{skw}}|^2 + k_3 |\mathbf{h}|^2 \operatorname{tr}(\mathbf{S})^2 + k_4 ([\mathbf{h}]_{\mathbf{X}} : (\mathbf{S})_{\text{skw}})^2 + 4k_5 |(\mathbf{S})_{\text{skw}} \mathbf{h}|^2,$$

see Section 1 for the definition of the matrix $[\cdot]_{\mathbf{X}}$.

We abbreviate the derivative of F with respect to \mathbf{h} by $F_{\mathbf{h}}$ and the derivative with respect to \mathbf{S} by $F_{\mathbf{S}}$, where

$$F_{\mathbf{S}} : \mathbb{R}^3 \times \mathbb{R}^{3 \times 3} \rightarrow \mathbb{R}^{3 \times 3}, \quad \text{and} \quad F_{\mathbf{h}} : \mathbb{R}^3 \times \mathbb{R}^{3 \times 3} \rightarrow \mathbb{R}^3.$$

These derivatives are given by

$$\begin{aligned} F_{\mathbf{S}}(\mathbf{h}, \mathbf{S}) &= k_1 \operatorname{tr}(\mathbf{S}) I + k_2 (\mathbf{S})_{\text{skw}} + k_3 \operatorname{tr}(\mathbf{S}) |\mathbf{h}|^2 I + k_4 [\mathbf{h}]_{\mathbf{X}} : ([\mathbf{h}]_{\mathbf{X}} : (\mathbf{S})_{\text{skw}}) \\ &\quad + 4k_5 ((\mathbf{S})_{\text{skw}} \mathbf{h} \otimes \mathbf{h})_{\text{skw}} \\ F_{\mathbf{h}}(\mathbf{h}, \mathbf{S}) &= k_3 \operatorname{tr}(\mathbf{S})^2 \mathbf{h} + 2k_4 ([\mathbf{h}]_{\mathbf{X}} : (\mathbf{S})_{\text{skw}}) [(\mathbf{S})_{\text{skw}}]_{-\mathbf{X}} + 4k_5 (\mathbf{S})_{\text{skw}}^T (\mathbf{S})_{\text{skw}} \mathbf{h}, \end{aligned} \quad (2.5)$$

see Section 1 for the definition of $[\cdot]_{-\mathbf{X}}$.

To abbreviate, we define the tensor of order 4 $\boldsymbol{\Lambda} \in \mathbb{R}^{3^4}$ and a tensor of order 6 $\boldsymbol{\Theta} \in \mathbb{R}^{3^6}$ via

$$\boldsymbol{\Lambda}_{ijkl} := k_1 \boldsymbol{\delta}_{ij} \boldsymbol{\delta}_{kl} + k_2 (\boldsymbol{\delta}_{ik} \boldsymbol{\delta}_{jl} - \boldsymbol{\delta}_{il} \boldsymbol{\delta}_{jk}), \quad (2.6)$$

and

$$\begin{aligned} \boldsymbol{\Theta}_{ijklmn} &:= k_3 \boldsymbol{\delta}_{ij} \boldsymbol{\delta}_{lm} \boldsymbol{\delta}_{kn} + k_5 (\boldsymbol{\delta}_{ii} \boldsymbol{\delta}_{mn} \boldsymbol{\delta}_{jk} - \boldsymbol{\delta}_{mi} \boldsymbol{\delta}_{ln} \boldsymbol{\delta}_{jk} - \boldsymbol{\delta}_{lj} \boldsymbol{\delta}_{mn} \boldsymbol{\delta}_{ik} + \boldsymbol{\delta}_{jm} \boldsymbol{\delta}_{ln} \boldsymbol{\delta}_{ik}) \\ &\quad + k_4 (\boldsymbol{\delta}_{kn} \boldsymbol{\delta}_{jm} \boldsymbol{\delta}_{il} + \boldsymbol{\delta}_{km} \boldsymbol{\delta}_{jl} \boldsymbol{\delta}_{in} + \boldsymbol{\delta}_{kl} \boldsymbol{\delta}_{jn} \boldsymbol{\delta}_{im} - \boldsymbol{\delta}_{kn} \boldsymbol{\delta}_{jl} \boldsymbol{\delta}_{im} - \boldsymbol{\delta}_{km} \boldsymbol{\delta}_{jn} \boldsymbol{\delta}_{il} - \boldsymbol{\delta}_{kl} \boldsymbol{\delta}_{jm} \boldsymbol{\delta}_{in}), \end{aligned} \quad (2.7)$$

respectively. The free energy can be written as

$$2F(\mathbf{d}, \nabla \mathbf{d}) = \nabla \mathbf{d} : \boldsymbol{\Lambda} : \nabla \mathbf{d} + \nabla \mathbf{d} \otimes \mathbf{d} : \boldsymbol{\Theta} : \nabla \mathbf{d} \otimes \mathbf{d}. \quad (2.8)$$

The partial derivatives (2.5) inserted in definition (2.2a) gives the variational derivative in the case of the Oseen–Frank energy via

$$\begin{aligned} \boldsymbol{q} &= -k_1 \nabla \nabla \cdot \mathbf{d} - k_2 \nabla \times \nabla \times \mathbf{d} - k_3 \nabla (\nabla \cdot \mathbf{d} |\mathbf{d}|^2) - k_4 \nabla \cdot ([\mathbf{d}]_{\mathbf{X}} (\mathbf{d} \cdot \nabla \times \mathbf{d})) - 4k_5 \nabla \cdot ((\nabla \mathbf{d})_{\text{skw}} \mathbf{d} \otimes \mathbf{d})_{\text{skw}} \\ &\quad + k_3 (\nabla \cdot \mathbf{d})^2 \mathbf{d} + k_4 (\mathbf{d} \cdot \nabla \times \mathbf{d}) \nabla \times \mathbf{d} + 4k_5 (\nabla \mathbf{d})_{\text{skw}}^T (\nabla \mathbf{d})_{\text{skw}} \mathbf{d} \\ &= -\Delta_{\boldsymbol{\Lambda}} \mathbf{d} - \nabla \cdot (\mathbf{d} : \boldsymbol{\Theta} : \nabla \mathbf{d} \otimes \mathbf{d}) + \nabla \mathbf{d} : \boldsymbol{\Theta} : \nabla \mathbf{d} \otimes \mathbf{d}. \end{aligned} \quad (2.9)$$

The Tensor Λ is strongly elliptic, i.e. there is a $\eta > 0$ such that $\mathbf{a} \otimes \mathbf{b} : \Lambda : \mathbf{a} \otimes \mathbf{b} \geq \eta |\mathbf{a}|^2 |\mathbf{b}|^2$ for all $\mathbf{a}, \mathbf{b} \in \mathbb{R}^3$. Indeed, it holds

$$\mathbf{a} \otimes \mathbf{b} : \Lambda : \mathbf{a} \otimes \mathbf{b} = k_1(\mathbf{a} \cdot \mathbf{b})^2 + k_2(|\mathbf{a}|^2 |\mathbf{b}|^2 - (\mathbf{a} \cdot \mathbf{b})^2) \geq \min\{k_1, k_2\} |\mathbf{a}|^2 |\mathbf{b}|^2.$$

In the next section, we introduce the concept of measure-valued solutions. The proof of existence of measure-valued solutions to the Ericksen–Leslie model equipped with the Oseen–Frank energy is done by the author in [29]. We also refer to [29] for a more extensive introduction into the concept of generalized gradient Young measures.

Measure-valued solutions

Definition 1 (measure-valued solutions). The tuple $((\mathbf{v}, \mathbf{d}), (\mathbf{v}^o, m, \mathbf{v}^\infty), (\mu, \mathbf{v}^\mu))$ consisting of the pair (\mathbf{v}, \mathbf{d}) of velocity field \mathbf{v} and director field \mathbf{d} , the generalized gradient Young measure (μ, \mathbf{v}^μ) and the defect measure (μ, \mathbf{v}^μ) (see below) is said to be a measure-valued solution to (2.1) if

$$\begin{aligned} \mathbf{v} &\in L^\infty(0, T; \mathbf{L}_\sigma^2) \cap L^2(0, T; \mathbf{H}_{0,\sigma}^1) \cap W^{1,2}(0, T; (\mathbf{W}_{0,\sigma}^{1,3}(\Omega))^*), \\ \mathbf{d} &\in L^\infty(0, T; \mathbf{H}^1) \cap W^{1,2}(0, T; (\mathbf{L}^{3/2})), \\ \{\mathbf{v}_{(\mathbf{x},t)}^o\} &\subset \mathcal{P}(\mathbb{R}^{3 \times 3}), \text{ a.e. in } \Omega \times (0, T), \\ \{m_t\} &\subset \mathcal{M}^+(\overline{\Omega}), \text{ a.e. in } (0, T), \\ \{\mathbf{v}_{(\mathbf{x},t)}^\infty\} &\subset \mathcal{P}(\overline{B}_3 \times \mathcal{S}^{3^2-1}), m_t\text{-a.e. in } \overline{\Omega} \text{ and a.e. in } (0, T), \\ \{\mu_t\} &\subset \mathcal{M}^+(\overline{\Omega}), \text{ a.e. in } (0, T), \\ \{\mathbf{v}_{(\mathbf{x},t)}^\mu\} &\subset \mathcal{P}(\mathcal{S}^{3^2-1}), \mu_t\text{-a.e. in } \overline{\Omega} \text{ and a.e. in } (0, T) \end{aligned} \tag{2.10}$$

and if

$$\begin{aligned} &\int_0^T (\partial_t \mathbf{v}(t), \boldsymbol{\varphi}(t)) dt + \int_0^T ((\mathbf{v}(t) \cdot \nabla) \mathbf{v}(t), \boldsymbol{\varphi}(t)) dt - \int_0^T \langle \langle \mathbf{v}_t, \mathbf{S}^T F_{\mathbf{S}}(\mathbf{h}, \mathbf{S}) : \nabla \boldsymbol{\varphi}(t) \rangle \rangle dt \\ &- 2 \int_0^T \langle \langle \mu_t, \boldsymbol{\Gamma} : (\boldsymbol{\Gamma} \cdot \nabla \boldsymbol{\varphi}(t)) \rangle \rangle dt + \int_0^T (\mathbf{T}^L(t) : \nabla \boldsymbol{\varphi}(t)) dt = \int_0^T \langle \mathbf{g}(t), \boldsymbol{\varphi}(t) \rangle dt, \end{aligned} \tag{2.11a}$$

as well as

$$\int_0^T (\mathbf{d}(t) \times (\partial_t \mathbf{d}(t) + (\mathbf{v}(t) \cdot \nabla) \mathbf{d}(t) - (\nabla \mathbf{v}(t))_{\text{skw}} \mathbf{d}(t) + \lambda (\nabla \mathbf{v}(t))_{\text{sym}} + \mathbf{q}), \boldsymbol{\psi}(t)) dt = 0 \tag{2.11b}$$

with

$$\begin{aligned} \int_0^T (\mathbf{d}(t) \times \mathbf{q}, \boldsymbol{\psi}(t)) dt &= \int_0^T ([\mathbf{d}(t)] \mathbf{x} F_{\mathbf{S}}(\mathbf{d}(t), \nabla \mathbf{d}(t)); \nabla \boldsymbol{\psi}(t)) dt \\ &+ \int_0^T \langle \langle \mathbf{v}_t, \boldsymbol{\Upsilon} : (\mathbf{S}(F_{\mathbf{S}}(\mathbf{h}, \mathbf{S}))^T) \cdot \boldsymbol{\psi}(t) \rangle \rangle dt + \int_0^T \langle \langle \mathbf{v}_t, (\mathbf{h} \times F_{\mathbf{h}}(\mathbf{h}, \mathbf{S})) \cdot \boldsymbol{\psi}(t) \rangle \rangle dt, \end{aligned} \tag{2.11c}$$

holds for all $\boldsymbol{\varphi} \in \mathcal{C}_c^\infty(\Omega \times (0, T); \mathbb{R}^3)$ with $\nabla \cdot \boldsymbol{\varphi} = 0$ and $\boldsymbol{\psi} \in \mathcal{C}_c^\infty(\Omega \times (0, T); \mathbb{R}^3)$, respectively. Additionally, the norm restriction of the director holds, i.e. $|\mathbf{d}(\mathbf{x}, t)| = 1$ for a.e. $(\mathbf{x}, t) \in \Omega \times (0, T)$, the oscillation measure of a linear function is the gradient of the director

$$\int_{\mathbb{R}^{3 \times 3}} \mathbf{S} \mathbf{v}_{(\mathbf{x},t)}^o(d\mathbf{S}) = \nabla \mathbf{d}(\mathbf{x}, t) \quad \text{for a.e. } (\mathbf{x}, t) \in \Omega \times (0, T), \tag{2.12}$$

and the initial conditions $(\mathbf{v}_0, \mathbf{d}_0) \in \mathbf{L}_\sigma^2 \times \mathbf{H}^2$ with $\mathbf{d}_0 \in \mathbf{H}^{7/2}(\partial\Omega)$ shall be fulfilled in the weak sense and the boundary conditions in the sense of the trace. We remark that the trace is well defined for the function $\mathbf{d} \in L^\infty(0, T; \mathbf{H}^1)$, which is the expected value of the oscillation measure \mathbf{v}^o .

The dual pairings are defined as

$$\langle \langle \mu_t, f \rangle \rangle := \int_{\overline{\Omega}} \int_{\mathcal{S}^{3^2-1}} \sum_{i,j=1}^3 f(\boldsymbol{\Gamma}) v_{(\mathbf{x},t)}^\mu(d\boldsymbol{\Gamma}) \mu_t(d\mathbf{x})$$

for $f \in \mathcal{C}(\mathcal{S}^{3-1}; \mathbb{R})$ and

$$\begin{aligned}\langle\langle v_t, f \rangle\rangle &:= \int_{\Omega} \int_{\mathbb{R}^{d \times d}} f(\mathbf{x}, \mathbf{d}(\mathbf{x}, t), \mathbf{S}) v_{(\mathbf{x}, t)}^o(\mathrm{d}\mathbf{S}) \mathrm{d}\mathbf{x} \\ &\quad + \int_{\bar{\Omega}} \int_{\mathcal{S}^{d-1} \times \bar{B}_d} \tilde{f}(\mathbf{x}, \tilde{\mathbf{h}}, \tilde{\mathbf{S}}) v_{(\mathbf{x}, t)}^\infty(\mathrm{d}\tilde{\mathbf{S}}, \mathrm{d}\tilde{\mathbf{h}}) m_t(\mathrm{d}\mathbf{x})\end{aligned}$$

for $f \in \mathcal{R}$ (see (2.13) below).

We refer to the section 1 for the definition of the tensor \mathbf{T} and to (2.13) for the definition of the transformed function \tilde{f} .

Remark 1. We often abuse the notation by writing $\langle\langle v_t, f(\mathbf{h}, \mathbf{S}) \rangle\rangle$. Thereby we mean the generalized Young measure applied to the continuous function $(\mathbf{h}, \mathbf{S}) \mapsto f(\mathbf{h}, \mathbf{S})$.

We refer to the section 1 for the definition of the tensor \mathbf{T} . The transformed function $\tilde{f} : \bar{\Omega} \times B_d \times B_{d \times d} \rightarrow \mathbb{R}$, the so-called recession function is given by

$$\tilde{f}(\mathbf{x}, \tilde{\mathbf{h}}, \tilde{\mathbf{S}}) := f(\mathbf{x}, \frac{\tilde{\mathbf{h}}}{\sqrt{1 - |\tilde{\mathbf{h}}|^2}}, \frac{\tilde{\mathbf{S}}}{\sqrt{1 - |\tilde{\mathbf{S}}|^2}})(1 - |\tilde{\mathbf{h}}|^2)(1 - |\tilde{\mathbf{S}}|^2). \quad (2.13)$$

The class of function for which the above representation is valid are those functions, $f \in \mathcal{C}(\bar{\Omega} \times \mathbb{R}^d \times \mathbb{R}^{d \times d})$ such that \tilde{f} admits a continuous extension on the closure of its domain

$$\begin{aligned}\mathcal{F} := \left\{ f \in \mathcal{C}(\bar{\Omega} \times [0, T] \times \mathbb{R}^d \times \mathbb{R}^{d \times d}) \mid \exists \tilde{g} \in \mathcal{C}(\bar{\Omega} \times [0, T] \times \bar{B}_d \times \bar{B}_{d \times d}; \right. \\ \left. \tilde{f} = \tilde{g} \text{ on } \bar{\Omega} \times [0, T] \times B_d \times B_{d \times d} \right\}.\end{aligned}$$

An straightforward calculation shows that function with quadratic growth in \mathbf{h} and \mathbf{S} is unchanged by the transformation (2.13). Most of the appearing terms in the above definition have this structure. This implies that the transformation of $\mathbf{h} \times F_{\mathbf{h}}(\mathbf{h}, \mathbf{S})$ remains the function itself. Only the linear terms in $F_{\mathbf{S}}$ are changed by multiplying them with $1 - |\tilde{\mathbf{h}}|^2$, such that for example

$$\widetilde{\mathbf{S}^T F_{\mathbf{S}}}(\tilde{\mathbf{h}}, \tilde{\mathbf{S}}) = \tilde{\mathbf{S}}^T F_{\mathbf{S}}(\tilde{\mathbf{h}}, \tilde{\mathbf{S}}) - k_1 |\tilde{\mathbf{h}}|^2 \operatorname{tr}(\tilde{\mathbf{S}} \tilde{\mathbf{S}}^T) - k_2 |\tilde{\mathbf{h}}|^2 \tilde{\mathbf{S}}^T (\tilde{\mathbf{S}})_{\text{skw}}.$$

Remark 2. We often use some abuse of the notation by writing $\langle\langle v_t, f(\mathbf{h}, \mathbf{S}) \rangle\rangle$. Thereby we mean the generalized Young measure applied to the continuous function $(\mathbf{h}, \mathbf{S}) \mapsto f(\mathbf{h}, \mathbf{S})$.

In comparison to weak solutions (see [15]) The Ericksen stress \mathbf{T}^E and the variational derivative $\mathbf{d} \times \mathbf{q}$ are in this measure-valued formulation represented by generalized Young measures. A **generalized Young measure** on $\bar{\Omega} \times [0, T]$ with values in $\mathbb{R}^d \times \mathbb{R}^{d \times d}$ is a triple $(v_{(\mathbf{x}, t)}^o, m_t, v_{(\mathbf{x}, t)}^\infty)$ consisting of

- a parametrized family of probability measures $\{v_{(\mathbf{x}, t)}^o\} \subset \mathcal{P}(\mathbb{R}^{d \times d})$ for a.e. $(\mathbf{x}, t) \in \Omega \times (0, T)$,
- a positive measure $\{m_t\} \subset \mathcal{M}^+(\bar{\Omega})$, for a.e. $t \in (0, T)$ and
- a parametrized family of probability measures $\{v_{(\mathbf{x}, t)}^\infty\} \subset \mathcal{P}(\bar{B}_d \times \mathcal{S}^{d-1})$, for m_t -a.e. $\mathbf{x} \in \bar{\Omega}$ and a.e. $t \in (0, T)$.

As in [28, page 552] we call v^o *oscillation measure*, m_t *concentration measure* and v^∞ the *concentration angle measure*.

Definition 2. A generalized Young measure is said to be a generalized gradient Young measure, if there exists a sequence of Functions $\{\mathbf{d}_k\} \subset L^\infty(0, T; \mathbf{H}^2)$ with $\sup_k \|\nabla \mathbf{d}_k(|\mathbf{d}_k| + 1)\|_{L^\infty(L^2)} < \infty$ and a function $\mathbf{d} \in L^\infty(0, T; \mathbf{H}^1)$ such that the convergence

$$\int_{\Omega} f(\mathbf{x}, t, \mathbf{d}_k(\mathbf{x}, t), \nabla \mathbf{d}_k(\mathbf{x}, t)) \mathrm{d}\mathbf{x} \rightarrow \langle\langle v_t, f \rangle\rangle,$$

and formula (2.12) holds for all functions $f \in \mathcal{F}$ and (2.12) holds for a.e. $(\mathbf{x}, t) \in \Omega \times (0, T)$.

Remark 3. The approximating sequence \mathbf{d}_k can be chosen such that it fulfills the prescribed Dirichlet boundary conditions (compare to [29, Thm. 2.1]).

In the case of the Ericksen stress, an additional defect measure is of need to describe the limit of the regularised system we considered in [15]. A **defect measure** on $\overline{\Omega} \times [0, T]$ with values in $\mathbb{R}^{d \times d \times d}$ is a pair (μ_t, v^μ) consisting of

- a positive measure $\mu_t \in \mathcal{M}^+(\overline{\Omega})$, for a.e. $t \in (0, T)$ and
- a parametrized family of probability measures $\{v_y^\mu\}_{y \in \overline{\Omega}} \in \mathcal{P}(\mathcal{S}^{d^3-1})$, for μ_t -a.e. $x \in \overline{\Omega}$ and a.e. $t \in (0, T)$.

We refer to [29] for more details on the convergence in the sense of generalized Young measures.

In order to prove the weak-strong uniqueness property, it is needed that the measure-valued solution fulfills an additional assumption, it has to satisfy an energy inequality.

Definition 3 (Suitable measure-valued solutions). A measure-valued solution is said to be a suitable measure-valued solution if it fulfills Definition 1 and additionally the energy inequality

$$\begin{aligned} & \frac{1}{2} \|\mathbf{v}(t)\|_{L^2}^2 + \langle \mathbf{v}_t, F \rangle + \frac{1}{2} \langle \mu_t, 1 \rangle + \int_0^t (\mu_1 + \lambda(\mu_2 + \mu_3)) \|\mathbf{d} \cdot (\nabla \mathbf{v})_{\text{sym}} \mathbf{d}\|_{L^2}^2 ds \\ & + \int_0^t [\mu_4 \|(\nabla \mathbf{v})_{\text{sym}}\|_{L^2}^2 + (\mu_5 + \mu_6 - \lambda(\mu_2 + \mu_3)) \|(\nabla \mathbf{v})_{\text{sym}} \mathbf{d}\|_{L^2}^2 + \|\mathbf{d} \times \mathbf{q}\|_{L^2}^2] ds \\ & \leq \left(\frac{1}{2} \|\mathbf{v}_0\|_{L^2}^2 + \mathcal{F}(\mathbf{d}_0) \right) + \int_0^t [\langle \mathbf{g}, \mathbf{v} \rangle + ((\mu_2 + \mu_3) - \lambda) (\mathbf{d} \times \mathbf{q}, \mathbf{d} \times (\nabla \mathbf{v})_{\text{sym}} \mathbf{d})] ds. \end{aligned} \quad (2.14)$$

a.e. in $(0, T)$.

Remark 4 (Existence of suitable measure-valued solutions). In our recent work, we proved the existence of measure-valued solutions to the system (2.1) in the sense of Definition 1. The proof relies on the existence of weak solutions $(\mathbf{v}_\delta, \mathbf{d}_\delta)$ to a regularised system, where the free energy (2.4) is changed by adding the regularising term $\delta |\Delta \mathbf{d}|^2$. We obtain generalized Young measure-valued solutions for vanishing regularisation.

In the case of Parodi's relation (2.2e), we can prove the following energy inequality for vanishing regularisation (see [29])

$$\begin{aligned} & \frac{1}{2} \|\mathbf{v}_\delta(t)\|_{L^2}^2 + \langle \mathbf{v}_t, F \rangle + \frac{1}{2} \langle \mu_t, 1 \rangle + (\mu_1 + \lambda^2) \int_0^t \|\mathbf{d}_\delta \cdot (\nabla \mathbf{v}_\delta)_{\text{sym}} \mathbf{d}_\delta\|_{L^2}^2 ds \\ & + \int_0^t [\mu_4 \|(\nabla \mathbf{v}_\delta)_{\text{sym}}\|_{L^2}^2 + (\mu_5 + \mu_6 - \lambda^2) \|(\nabla \mathbf{v}_\delta)_{\text{sym}} \mathbf{d}_\delta\|_{L^2}^2 + \|\mathbf{d} \times \mathbf{q}\|_{L^2}^2 - \langle \mathbf{g}, \mathbf{v}_\delta \rangle] ds \\ & \leq \frac{1}{2} \|\mathbf{v}_\delta(0)\|_{L^2}^2 + \frac{\delta}{2} \|\Delta \mathbf{d}_\delta(0)\|_{L^2}^2 + \mathcal{F}(\mathbf{d}(0)). \end{aligned} \quad (2.15)$$

In the case that Parodi's relation (2.2e) does not hold, we can not prove the existence of suitable measure-valued solutions.

Strong solutions and main theorem

Definition 4 (strong solution). A pair $(\tilde{\mathbf{v}}, \tilde{\mathbf{d}})$ is called a strong solution if it fulfills system (2.1) and exhibits the regularity

$$\begin{aligned} \tilde{\mathbf{v}} & \in L^\infty(0, T; \mathbf{L}_\sigma^2) \cap L^2(0, T; \mathbf{L}^\infty) \cap L^2(0, T; \mathbf{W}_{0,\sigma}^{1,3}) \cap L^1(0, T; \mathbf{W}^{1,\infty}) \cap W^{1,2}(0, T; (\mathbf{H}_{0,\sigma}^1)^*), \\ \tilde{\mathbf{d}} & \in L^\infty(0, T; \mathbf{W}^{1,\infty}) \cap L^2(0, T; \mathbf{W}^{2,3}) \cap L^4(0, T; \mathbf{W}^{1,6}) \cap W^{1,1}(0, T; \mathbf{W}^{1,3} \cap \mathbf{L}^\infty), \end{aligned}$$

as well as

$$\nabla \tilde{\mathbf{v}} \in \mathcal{C}(\overline{\Omega} \times [0, T]) \quad \text{and} \quad \nabla \tilde{\mathbf{d}} \in \mathcal{C}(\overline{\Omega} \times [0, T]). \quad (2.16)$$

Remark 5. The continuity assumptions on the solution are especially needed to be able to insert the formulation of the measure-valued solution (see Definition 1). This assumptions can presumably be generalized. The assumptions on the differentiability of $\tilde{\mathbf{v}}$ with respect to t follows from equation (2.1a) and the other assumptions on $\tilde{\mathbf{v}}$.

Remark 6. To our best knowledge, there is no existence result in the class of strong solutions for the Ericksen-Leslie model in its full generality. But there are similar results for simpler models, and they can possibly be generalized to the case presented here.

Theorem 2.1. Let (\mathbf{v}, \mathbf{d}) be a suitable measure-valued solution to the Ericksen–Leslie model according to Definition 1 and $(\tilde{\mathbf{v}}, \tilde{\mathbf{d}})$ a strong solution according to definition 4 to the same initial and boundary values. Then both solutions coincide, i.e. $\mathbf{v} = \tilde{\mathbf{v}}$ and $\mathbf{d} = \tilde{\mathbf{d}}$ in $\Omega \times (0, T)$ such that

$$\mathbf{v}^o = \delta_{\nabla \tilde{\mathbf{d}}}, \quad \mu_t \equiv 0, \quad \text{and } m_t \equiv 0, .$$

The proof of this main result is carried out in the following sections. It is a simple consequence of Proposition 6.1.

Remark 7. In Proposition 6.1 even a continuous dependence on the initial values is proven.

3 Preliminaries

This section collects some auxiliary lemmas and the energy equality for the strong solution.

Lemma 3.1. Let $f, g \in L^1(0, T)$ be fulfilling

$$-\int_0^T \phi'(t)g(t) dt \leq \int_0^T \phi(t)f(t) dt \quad \text{for all } \phi \in \mathcal{C}_c^\infty(0, T) \text{ with } \phi \geq 0.$$

Then it holds

$$g(t) - g(s) \leq \int_s^t f(\tau) d\tau \quad \text{for a.e. } t, s \in (0, T).$$

Proof. The proof is similar to the proof of the standard variational lemma, just by using the sequence of smooth functions $\phi_\varepsilon \in \mathcal{C}_c^\infty(0, T)$ defined by

$$\phi_\varepsilon(\tau) := (\rho_\varepsilon * \chi_{[t,s]})(\tau) = \int_0^T \rho(\tau - r) \chi_{[t,s]}(r) dr.$$

Here the functions ρ_ε are the usual mollifier functions (see Emmrich [13, Definition 3.1.6]) and $\chi_{[t,s]}$ denotes the characteristic function. The limit $\varepsilon \rightarrow 0$ gives the assertion. \square

Lemma 3.2. For the measure-valued solution, there holds

$$\partial_t \mathbf{d} = [\mathbf{d}]_X^T [\mathbf{d}]_X \partial_t \mathbf{d} \quad \text{in } L^2(0, T; \mathbf{L}^{3/3}).$$

For the strong solution, there holds

$$\partial_t \tilde{\mathbf{d}} = [\tilde{\mathbf{d}}]_X^T [\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}} \quad \text{in } L^\infty(0, T; \mathbf{L}^\infty) \cap L^\infty(0, T; \mathbf{W}^{1,3}).$$

Proof. First, we remark that $\mathbf{d} \in L^\infty(0, T; \mathbf{L}^\infty)$. Indeed, $|\mathbf{d}(\mathbf{x}, t)| = 1$ a.e. in $\Omega \times (0, T)$. For a test function $\boldsymbol{\varphi} \in \mathcal{C}_c^\infty(\Omega \times (0, T))$ we can calculate

$$\begin{aligned} \int_0^T [[\mathbf{d}]_X^T [\mathbf{d}]_X \partial_t \mathbf{d}, \boldsymbol{\varphi}] dt &= \int_0^T [(|\mathbf{d}|^2 I - \mathbf{d} \otimes \mathbf{d}) \partial_t \mathbf{d}, \boldsymbol{\varphi}] dt \\ &= \int_0^T [\partial_t \mathbf{d}, \boldsymbol{\varphi}] dt + \frac{1}{2} \int_0^T [|\mathbf{d}|^2, \partial_t (\mathbf{d} \cdot \boldsymbol{\varphi})] dt = \int_0^T [\partial_t \mathbf{d}, \boldsymbol{\varphi}] dt. \end{aligned}$$

The weak time derivative of $|\mathbf{d}|^2$ vanishes, since it is constant a.e. in $\Omega \times (0, T)$. The density of the test functions in $L^2(0, T; \mathbf{L}^6)$ proves the assertion.

The proof for the strong solution is similar. The additional regularity is due to (2.16). \square

Lemma 3.3. For the measure-valued solution, there holds

$$(\mathbf{v} \cdot \nabla) \mathbf{d} = [\mathbf{d}]_X^T [\mathbf{d}]_X (\mathbf{v} \cdot \nabla) \mathbf{d} \quad \text{in } L^2(0, T; \mathbf{L}^{3/2}).$$

For the strong solution, there holds

$$(\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} = [\tilde{\mathbf{d}}]_X^T [\tilde{\mathbf{d}}]_X (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} \quad \text{in } L^\infty(0, T; \mathbf{L}^\infty) \cap L^\infty(0, T; \mathbf{W}^{1,3}).$$

Proof. For a test function $\psi \in \mathcal{C}_c^\infty(\Omega \times (0, T))$ we can calculate

$$\begin{aligned} \int_0^T ([\mathbf{d}]_X^T [\mathbf{d}]_X (\mathbf{v} \cdot \nabla) \mathbf{d}, \psi) dt &= \int_0^T (|\mathbf{d}|^2 (\mathbf{v} \cdot \nabla) \mathbf{d}, \psi) dt - \frac{1}{2} \int_0^T ((\mathbf{v} \cdot \nabla) |\mathbf{d}|^2, \psi \cdot \mathbf{d}) dt \\ &= \int_0^T ((\mathbf{v} \cdot \nabla) \mathbf{d}, \psi) dt + \frac{1}{2} \int_0^T [(|\mathbf{d}|^2, \nabla \cdot (\mathbf{v} \psi \cdot \mathbf{d}))] dt. \end{aligned}$$

Since $|\mathbf{d}| = 1$ a. e. in $\Omega \times (0, T)$, the weak derivatives of $|\mathbf{d}|^2$ in the last term on the right-hand side of the forgoing equation vanishes. The density of the test functions in $L^2(0, T; \mathbf{L}^3)$ proves the assertion.

The proof for the strong solution is similar. The additional regularity is due to (2.16). \square

Corollary 3.1. For the measure-valued solution, there holds

$$\mathbf{e} = \partial_t \mathbf{d} + (\mathbf{v} \cdot \nabla) \mathbf{d} - (\nabla \mathbf{v})_{\text{skw}} \mathbf{d} = [\mathbf{d}]_X^T [\mathbf{d}]_X \mathbf{e} = -[\mathbf{d}]_X^T [\mathbf{d}]_X (\lambda (\nabla \mathbf{v})_{\text{sym}} \mathbf{d} + \mathbf{q}) \quad \text{in } L^2(0, T; \mathbf{L}^2).$$

For the strong solution, there holds

$$\tilde{\mathbf{e}} = \partial_t \tilde{\mathbf{d}} + (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} - (\nabla \tilde{\mathbf{v}})_{\text{skw}} \tilde{\mathbf{d}} = [\tilde{\mathbf{d}}]_X^T [\tilde{\mathbf{d}}]_X \tilde{\mathbf{e}} = -[\tilde{\mathbf{d}}]_X^T [\tilde{\mathbf{d}}]_X (\lambda (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}} + \tilde{\mathbf{q}}) \quad \text{in } L^\infty(0, T; \mathbf{L}^3).$$

Proof. Since $[\mathbf{d}]_X^T [\mathbf{d}]_X = I - \mathbf{d} \otimes \mathbf{d}$ and $\mathbf{d} \cdot (\nabla \mathbf{v})_{\text{skw}} \mathbf{d} = 0$, the identity $\mathbf{e} = [\mathbf{d}]_X^T [\mathbf{d}]_X \mathbf{e}$ in $L^2(0, T; \mathbf{L}^{6/5})$ is obvious from Lemma 3.2 and Lemma 3.3.

The term \mathbf{e} can be estimated by

$$\|\mathbf{e}\|_{L^2(L^{3/2})} \leq \|\partial_t \mathbf{d}\|_{L^2(L^{3/2})} + \|\mathbf{v}\|_{L^2(\mathbf{L}^6)} \|\mathbf{d}\|_{L^\infty(\mathbf{W}^{1,2})} + \|\mathbf{v}\|_{L^2(\mathbf{H}^1)} \|\mathbf{d}\|_{L^\infty(\mathbf{L}^6)}.$$

Due to the energy inequality (2.14) and equation (2.11b) we can estimate

$$\begin{aligned} \|[\mathbf{d}]_X^T [\mathbf{d}]_X \mathbf{e}\|_{L^2(\mathbf{L}^2)} &\leq \|\mathbf{d}\|_{L^\infty(\mathbf{L}^6)} \left(\|\lambda \mathbf{d} \times (\nabla \mathbf{v})_{\text{sym}} \mathbf{d}\|_{L^2(\mathbf{L}^2)} + \|\mathbf{d} \times \mathbf{q}\|_{L^2(\mathbf{L}^2)} \right) \\ &\leq c \|\mathbf{d}\|_{L^\infty(\mathbf{L}^\infty)} \left(\|\mathbf{d} \cdot (\nabla \mathbf{v})_{\text{sym}} \mathbf{d}\|_{L^2(\mathbf{L}^2)} + \|(\nabla \mathbf{v})_{\text{sym}} \mathbf{d}\|_{L^2(\mathbf{L}^2)} + \|\mathbf{d} \times \mathbf{q}\|_{L^2(\mathbf{L}^2)} \right). \end{aligned}$$

The density of the test functions in $L^2(\mathbf{L}^2)$ and equation (2.11b) gives the first assertion. The second one follows similar. \square

Proposition 3.1. Let $(\tilde{\mathbf{v}}, \tilde{\mathbf{d}})$ be a strong solution according to Definition 4. Then it fulfils the following energy equality

$$\begin{aligned} &\frac{1}{2} \|\tilde{\mathbf{v}}(t)\|_{\mathbf{L}^2}^2 + \mathcal{F}(\tilde{\mathbf{d}}(t)) + \int_0^t [(\mu_1 + \lambda(\mu_2 + \mu_3)) \|\tilde{\mathbf{d}} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}\|_{\mathbf{L}^2}^2 + \mu_4 \|(\nabla \tilde{\mathbf{v}})_{\text{sym}}\|_{\mathbf{L}^2}^2] ds \\ &+ \int_0^t [(\mu_5 + \mu_6 - \lambda(\mu_2 + \mu_3)) \|(\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}\|_{\mathbf{L}^2}^2 + \|\tilde{\mathbf{d}} \times \tilde{\mathbf{q}}\|_{\mathbf{L}^2}^2] ds \\ &= \left(\frac{1}{2} \|\tilde{\mathbf{v}}(0)\|_{\mathbf{L}^2}^2 + \mathcal{F}(\tilde{\mathbf{d}}(0)) \right) + \int_0^t [\langle \mathbf{g}, \tilde{\mathbf{v}} \rangle + ((\mu_2 + \mu_3) - \lambda) (\tilde{\mathbf{d}} \times \tilde{\mathbf{q}}, \tilde{\mathbf{d}} \times (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}})] ds. \end{aligned}$$

Proof. The proof is similar to the proof of Proposition 4.1 in [15]. Here we only focus on the necessary modification. Equation (2.1a) is tested with $\tilde{\mathbf{v}}$ and equation (2.1b) with $\tilde{\mathbf{d}} \times \tilde{\mathbf{q}}$. Remark that the identities $[\tilde{\mathbf{d}}]_X^T [\tilde{\mathbf{d}}]_X = I - \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}$, $\partial_t |\tilde{\mathbf{d}}|^2 = 0$ as well as $(\tilde{\mathbf{v}} \cdot \nabla) |\tilde{\mathbf{d}}|^2 = 0$ hold for the unit vector $\tilde{\mathbf{d}}$. For the Leslie-stress \mathbf{T}^L tested with $(\nabla \tilde{\mathbf{v}})_{\text{sym}}$, we use Corollary 3.1 to be able to insert equation (2.1b) twice

$$\begin{aligned} &(\tilde{\mathbf{T}}^L; (\nabla \tilde{\mathbf{v}})_{\text{sym}}) - (\tilde{\mathbf{d}} \times (\nabla \tilde{\mathbf{v}})_{\text{skw}} \tilde{\mathbf{d}}, \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}) \\ &= \mu_1 |\tilde{\mathbf{d}} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}|^2 + \mu_4 |(\nabla \tilde{\mathbf{v}})_{\text{sym}}|^2 + (\mu_5 + \mu_6) |(\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}|^2 + (\mu_2 + \mu_3) (\tilde{\mathbf{e}}, (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}) \\ &\quad + \lambda ((\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, (\nabla \tilde{\mathbf{v}})_{\text{skw}} \tilde{\mathbf{d}}) + (\tilde{\mathbf{e}}, (\nabla \tilde{\mathbf{v}})_{\text{skw}} \tilde{\mathbf{d}}) - (\tilde{\mathbf{d}} \times (\nabla \tilde{\mathbf{v}})_{\text{skw}} \tilde{\mathbf{d}}, \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}) \\ &= (\mu_1 + \lambda(\mu_2 + \mu_3)) |\tilde{\mathbf{d}} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}|^2 + \mu_4 |(\nabla \tilde{\mathbf{v}})_{\text{sym}}|^2 + (\mu_5 + \mu_6 - \lambda(\mu_2 + \mu_3)) |(\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}|^2 \\ &\quad - (\mu_2 + \mu_3) (\tilde{\mathbf{d}} \times \tilde{\mathbf{q}}, \tilde{\mathbf{d}} \times (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}). \end{aligned}$$

Summing up both tested equations and integrating in time gives the desired energy equality. \square

Lemma 3.4. Let \mathbf{v} be a measure-valued solution (see Definition 1) and $\tilde{\mathbf{v}}$ a strong solution (see Definition 4). It holds that

$$\|\mathbf{v} - \tilde{\mathbf{v}}\|_{\mathbf{L}^6}^2 + \|(\nabla \mathbf{v})_{\text{skw}} - (\nabla \tilde{\mathbf{v}})_{\text{skw}}\|_{\mathbf{L}^2}^2 \leq c \|(\nabla \mathbf{v})_{\text{sym}} - (\nabla \tilde{\mathbf{v}})_{\text{sym}}\|_{\mathbf{L}^2}^2.$$

Proof. This Lemma is a simple application of the embedding in three dimensions and Korn's inequality (see McLean [37, Theorem 10.1]). \square

4 Relative energy

The relative energy is defined by

$$\begin{aligned} \mathcal{E}(t) := & \frac{1}{2} \|\mathbf{v}(t) - \tilde{\mathbf{v}}(t)\|_{\mathbf{L}^2}^2 + \frac{1}{2} \langle\langle \mu_t, 1 \rangle\rangle + \frac{1}{2} \langle\langle \mathbf{v}_t, (\mathbf{S} - \nabla \tilde{\mathbf{d}}(t)) : \boldsymbol{\Lambda} : (\mathbf{S} - \nabla \tilde{\mathbf{d}}(t)) \rangle\rangle \\ & + \frac{1}{2} \langle\langle \mathbf{v}_t, (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}}(t) \otimes \tilde{\mathbf{d}}(t)) : \boldsymbol{\Theta} : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}}(t) \otimes \tilde{\mathbf{d}}(t)) \rangle\rangle \end{aligned} \quad (4.1)$$

and the relative dissipation by

$$\begin{aligned} \mathcal{W}(t) := & (\mu_1 + \lambda(\mu_2 + \mu_3)) \|\mathbf{d}(t) \cdot (\nabla \mathbf{v}(t))_{\text{sym}} \mathbf{d}(t) - \tilde{\mathbf{d}}(t) \cdot (\nabla \tilde{\mathbf{v}}(t))_{\text{sym}} \tilde{\mathbf{d}}(t)\|_{\mathbf{L}^2}^2 \\ & + (\mu_5 + \mu_6 - \lambda(\mu_2 + \mu_3)) \|(\nabla \mathbf{v}(t))_{\text{sym}} \mathbf{d}(t) - (\nabla \tilde{\mathbf{v}}(t))_{\text{sym}} \tilde{\mathbf{d}}(t)\|_{\mathbf{L}^2}^2 \\ & + \mu_4 \|(\nabla \mathbf{v}(t))_{\text{sym}} - (\nabla \tilde{\mathbf{v}}(t))_{\text{sym}}\|_{\mathbf{L}^2}^2 + \|\mathbf{d}(t) \times \mathbf{q}(t) - \tilde{\mathbf{d}}(t) \times \tilde{\mathbf{q}}(t)\|_{\mathbf{L}^2}^2. \end{aligned} \quad (4.2)$$

Inserting the definitions of the tensors $\boldsymbol{\Lambda}$ and $\boldsymbol{\Theta}$, the relative energy can be expressed as

$$\begin{aligned} \mathcal{E}(t) = & \frac{1}{2} \int_{\Omega} \mu_t(\mathbf{d}) + \frac{1}{2} \langle\langle \mathbf{v}_t, k_1(\text{tr}(\mathbf{S}) - \text{tr}(\nabla \tilde{\mathbf{d}}(t)))^2 + 2k_2 |(\mathbf{S})_{\text{skw}} - (\nabla \tilde{\mathbf{d}}(t))_{\text{skw}}|^2 \rangle\rangle \\ & + \frac{1}{2} \langle\langle \mathbf{v}_t, k_3 |\text{tr}(\mathbf{S}) \mathbf{h} - (\nabla \cdot \tilde{\mathbf{d}}(t)) \tilde{\mathbf{d}}(t)|^2 + k_4 ((\mathbf{S})_{\text{skw}} : [\mathbf{h}]_{\mathbf{x}} - (\nabla \tilde{\mathbf{d}}(t))_{\text{skw}} : [\tilde{\mathbf{d}}(t)]_{\mathbf{x}})^2 \rangle\rangle \\ & + \frac{1}{2} \langle\langle \mathbf{v}_t, 4k_5 |(\mathbf{S})_{\text{skw}} \mathbf{h} - (\nabla \tilde{\mathbf{d}}(t))_{\text{skw}} \tilde{\mathbf{d}}(t)|^2 \rangle\rangle + \frac{1}{2} \|\mathbf{v}(t) - \tilde{\mathbf{v}}(t)\|_{\mathbf{L}^2}^2. \end{aligned}$$

We remark that due to the regularity shown in [29], there holds $\mathcal{E} \in L^\infty(0, T)$ and $\mathcal{W} \in L^1(0, T)$. The overall goal is to apply Gronwalls estimate to the inequality of the form

$$\mathcal{E}(t) + \int_0^t \mathcal{W}(s) \, ds \leq c_0 + \zeta \int_0^t \mathcal{W}(s) \, ds + \int_0^t \mathcal{K}(s) \mathcal{E}(s) \, ds,$$

with $\zeta < 1$ and $\mathcal{K} \in L^1(0, T)$. The constant c_0 is a constant depending on the initial values of the solutions (\mathbf{v}, \mathbf{d}) and $(\tilde{\mathbf{v}}, \tilde{\mathbf{d}})$, which vanishes if those are the same.

The term \mathcal{K} is given by

$$\begin{aligned} \mathcal{K}(s) = & C_{\delta} (\|\tilde{\mathbf{v}}(s)\|_{\mathbf{L}^\infty}^2 + \|\tilde{\mathbf{v}}(s)\|_{\mathbf{W}^{1,3}}^2 + \|\tilde{\mathbf{d}}(s)\|_{\mathbf{W}^{2,3}}^2 + \|\tilde{\mathbf{d}}(s)\|_{\mathbf{W}^{1,6}}^4 + \|\partial_t \tilde{\mathbf{d}}(s)\|_{\mathbf{L}^\infty} + \|\partial_t \tilde{\mathbf{d}}(s)\|_{\mathbf{W}^{1,3}}) \\ & + C_{\delta} (\|(\nabla \tilde{\mathbf{v}}(s))_{\text{sym}}\|_{\mathbf{L}^\infty} + 1), \end{aligned}$$

where C_{δ} is a possible large constant depending on the norms $\|\mathbf{d}\|_{L^\infty(\mathbf{L}^\infty)}$, $\|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{W}^{1,\infty})}$ and δ . It is obvious that \mathcal{K} is bounded in $L^1(0, T)$ due to the regularity of the strong solution. For the generalized Young measures, a result similar to the Sobolev embedding theorem holds true.

Lemma 4.1. Let \mathbf{d} be a measure-valued solution (see Definition 1) and $(\mathbf{v}, m, \mathbf{v}^\infty)$ the associated generalized Young measure and $\tilde{\mathbf{d}}$ a strong solution (see Definition 4) to the same initial and boundary values. Let the associated relative Energy be given as above (see (4.1)). Then there exists a constant $c > 0$ such that

$$\int_0^t \|\mathbf{d}(s) - \tilde{\mathbf{d}}(s)\|_{\mathbf{L}^6}^2 \, ds + \int_0^t \langle\langle \mathbf{v}_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}(s)|^2 \rangle\rangle \, ds \leq c \int_0^t \mathcal{E}(s) \, ds, \quad (4.3a)$$

$$\int_0^t \langle\langle \mathbf{v}_s, |\boldsymbol{\Theta} : (\mathbf{S} \otimes \mathbf{d} - \nabla \tilde{\mathbf{d}}(s) \otimes \tilde{\mathbf{d}}(s))|^2 \rangle\rangle \, ds \leq c \int_0^t \mathcal{E}(s) \, ds, \quad (4.3b)$$

$$\int_0^t \langle\langle \mathbf{v}_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}(s)|^2 |\mathbf{h} - \tilde{\mathbf{d}}(s)|^2 \rangle\rangle \, ds \leq c (1 + \|\nabla \tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^3)}^2 + \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^2) \int_0^t \mathcal{E}(s) \, ds, \quad (4.3c)$$

$$\int_0^t \|\mathbf{d}(s) - \tilde{\mathbf{d}}(s)\|_{\mathbf{L}^{12}}^4 \, ds \leq c (1 + \|\nabla \tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^3)}^2 + \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^2) \int_0^t \mathcal{E}(s) \, ds. \quad (4.3d)$$

Proof. Since \mathbf{d} and $\tilde{\mathbf{d}}$ are measure and strong solutions to the same boundary data, the difference fulfills Dirichlet boundary conditions, i. e. $\mathbf{d}(t) - \tilde{\mathbf{d}}(t) \in \mathbf{H}_0^1$. The function $\mathbf{d} - \tilde{\mathbf{d}}$ gives rise to a generalized gradient Young measure. Due to Definition 2 there exists a sequence $\{\mathbf{d}_n\}$ such that $\mathbf{d}_n \rightarrow \mathbf{d} - \tilde{\mathbf{d}}$ in the sense of generalized Young measures. Due to Remark 3 the sequence fulfills homogeneous boundary conditions. For \mathbf{d}_n , [29, Proposition 4.2] yields

$$\begin{aligned} \|\mathbf{d}_n(t)\|_{\mathbf{H}_0^1}^2 &\leq c(\nabla \mathbf{d}_n(t); \boldsymbol{\Lambda} : \nabla \mathbf{d}_n(t)), \\ \int_{\Omega} |\mathbf{d}_n(\mathbf{x}, t)|^2 |\nabla \mathbf{d}_n(\mathbf{x}, t)|^2 d\mathbf{x} &\leq c(\nabla \mathbf{d}_n(t) \otimes \mathbf{d}_n(t); \boldsymbol{\Theta} : \nabla \mathbf{d}_n(t) \otimes \mathbf{d}_n(t)). \end{aligned} \quad (4.4)$$

The convergence result for generalized gradient Young measures for the sequence $\{\mathbf{d}_n\}$ applied to the test functions $(\mathbf{h}, \mathbf{S}) \mapsto |\mathbf{S}|^2$ gives

$$\int_0^T \phi(t) |\nabla \mathbf{d}_n(t)|^2 dt \rightarrow \int_0^T \phi(t) \langle\langle v_t, |\mathbf{S} - \nabla \tilde{\mathbf{d}}(t)|^2 \rangle\rangle dt$$

and for the test function $(\mathbf{h}, \mathbf{S}) \mapsto \mathbf{S} : \boldsymbol{\Lambda} : \mathbf{S}$ the convergence

$$\int_0^T \phi(t) (\nabla \mathbf{d}_n(t); \boldsymbol{\Lambda} : \nabla \mathbf{d}_n(t)) dt \rightarrow \int_0^T \phi(t) \langle\langle v_t, (\mathbf{S} - \nabla \tilde{\mathbf{d}}(t)) : \boldsymbol{\Lambda} : (\mathbf{S} - \nabla \tilde{\mathbf{d}}(t)) \rangle\rangle dt$$

Since m_t is a positive measure and v^o is a probability measure, we get with Jensen's inequality

$$\begin{aligned} \int_0^T \phi(t) \langle\langle v_t, |\mathbf{S} - \nabla \tilde{\mathbf{d}}(t)|^2 \rangle\rangle dt &\geq \int_0^T \phi(t) \int_{\Omega} \int_{\mathbb{R}^{3 \times 3}} |\mathbf{S} - \nabla \tilde{\mathbf{d}}(\mathbf{x}, t)|^2 v_{(\mathbf{x}, t)}^o(d\mathbf{S}) d\mathbf{x} dt \\ &\geq \int_0^T \phi(t) \int_{\Omega} \left| \int_{\mathbb{R}^{3 \times 3}} (\mathbf{S} - \nabla \tilde{\mathbf{d}}(\mathbf{x}, t)) v_{(\mathbf{x}, t)}^o(d\mathbf{S}) \right|^2 d\mathbf{x} dt. \end{aligned}$$

The right-hand side can be identified due to (2.12) with the weak limit $\nabla \mathbf{d}$, such that with the embedding in three dimensions

$$\int_0^T \phi(t) \langle\langle v_t, |\mathbf{S} - \nabla \tilde{\mathbf{d}}(t)|^2 \rangle\rangle dt \geq \int_0^T \phi(t) \|\nabla \mathbf{d}(t) - \nabla \tilde{\mathbf{d}}(t)\|_{\mathbf{L}^2}^2 dt \geq c \int_0^T \phi(t) \|\mathbf{d}(t) - \tilde{\mathbf{d}}(t)\|_{\mathbf{L}^6}^2 dt.$$

With the fundamental lemma of variational calculus the estimate (4.3a) is obvious.

The estimate (4.3b) follows, if one uses the purely algebraic inequality

$$|\boldsymbol{\Theta} : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})|^2 \leq c(\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) : \boldsymbol{\Theta} : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}). \quad (4.5)$$

The proof of the inequality (4.5) is deferred to the Appendix.

Choosing the same approximating sequence as above, application of the convergence result for generalized gradient Young measures applied to the test function $(\mathbf{h}, \mathbf{S}) \mapsto |\mathbf{h}|^2 |\mathbf{S}|^2$ gives

$$\int_0^T \phi(t) \|\nabla \mathbf{d}_n(t) \otimes \mathbf{d}_n(t)\|^2 dt \rightarrow \int_0^T \phi(t) \langle\langle v_t, |\mathbf{S} - \nabla \tilde{\mathbf{d}}(t)|^2 |\mathbf{h} - \tilde{\mathbf{d}}(t)|^2 \rangle\rangle dt,$$

and for the test function $(\mathbf{h}, \mathbf{S}) \mapsto \mathbf{S} \otimes \mathbf{h} : \boldsymbol{\Theta} : \mathbf{S} \otimes \mathbf{h}$ the convergence

$$\begin{aligned} \int_0^T \phi(t) (\nabla \mathbf{d}_n(t) \otimes \mathbf{d}_n(t); \boldsymbol{\Theta} : \nabla \mathbf{d}_n(t) \otimes \mathbf{d}_n(t)) dt &\rightarrow \\ \int_0^T \phi(t) \langle\langle v_t, (\mathbf{S} - \nabla \tilde{\mathbf{d}}(t)) \otimes (\mathbf{h} - \tilde{\mathbf{d}}(t)) : \boldsymbol{\Theta} : (\mathbf{S} - \nabla \tilde{\mathbf{d}}(t)) \otimes (\mathbf{h} - \tilde{\mathbf{d}}(t)) \rangle\rangle dt. \end{aligned}$$

Together with (4.4) we get

$$\langle\langle v_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 |\mathbf{h} - \tilde{\mathbf{d}}|^2 \rangle\rangle ds \leq c \langle\langle v_s, (\mathbf{S} - \nabla \tilde{\mathbf{d}}) \otimes (\mathbf{h} - \tilde{\mathbf{d}}) : \boldsymbol{\Theta} : (\mathbf{S} - \nabla \tilde{\mathbf{d}}) \otimes (\mathbf{h} - \tilde{\mathbf{d}}) \rangle\rangle$$

a. e. in $(0, T)$. A regrouping of the terms shows

$$\begin{aligned} (\mathbf{S} - \nabla \tilde{\mathbf{d}}) \otimes (\mathbf{h} - \tilde{\mathbf{d}}) &= \mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \mathbf{h} - \mathbf{S} \otimes \tilde{\mathbf{d}} + \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} \\ &= (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + (\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} - \nabla \tilde{\mathbf{d}} \otimes \mathbf{h}) + (\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} - \mathbf{S} \otimes \tilde{\mathbf{d}}) \\ &= (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} \otimes (\tilde{\mathbf{d}} - \mathbf{h}) + (\nabla \tilde{\mathbf{d}} - \mathbf{S}) \otimes \tilde{\mathbf{d}} \end{aligned}$$

and, thus

$$\begin{aligned} & \langle\langle v_s, (\mathbf{S} - \nabla \tilde{\mathbf{d}}) \otimes (\mathbf{h} - \tilde{\mathbf{d}}) : \Theta : (\mathbf{S} - \nabla \tilde{\mathbf{d}}) \otimes (\mathbf{h} - \tilde{\mathbf{d}}) \rangle\rangle \\ & \leq \langle\langle v_s, (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) : \Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \rangle\rangle + \langle\langle v_s, |\Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})|^2 \rangle\rangle \\ & \quad + c \|\nabla \tilde{\mathbf{d}}\|_{L^3}^2 \|\tilde{\mathbf{d}} - \mathbf{d}\|_{L^6}^2 + c \|\tilde{\mathbf{d}}\|_{L^\infty}^2 \langle\langle v_s, |\nabla \tilde{\mathbf{d}} - \mathbf{S}|^2 \rangle\rangle \end{aligned}$$

a. e. in $(0, T)$. The estimates (4.3a) and (4.3b) imply inequality (4.3c).

To prove the last inequality, we consider the left-hand side of (4.3c). Since m_t is a positive measure and v^o is a probability measure, we get with Jensen's inequality

$$\begin{aligned} \int_0^T \phi(t) \langle\langle v_t, |\mathbf{S} - \nabla \tilde{\mathbf{d}}(t)|^2 |\mathbf{h} - \tilde{\mathbf{d}}(t)|^2 \rangle\rangle dt & \geq \int_0^T \phi(t) \int_\Omega \int_{\mathbb{R}^{3 \times 3}} |\mathbf{S} - \nabla \tilde{\mathbf{d}}(\mathbf{x}, t)|^2 v_{(\mathbf{x}, t)}^o(d\mathbf{S}) |\mathbf{d}(\mathbf{x}, t) - \tilde{\mathbf{d}}(\mathbf{x}, t)|^2 d\mathbf{x} dt \\ & \geq \int_0^T \phi(t) \int_\Omega \left| \int_{\mathbb{R}^{3 \times 3}} (\mathbf{S} - \nabla \tilde{\mathbf{d}}(\mathbf{x}, t)) v_{(\mathbf{x}, t)}^o(d\mathbf{S}) \right|^2 |\mathbf{d}(\mathbf{x}, t) - \tilde{\mathbf{d}}(\mathbf{x}, t)|^2 d\mathbf{x} dt. \end{aligned}$$

The right-hand side can be identified due to (2.12) with the weak limit $\nabla \mathbf{d}$, such that with the embedding in three dimensions

$$\begin{aligned} \int_0^T \phi(t) \langle\langle v_t, |\mathbf{S} - \nabla \tilde{\mathbf{d}}(t)|^2 |\mathbf{h} - \tilde{\mathbf{d}}(t)|^2 \rangle\rangle dt & \geq \int_0^T \phi(t) \|\nabla \mathbf{d}(t) - \nabla \tilde{\mathbf{d}}(t)\| |\mathbf{d}(t) - \tilde{\mathbf{d}}(t)| \|_{L^2}^2 dt \\ & \geq \int_0^T \phi(t) \|\nabla |\mathbf{d}(t) - \tilde{\mathbf{d}}(t)|^2\|_{L^2}^2 dt \\ & \geq c \int_0^T \phi(t) \||\mathbf{d}(t) - \tilde{\mathbf{d}}(t)|^2\|_{L^6}^2 dt \\ & \geq c \int_0^T \phi(t) \||\mathbf{d}(t) - \tilde{\mathbf{d}}(t)|^4\|_{L^{12}}^4 dt. \end{aligned}$$

Remark that the difference of both solutions fulfills homogeneous Dirichlet boundary conditions. With the fundamental lemma of variational calculus the estimate (4.3d) is obvious.

□

Lemma 4.2. Let \mathbf{d} and $\tilde{\mathbf{d}}$ be given as in Lemma 4.1 as well as $\mathbf{a} \in L^1(0, T, \mathbf{L}^3)$. Then it holds

$$\int_0^t ([\mathbf{d}(s)]_X \partial_t \mathbf{d}(s) - [\tilde{\mathbf{d}}(s)]_X \partial_t \tilde{\mathbf{d}}(s), ([\mathbf{d}(s)]_X - [\tilde{\mathbf{d}}(s)]_X) \mathbf{a}(s)) ds \leq \int_0^t \mathcal{K}(s) \mathcal{E}(s) + \delta \mathcal{W}(s) ds.$$

Here \mathcal{K} is given by $\mathcal{K} = C_\delta \left(\|\tilde{\mathbf{v}}\|_{L^\infty}^2 + \|\tilde{\mathbf{v}}\|_{W^{1,3}}^2 + \|\mathbf{a}\|_{L^3}^2 \right) (\|\mathbf{d}\|_{L^\infty(L^\infty)} \|\tilde{\mathbf{d}}\|_{L^\infty(L^\infty)}^2 \|\tilde{\mathbf{d}}\|_{L^\infty(W^{1,3})}^2 + 1)$ and C_δ is a constant depending on δ .

Proof. First we insert equation (2.11b) and (2.1b) for the measure-valued solution and the strong solution respectively. This gives

$$\int_0^t ([\mathbf{d} - \tilde{\mathbf{d}}]_X^T ([\mathbf{d}]_X \partial_t \mathbf{d} - [\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}}), \mathbf{a}) ds = \int_0^t \left(([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X)^T ([\tilde{\mathbf{d}}]_X (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} - [\mathbf{d}]_X (\mathbf{v} \cdot \nabla) \mathbf{d}), \mathbf{a} \right) ds \quad (4.6a)$$

$$+ \int_0^t \left(([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X)^T ([\mathbf{d}]_X (\nabla \mathbf{v})_{\text{skw}} \mathbf{d} - [\tilde{\mathbf{d}}]_X (\nabla \tilde{\mathbf{v}})_{\text{skw}} \tilde{\mathbf{d}}), \mathbf{a} \right) ds \quad (4.6b)$$

$$+ \lambda \int_0^t \left(([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X)^T ([\tilde{\mathbf{d}}]_X (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}} - [\mathbf{d}]_X (\nabla \mathbf{v})_{\text{sym}} \mathbf{d}), \mathbf{a} \right) ds \quad (4.6c)$$

$$+ \int_0^t \left(([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X)^T ([\tilde{\mathbf{d}}]_X \tilde{\mathbf{q}} - [\mathbf{d}]_X \mathbf{q}), \mathbf{a} \right) ds. \quad (4.6d)$$

The dependence on s is not written out to remain the lucidity. We are going to estimate the lines individually. For the

line (4.6a) we observe

$$\begin{aligned}
& \int_0^t \left(([\mathbf{d}]_{\mathbf{X}} - [\tilde{\mathbf{d}}]_{\mathbf{X}})^T ([\tilde{\mathbf{d}}]_{\mathbf{X}}(\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} - [\mathbf{d}]_{\mathbf{X}}(\mathbf{v} \cdot \nabla) \mathbf{d}), \mathbf{a} \right) ds \\
&= \int_0^t \left(([\mathbf{d} - \tilde{\mathbf{d}}]_{\mathbf{X}})^T (([\tilde{\mathbf{d}} - \mathbf{d}]_{\mathbf{X}})(\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} + [\mathbf{d}]_{\mathbf{X}}(\tilde{\mathbf{v}} \cdot \nabla)(\tilde{\mathbf{d}} - \mathbf{d}) + [\mathbf{d}]_{\mathbf{X}}((\tilde{\mathbf{v}} - \mathbf{v}) \cdot \nabla) \mathbf{d}), \mathbf{a} \right) ds \\
&= \int_0^t \left(([\mathbf{d} - \tilde{\mathbf{d}}]_{\mathbf{X}})^T (([\tilde{\mathbf{d}} - \mathbf{d}]_{\mathbf{X}})(\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} + [\mathbf{d} - \tilde{\mathbf{d}}]_{\mathbf{X}}(\tilde{\mathbf{v}} \cdot \nabla)(\tilde{\mathbf{d}} - \mathbf{d}) + [\tilde{\mathbf{d}}]_{\mathbf{X}}(\tilde{\mathbf{v}} \cdot \nabla)(\tilde{\mathbf{d}} - \mathbf{d})), \mathbf{a} \right) ds \\
&\quad + \int_0^t \left(([\mathbf{d} - \tilde{\mathbf{d}}]_{\mathbf{X}})^T ([\mathbf{d}]_{\mathbf{X}}((\tilde{\mathbf{v}} - \mathbf{v}) \cdot \nabla)(\mathbf{d} - \tilde{\mathbf{d}})), \mathbf{a} \right) ds \\
&\quad + \int_0^t \left(([\mathbf{d} - \tilde{\mathbf{d}}]_{\mathbf{X}})^T ([\mathbf{d} - \tilde{\mathbf{d}}]_{\mathbf{X}}((\tilde{\mathbf{v}} - \mathbf{v}) \cdot \nabla) \tilde{\mathbf{d}} + [\tilde{\mathbf{d}}]_{\mathbf{X}}((\tilde{\mathbf{v}} - \mathbf{v}) \cdot \nabla) \tilde{\mathbf{d}}), \mathbf{a} \right) ds,
\end{aligned}$$

which can be estimated by

$$\begin{aligned}
& \int_0^t \left(([\mathbf{d}]_{\mathbf{X}} - [\tilde{\mathbf{d}}]_{\mathbf{X}})^T ([\tilde{\mathbf{d}}]_{\mathbf{X}}(\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} - [\mathbf{d}]_{\mathbf{X}}(\mathbf{v} \cdot \nabla) \mathbf{d}), \mathbf{a} \right) ds \\
&\leq c \int_0^t \|\tilde{\mathbf{v}}\|_{L^\infty} \|\nabla \tilde{\mathbf{d}}\|_{L^3} \|\mathbf{a}\|_{L^3} \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2 ds \\
&\quad + c \int_0^t \|\tilde{\mathbf{v}}\|_{L^\infty} \|\mathbf{a}\|_{L^3} (\|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^{12}}^4 + \|\tilde{\mathbf{d}}\|_{L^\infty} \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2 + \langle \langle v_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 \rangle \rangle) ds \\
&\quad + \delta \int_0^t \|\tilde{\mathbf{v}} - \mathbf{v}\|_{L^6}^2 ds + C_\delta + \int_0^t \|\mathbf{d}\|_{L^\infty}^2 \|\mathbf{a}\|_{L^3}^2 \langle \langle v_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 |\mathbf{h} - \tilde{\mathbf{d}}|^2 \rangle \rangle ds \\
&\quad + C_\delta \int_0^t \|\mathbf{d}\|_{W^{1,3}}^2 \|\mathbf{a}\|_{L^3}^2 (\|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^{12}}^4 + \|\tilde{\mathbf{d}}\|_{L^\infty}^2 \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2) ds.
\end{aligned}$$

For line (4.6b), we get a similar rearrangement

$$\begin{aligned}
& \int_0^t \left(([\mathbf{d}]_{\mathbf{X}} - [\tilde{\mathbf{d}}]_{\mathbf{X}})^T ([\mathbf{d}]_{\mathbf{X}}(\nabla \mathbf{v})_{\text{skw}} \mathbf{d} - [\tilde{\mathbf{d}}]_{\mathbf{X}}(\nabla \tilde{\mathbf{v}})_{\text{skw}} \tilde{\mathbf{d}}), \mathbf{a} \right) ds \\
&= \int_0^t \left(([\mathbf{d}]_{\mathbf{X}} - [\tilde{\mathbf{d}}]_{\mathbf{X}})^T ([\mathbf{d}]_{\mathbf{X}}((\nabla \mathbf{v})_{\text{skw}} - (\nabla \tilde{\mathbf{v}})_{\text{skw}}) \mathbf{d} + [\mathbf{d}]_{\mathbf{X}}(\nabla \tilde{\mathbf{v}})_{\text{skw}}(\mathbf{d} - \tilde{\mathbf{d}}) + ([\mathbf{d}]_{\mathbf{X}} - [\tilde{\mathbf{d}}]_{\mathbf{X}})(\nabla \tilde{\mathbf{v}})_{\text{skw}} \tilde{\mathbf{d}}), \mathbf{a} \right) ds \\
&= \int_0^t \left(([\mathbf{d}]_{\mathbf{X}} - [\tilde{\mathbf{d}}]_{\mathbf{X}})^T ([\mathbf{d}]_{\mathbf{X}}((\nabla \mathbf{v})_{\text{skw}} - (\nabla \tilde{\mathbf{v}})_{\text{skw}})(\mathbf{d} - \tilde{\mathbf{d}}) + [\mathbf{d}]_{\mathbf{X}}((\nabla \mathbf{v})_{\text{skw}} - (\nabla \tilde{\mathbf{v}})_{\text{skw}}) \tilde{\mathbf{d}}), \mathbf{a} \right) ds \\
&\quad + \int_0^t \left(([\mathbf{d}]_{\mathbf{X}} - [\tilde{\mathbf{d}}]_{\mathbf{X}})^T ([\mathbf{d}]_{\mathbf{X}}(\nabla \tilde{\mathbf{v}})_{\text{skw}}(\mathbf{d} - \tilde{\mathbf{d}}) + ([\mathbf{d}]_{\mathbf{X}} - [\tilde{\mathbf{d}}]_{\mathbf{X}})(\nabla \tilde{\mathbf{v}})_{\text{skw}} \tilde{\mathbf{d}}), \mathbf{a} \right) ds
\end{aligned}$$

such that the terms of line (4.6b) can be estimated by

$$\begin{aligned}
& \int_0^t \left(([\mathbf{d}]_{\mathbf{X}} - [\tilde{\mathbf{d}}]_{\mathbf{X}})^T ([\mathbf{d}]_{\mathbf{X}}(\nabla \mathbf{v})_{\text{skw}} \mathbf{d} - [\tilde{\mathbf{d}}]_{\mathbf{X}}(\nabla \tilde{\mathbf{v}})_{\text{skw}} \tilde{\mathbf{d}}), \mathbf{a} \right) ds \\
&\leq \delta \int_0^t \|(\nabla \mathbf{v})_{\text{skw}} - (\nabla \tilde{\mathbf{v}})_{\text{skw}}\|_{L^2}^2 ds + C_\delta \int_0^t \|\mathbf{d}\|_{L^\infty}^2 \|\mathbf{a}\|_{L^3}^2 (\|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^{12}}^4 + \|\tilde{\mathbf{d}}\|_{L^\infty}^2 \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2) ds \\
&\quad + c \int_0^t (\|\mathbf{d}\|_{L^\infty} + \|\tilde{\mathbf{d}}\|_{L^\infty}) \|(\nabla \tilde{\mathbf{v}})_{\text{skw}}\|_{L^3} \|\mathbf{a}\|_{L^3} \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2 ds.
\end{aligned}$$

For line (4.6c) we observe

$$\begin{aligned}
& \lambda \int_0^t \left(([\mathbf{d}]_{\mathbf{X}} - [\tilde{\mathbf{d}}]_{\mathbf{X}})^T ([\tilde{\mathbf{d}}]_{\mathbf{X}}(\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}} - [\mathbf{d}]_{\mathbf{X}}(\nabla \mathbf{v})_{\text{sym}} \mathbf{d}), \mathbf{a} \right) ds \\
&= \lambda \int_0^t \left(([\mathbf{d} - \tilde{\mathbf{d}}]_{\mathbf{X}})^T (([\tilde{\mathbf{d}} - \mathbf{d}]_{\mathbf{X}})(\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}} + [\mathbf{d}]_{\mathbf{X}}((\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}} - (\nabla \mathbf{v})_{\text{sym}} \mathbf{d})), \mathbf{a} \right) ds \\
&= \lambda \int_0^t \left(([\mathbf{d} - \tilde{\mathbf{d}}]_{\mathbf{X}})^T (((\tilde{\mathbf{d}} - \mathbf{d})_{\mathbf{X}})(\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}} + [\mathbf{d} - \tilde{\mathbf{d}}]_{\mathbf{X}}((\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}} - (\nabla \mathbf{v})_{\text{sym}} \mathbf{d})), \mathbf{a} \right) ds \\
&\quad + \lambda \int_0^t \left(([\mathbf{d} - \tilde{\mathbf{d}}]_{\mathbf{X}})^T [\tilde{\mathbf{d}}]_{\mathbf{X}}((\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}} - (\nabla \mathbf{v})_{\text{sym}} \mathbf{d}), \mathbf{a} \right) ds.
\end{aligned}$$

Thus, the terms of lines (4.6c) and (4.6d) can be estimated by

$$\begin{aligned} & \int_0^t \left(([\mathbf{d}]_{\mathbf{X}} - [\tilde{\mathbf{d}}]_{\mathbf{X}})^T (\lambda [\tilde{\mathbf{d}}]_{\mathbf{X}} (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}} - \lambda [\mathbf{d}]_{\mathbf{X}} (\nabla \mathbf{v})_{\text{sym}} \mathbf{d}) + [\tilde{\mathbf{d}}]_{\mathbf{X}} \tilde{\mathbf{q}} - [\mathbf{d}]_{\mathbf{X}} \mathbf{q}, \mathbf{a} \right) \mathrm{d}s \\ & \leq c \int_0^t \|\tilde{\mathbf{d}}\|_{\mathbf{L}^\infty} \|(\nabla \tilde{\mathbf{v}})_{\text{sym}}\|_{\mathbf{L}^3} \|\mathbf{a}\|_{\mathbf{L}^3} \|\mathbf{d} - \tilde{\mathbf{d}}\|_{\mathbf{L}^6}^2 \mathrm{d}s \\ & + \delta \int_0^t \|(\nabla \mathbf{v})_{\text{sym}} \mathbf{d} - (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}\|_{\mathbf{L}^2}^2 \mathrm{d}s + C_\delta \int_0^t \|\mathbf{a}\|_{\mathbf{L}^3}^2 (\|\mathbf{d} - \tilde{\mathbf{d}}\|_{\mathbf{L}^{12}}^4 + \|\tilde{\mathbf{d}}\|_{\mathbf{L}^\infty} \|\mathbf{d} - \tilde{\mathbf{d}}\|_{\mathbf{L}^6}^2) \mathrm{d}s \\ & + \delta \int_0^t \|\mathbf{d} \times \mathbf{q} - \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}\|_{\mathbf{L}^2}^2 \mathrm{d}s + C_\delta \int_0^t \|\mathbf{a}\|_{\mathbf{L}^3}^2 \|\mathbf{d} - \tilde{\mathbf{d}}\|_{\mathbf{L}^6}^2 \mathrm{d}s \end{aligned}$$

The above estimates together with Lemma 3.4 and the definitions (4.1) and (4.2) prove the assertion. \square

5 Integration-by-parts formulae

Proposition 5.1 (Integration-by-parts formula). Let (\mathbf{v}, \mathbf{d}) be a measure-valued solution and $(\tilde{\mathbf{v}}, \tilde{\mathbf{d}})$ a strong solution. Then the following integration-by-parts formulae hold for a.e. $t, s \in (0, T)$.

$$(\mathbf{v}(t), \tilde{\mathbf{v}}(t)) - (\mathbf{v}(s), \tilde{\mathbf{v}}(s)) = \int_s^t (\mathbf{v}(\tau), \partial_t \tilde{\mathbf{v}}(\tau)) + (\partial_t \mathbf{v}(\tau), \tilde{\mathbf{v}}(\tau)) \mathrm{d}\tau, \quad (5.1)$$

$$\|\mathbf{d}(t) - \tilde{\mathbf{d}}(t)\|_{\mathbf{L}^2}^2 - \|\mathbf{d}(s) - \tilde{\mathbf{d}}(s)\|_{\mathbf{L}^2}^2 \leq \delta \int_s^t \mathcal{W}(\tau) \mathrm{d}\tau + C_\delta \int_s^t \mathcal{E}(\tau) \mathrm{d}\tau, \quad (5.2)$$

$$\begin{aligned} & (\nabla \mathbf{d}(t); \boldsymbol{\Lambda} : \nabla \tilde{\mathbf{d}}(t)) - (\nabla \mathbf{d}(s); \boldsymbol{\Lambda} : \nabla \tilde{\mathbf{d}}(s)) \\ & \leq \int_s^t [\langle\langle \mathbf{v}_t, \mathbf{S} : \boldsymbol{\Lambda} : (\mathbf{T} : ([\tilde{\mathbf{d}}]_{\mathbf{X}} \partial_t \tilde{\mathbf{d}} \otimes \mathbf{S})) \rangle\rangle + (\nabla \mathbf{d}; \boldsymbol{\Lambda} : [\mathbf{d}]_{\mathbf{X}}^T \nabla ([\tilde{\mathbf{d}}]_{\mathbf{X}} \partial_t \tilde{\mathbf{d}}))] \mathrm{d}\tau \\ & - \int_s^t ([\tilde{\mathbf{d}}]_{\mathbf{X}} \Delta_{\boldsymbol{\Lambda}} \tilde{\mathbf{d}}, [\mathbf{d}]_{\mathbf{X}} \partial_t \mathbf{d}) \mathrm{d}\tau + \delta \int_s^t \mathcal{W}(\tau) \mathrm{d}\tau + c \int_s^t \mathcal{E}(\tau) \mathrm{d}\tau, \end{aligned} \quad (5.3)$$

and

$$\begin{aligned} & (\nabla \mathbf{d}(t) \otimes \mathbf{d}(t); \boldsymbol{\Theta} : \nabla \tilde{\mathbf{d}}(t) \otimes \tilde{\mathbf{d}}(t)) - (\nabla \mathbf{d}(s) \otimes \mathbf{d}(s); \boldsymbol{\Theta} : \nabla \tilde{\mathbf{d}}(s) \otimes \tilde{\mathbf{d}}(s)) \\ & - ((\nabla \mathbf{d}(t) - \nabla \tilde{\mathbf{d}}(t)) \otimes (\mathbf{d}(t) - \tilde{\mathbf{d}}(t))) : \boldsymbol{\Theta} : (\nabla \tilde{\mathbf{d}}(t) \otimes \tilde{\mathbf{d}}(t)) \\ & + ((\nabla \mathbf{d}(s) - \nabla \tilde{\mathbf{d}}(s)) \otimes (\mathbf{d}(s) - \tilde{\mathbf{d}}(s))) : \boldsymbol{\Theta} : (\nabla \tilde{\mathbf{d}}(s) \otimes \tilde{\mathbf{d}}(s)) \end{aligned} \quad (5.4a)$$

$$\leq \int_s^t ([\mathbf{d}]_{\mathbf{X}} \partial_t \mathbf{d}, [\tilde{\mathbf{d}}]_{\mathbf{X}} (-\nabla \cdot (\tilde{\mathbf{d}} : \boldsymbol{\Theta} : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \boldsymbol{\Theta} : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) \mathrm{d}\tau \quad (5.4b)$$

$$+ \int_s^t \langle\langle \mathbf{v}_\tau, \mathbf{S} \otimes \mathbf{h} : \boldsymbol{\Theta} : (\mathbf{T} : ([\tilde{\mathbf{d}}]_{\mathbf{X}}^T \partial_t \tilde{\mathbf{d}} \otimes \mathbf{S})) \otimes \mathbf{h} \rangle\rangle \mathrm{d}\tau \quad (5.4c)$$

$$+ \int_s^t (\nabla \mathbf{d} \otimes \mathbf{d}; \boldsymbol{\Theta} : [\mathbf{d}]_{\mathbf{X}}^T \nabla ([\tilde{\mathbf{d}}]_{\mathbf{X}} \partial_t \tilde{\mathbf{d}}) \otimes \mathbf{d}) \mathrm{d}\tau \quad (5.4d)$$

$$+ \int_s^t \langle\langle \mathbf{v}_\tau, \mathbf{S} \otimes \mathbf{h} : \boldsymbol{\Theta} : \mathbf{S} \otimes [\mathbf{h}]_{\mathbf{X}}^T [\tilde{\mathbf{d}}]_{\mathbf{X}} \partial_t \tilde{\mathbf{d}} \rangle\rangle \mathrm{d}\tau + \delta \int_s^t \mathcal{W} \mathrm{d}\tau + c \int_s^t \mathcal{E} \mathrm{d}\tau. \quad (5.4e)$$

Proof. Let $\phi \in \mathcal{C}_c^\infty(0, T)$ we calculate the weak time derivative and neglect the dependence on the time variable t under the integral for more convenient writing.

Remark that \mathbf{v} and $\tilde{\mathbf{v}}$ fulfill the regularity assumptions

$$\mathbf{v} \in L^2(0, T; \mathbf{H}_{0,\sigma}^1) \cap W^{1,2}(0, T; (\mathbf{W}_{0,\sigma}^{1,3})^*) \quad \text{and} \quad \tilde{\mathbf{v}} \in L^2(0, T; (\mathbf{W}_{0,\sigma}^{1,3})^*) \cap W^{1,2}(0, T; (\mathbf{H}_{0,\sigma}^1)^*).$$

We consider the weak derivative of the L^2 -product of \mathbf{v} and $\tilde{\mathbf{v}}$

$$\int_0^T \phi'(t) (\mathbf{v}(t), \tilde{\mathbf{v}}(t)) \mathrm{d}t = \int_0^T \phi(t) ((\partial_t \mathbf{v}(t), \tilde{\mathbf{v}}(t)) + (\mathbf{v}(t), \partial_t \tilde{\mathbf{v}}(t))) \mathrm{d}t.$$

The variational lemma implies the assertion (5.1).

To prove the formula (5.2) one can simply calculate the weak derivative, since the regularity of \mathbf{d} is sufficient to do so:

$$\int_0^T \phi' \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^2}^2 dt = 2 \int_0^T \phi (\partial_t \mathbf{d} - \partial_t \tilde{\mathbf{d}}, \mathbf{d} - \tilde{\mathbf{d}}) dt$$

Applying Lemma 3.2 and regrouping the terms shows

$$\begin{aligned} \int_0^T \phi' \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^2}^2 dt &= 2 \int_0^T \phi ([\mathbf{d}]_X^T [\mathbf{d}]_X \partial_t \mathbf{d} - [\tilde{\mathbf{d}}]_X^T [\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}}, \mathbf{d} - \tilde{\mathbf{d}}) dt \\ &= 2 \int_0^T \phi ([\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}}, ([\tilde{\mathbf{d}}]_X - [\mathbf{d}]_X) (\mathbf{d} - \tilde{\mathbf{d}})) dt \\ &\quad + 2 \int_0^T \phi ([\mathbf{d}]_X \partial_t \mathbf{d} - [\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}}, ([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X) \tilde{\mathbf{d}}) dt \end{aligned}$$

The right hand side can be estimated by Lemma 4.2 and by

$$\int_0^T \phi ([\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}}, ([\tilde{\mathbf{d}}]_X - [\mathbf{d}]_X) (\mathbf{d} - \tilde{\mathbf{d}})) dt \leq \|\tilde{\mathbf{d}}\|_{L^\infty(L^\infty)} \|\partial_t \tilde{\mathbf{d}}\|_{L^\infty(L^\infty)} \int_0^T \phi \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^2}^2 dt.$$

Lemma 3.1 implies the formula (5.2).

For the next integration-by-parts formula (5.3), the weak time derivative of the quadratic term of the gradient of the director is calculated. Remark that there is no additional regularity known for the time derivative of $\nabla \partial_t \mathbf{d}$. We write

$$-\int_0^T \phi' (\nabla \mathbf{d}; \boldsymbol{\Lambda} : \nabla \tilde{\mathbf{d}}) dt = \int_0^T [\phi (\nabla \mathbf{d}; \boldsymbol{\Lambda} : \nabla \partial_t \tilde{\mathbf{d}}) - (\nabla \mathbf{d}; \boldsymbol{\Lambda} : \partial_t (\phi \nabla \tilde{\mathbf{d}}))] dt$$

The measure-valued solution is approximated by smooth functions $\{\mathbf{d}_n\}$ fulfilling the same boundary conditions, such that $\mathbf{d}_n \rightarrow \mathbf{d}$ in $L^2(0, T; \mathbf{H}^1) \cap W^{1,2}(0, T; \mathbf{L}^{3/2})$. For this approximation, the following integration-by-parts holds true

$$\begin{aligned} -\int_0^T [\nabla \mathbf{d}_n; \partial_t (\phi \boldsymbol{\Lambda} : \nabla \tilde{\mathbf{d}})] dt &= \int_0^T \phi (\nabla \partial_t \mathbf{d}_n; \boldsymbol{\Lambda} : \nabla \tilde{\mathbf{d}}) dt \\ &= \int_0^T \phi (\partial_t \mathbf{d}_n, -\Delta_{\boldsymbol{\Lambda}} \tilde{\mathbf{d}}) dt. \end{aligned}$$

The boundary values vanish since $\phi \in \mathcal{C}_c^\infty(0, T)$ and since the boundary values of \mathbf{d}_n are constant in time and, hence, its time derivative vanishes on the boundary. Going to the limit in the approximation gives

$$-\int_0^T (\nabla \mathbf{d}; \partial_t (\phi \boldsymbol{\Lambda} : \nabla \tilde{\mathbf{d}})) dt = \int_0^T \phi (\partial_t \mathbf{d}, -\Delta_{\boldsymbol{\Lambda}} \tilde{\mathbf{d}}) dt.$$

Lemma 3.2 implies that

$$\begin{aligned} -\int_0^T \phi' (\nabla \mathbf{d}; \boldsymbol{\Lambda} : \nabla \tilde{\mathbf{d}}) dt &= \int_0^T \phi ((\nabla \mathbf{d}; \boldsymbol{\Lambda} : \nabla ([\tilde{\mathbf{d}}]_X^T [\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}})) + ([\mathbf{d}]_X^T [\mathbf{d}]_X \partial_t \mathbf{d}, -\Delta_{\boldsymbol{\Lambda}} \tilde{\mathbf{d}})) dt \\ &= \int_0^T \phi ((\nabla \mathbf{d}; \boldsymbol{\Lambda} : (\boldsymbol{\Upsilon} : ([\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}}) \otimes \nabla \tilde{\mathbf{d}} + [\tilde{\mathbf{d}}]_X^T \nabla ([\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}}))) + ([\mathbf{d}]_X^T [\mathbf{d}]_X \partial_t \mathbf{d}, -\Delta_{\boldsymbol{\Lambda}} \tilde{\mathbf{d}})) dt \end{aligned} \tag{5.5}$$

To show the formula (5.3), we add and subtract the desired terms of the right hand side of (5.3). Some additional rearrangement of the terms show

$$\begin{aligned} -\int_0^T \phi' (\nabla \mathbf{d}; \boldsymbol{\Lambda} : \nabla \tilde{\mathbf{d}}) dt &= \int_0^T \phi (\langle \langle \mathbf{v}_t, \mathbf{S} : \boldsymbol{\Lambda} : \boldsymbol{\Upsilon} : ([\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}} \otimes \mathbf{S}) \rangle \rangle + (\nabla \mathbf{d}; \boldsymbol{\Lambda} : [\mathbf{d}]_X^T \nabla ([\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}})) - ([\tilde{\mathbf{d}}]_X \Delta_{\boldsymbol{\Lambda}} \tilde{\mathbf{d}}, [\mathbf{d}]_X \partial_t \mathbf{d})) dt \\ &\quad + \int_0^T \phi (\langle \langle \mathbf{v}_t, \mathbf{S} : \boldsymbol{\Lambda} : \boldsymbol{\Upsilon} : [\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}} \otimes (\nabla \tilde{\mathbf{d}} - \mathbf{S}) \rangle \rangle + (\nabla \mathbf{d}; \boldsymbol{\Lambda} : ([\tilde{\mathbf{d}}]_X - [\mathbf{d}]_X) \nabla ([\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}}))) dt \\ &\quad + \int_0^T \phi (([\tilde{\mathbf{d}}]_X - [\mathbf{d}]_X)^T [\mathbf{d}]_X \partial_t \mathbf{d}, \Delta_{\boldsymbol{\Lambda}} \tilde{\mathbf{d}}) dt \end{aligned} \tag{5.6}$$

Adding and subtracting the term

$$\begin{aligned} \int_0^T \phi'(([d]_X - [d]_X)^T [\tilde{d}]_X \partial_t \tilde{d}, -\Delta_\Lambda \tilde{d}) dt &= \int_0^T \phi(\nabla \tilde{d} : \boldsymbol{\Lambda} : \boldsymbol{\Upsilon} : ([d]_X \partial_t \tilde{d} \otimes (\nabla \tilde{d} - \nabla d))) dt \\ &\quad + \int_0^T \phi(\nabla \tilde{d} : \boldsymbol{\Lambda} : (([d]_X - [d]_X) \nabla ([d]_X \partial_t \tilde{d}))) dt \end{aligned}$$

gives

$$\begin{aligned} &- \int_0^T \phi'(\nabla d : \boldsymbol{\Lambda} : \nabla \tilde{d}) dt \\ &= \int_0^T \phi(\langle v_t, \mathbf{S} : \boldsymbol{\Lambda} : \boldsymbol{\Upsilon} : ([\tilde{d}]_X \partial_t \tilde{d} \otimes \mathbf{S}) \rangle + (\nabla d : \boldsymbol{\Lambda} : [\tilde{d}]_X^T \nabla ([\tilde{d}]_X \partial_t \tilde{d})) - ([\tilde{d}]_X \Delta_\Lambda \tilde{d}, [d]_X \partial_t d)) dt \end{aligned} \quad (5.7a)$$

$$\begin{aligned} &+ \int_0^T \phi(\langle v_t, (\mathbf{S} - \nabla \tilde{d}) : \boldsymbol{\Lambda} : \boldsymbol{\Upsilon} : [\tilde{d}]_X \partial_t \tilde{d} \otimes (\nabla \tilde{d} - \mathbf{S}) \rangle + (\nabla d - \nabla \tilde{d} : \boldsymbol{\Lambda} : (([\tilde{d}]_X - [d]_X) \nabla ([\tilde{d}]_X \partial_t \tilde{d}))) dt \\ &+ \int_0^T \phi(([d]_X - [d]_X)^T ([\tilde{d}]_X \partial_t \tilde{d} - [d]_X \partial_t d), -\Delta_\Lambda \tilde{d}) dt \end{aligned} \quad (5.7b)$$

We can estimate the terms of line (5.7a) by

$$\begin{aligned} &\int_0^T \phi(\langle v_t, (\mathbf{S} - \nabla \tilde{d}) : \boldsymbol{\Lambda} : \boldsymbol{\Upsilon} : [\tilde{d}]_X \partial_t \tilde{d} \otimes (\nabla \tilde{d} - \mathbf{S}) \rangle + (\nabla d - \nabla \tilde{d} : \boldsymbol{\Lambda} : (([\tilde{d}]_X - [d]_X) \nabla ([\tilde{d}]_X \partial_t \tilde{d}))) dt \\ &\leq c \|\tilde{d}\|_{L^\infty(\mathbf{L}^\infty)} \int_0^T \phi \|\partial_t \tilde{d}\|_{\mathbf{L}^\infty} \langle v_t, |\mathbf{S} - \nabla \tilde{d}|^2 \rangle dt \\ &\quad + c \int_0^T \phi(\|\nabla \tilde{d}\|_{L^\infty(\mathbf{L}^3)} \|\partial_t \tilde{d}\|_{\mathbf{L}^\infty} + \|\tilde{d}\|_{L^\infty(\mathbf{L}^\infty)} \|\nabla \partial_t \tilde{d}\|_{\mathbf{L}^3}) (\langle v_t, |\mathbf{S} - \nabla \tilde{d}|^2 \rangle + \|d - \tilde{d}\|_{\mathbf{L}^6}^2) dt \end{aligned}$$

and terms of line (5.7b) by Lemma 4.2. Inserting those estimates back into (5.7) and applying Lemma 3.1 implies the integration-by-parts formula (5.3).

Finally, we are going to prove the integration-by-parts formula (5.4). Consider the following term

$$\begin{aligned} &- \int_0^T \phi'(\nabla d \otimes d ; \boldsymbol{\Theta} : \nabla \tilde{d} \otimes \tilde{d}) + \int_0^T ((\nabla d - \nabla \tilde{d}) \otimes (d - \tilde{d}) ; \boldsymbol{\Theta} : \partial_t (\phi \nabla \tilde{d} \otimes \tilde{d})) dt \\ &= - \int_0^T (\nabla d \otimes d ; \boldsymbol{\Theta} : \partial_t (\phi \nabla \tilde{d} \otimes \tilde{d})) dt + \int_0^T \phi(\nabla d \otimes d ; \boldsymbol{\Theta} : \partial_t (\nabla \tilde{d} \otimes \tilde{d})) dt \\ &\quad + \int_0^T (\nabla d \otimes d ; \boldsymbol{\Theta} : \partial_t (\phi \nabla \tilde{d} \otimes \tilde{d})) dt - \int_0^T (\nabla \tilde{d} \otimes d ; \boldsymbol{\Theta} : \partial_t (\phi \nabla \tilde{d} \otimes \tilde{d})) dt \\ &\quad - \int_0^T (\nabla d \otimes \tilde{d} ; \boldsymbol{\Theta} : \partial_t (\phi \nabla \tilde{d} \otimes \tilde{d})) dt + \int_0^T (\nabla \tilde{d} \otimes \tilde{d} ; \boldsymbol{\Theta} : \partial_t (\phi \nabla \tilde{d} \otimes \tilde{d})) dt \end{aligned}$$

The first and the third term on the right-hand cancel each other. Calculating the weak time derivatives of the products gives

$$\begin{aligned} &- \int_0^T \phi'(\nabla d \otimes d ; \boldsymbol{\Theta} : \nabla \tilde{d} \otimes \tilde{d}) dt + \int_0^T ((\nabla d - \nabla \tilde{d}) \otimes (d - \tilde{d}) ; \boldsymbol{\Theta} : \partial_t (\phi \nabla \tilde{d} \otimes \tilde{d})) dt \\ &= \int_0^T \phi(\nabla d \otimes d ; \boldsymbol{\Theta} : \nabla \partial_t \tilde{d} \otimes \tilde{d}) dt + \int_0^T \phi(\nabla d \otimes d ; \boldsymbol{\Theta} : \nabla \tilde{d} \otimes \partial_t \tilde{d}) dt \\ &\quad + \int_0^T \phi(\nabla \tilde{d} \otimes \partial_t d ; \boldsymbol{\Theta} : \nabla \tilde{d} \otimes \tilde{d}) dt + \int_0^T \phi(\nabla \partial_t \tilde{d} \otimes d ; \boldsymbol{\Theta} : \nabla \tilde{d} \otimes \tilde{d}) dt \\ &\quad - \int_0^T (\nabla d ; \partial_t (\phi \tilde{d} \cdot \boldsymbol{\Theta} : \nabla \tilde{d} \otimes \tilde{d})) dt + \int_0^T \phi(\nabla d \otimes \partial_t \tilde{d} ; \boldsymbol{\Theta} : \nabla \tilde{d} \otimes \tilde{d}) dt \\ &\quad - \int_0^T \phi(\nabla \partial_t \tilde{d} \otimes \tilde{d} ; \boldsymbol{\Theta} : \nabla \tilde{d} \otimes \tilde{d}) dt - \int_0^T \phi(\nabla \tilde{d} \otimes \partial_t \tilde{d} ; \boldsymbol{\Theta} : \nabla \tilde{d} \otimes \tilde{d}) dt \end{aligned} \quad (5.8)$$

In the the first term of the line (5.8) the direktor of the measure-valued solution is again approximated by a sequence $\{\mathbf{d}_n\} \subset \mathcal{C}^1(\Omega \times (0, T))$ fulfilling the same boundary values, such that $\mathbf{d}_n \rightarrow \mathbf{d}$ in $L^2(0, T; \mathbf{H}^1) \cap W^{1,2}(0, T; \mathbf{L}^{3/2})$. For this approximations, the following integration-by-parts is valid

$$\begin{aligned} &- \int_0^T (\nabla \mathbf{d}_n ; \partial_t (\phi \tilde{d} \cdot \boldsymbol{\Theta} : (\nabla \tilde{d} \otimes \tilde{d}))) dt = \int_0^T (\nabla \partial_t \mathbf{d}_n, \phi \tilde{d} \cdot \boldsymbol{\Theta} : (\nabla \tilde{d} \otimes \tilde{d})) dt \\ &= \int_0^T \phi(\partial_t \mathbf{d}_n, -\nabla \cdot (\tilde{d} \cdot \boldsymbol{\Theta} : (\nabla \tilde{d} \otimes \tilde{d}))) dt. \end{aligned} \quad (5.9)$$

The boundary terms vanish, since $\phi \in \mathcal{C}_c^\infty(0, T)$ and since the prescribed boundary value for \mathbf{d}_n is constant in time, so that is time derivative vanishes. Going to the limit in the approximation gives

$$-\int_0^T (\nabla \mathbf{d}; \partial_t (\phi \tilde{\mathbf{d}} \cdot \Theta : (\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}))) dt = \int_0^T \phi (\partial_t \mathbf{d}, -\nabla \cdot (\tilde{\mathbf{d}} \cdot \Theta : (\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}))) dt.$$

Adding and subtracting the terms

$$\int_0^T \phi (\nabla \mathbf{d} \otimes \mathbf{d}; \Theta : \nabla \partial_t \tilde{\mathbf{d}} \otimes \mathbf{d}) dt + \int_0^T \phi \langle v_t, \mathbf{S} \otimes \mathbf{h} : \Theta : \mathbf{S} \otimes \partial_t \tilde{\mathbf{d}} \rangle dt$$

to (5.8) shows

$$\begin{aligned} & -\int_0^T \phi' (\nabla \mathbf{d} \otimes \mathbf{d}; \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) dt + \int_0^T \phi' ((\nabla \mathbf{d} - \nabla \tilde{\mathbf{d}}) \otimes (\mathbf{d} - \tilde{\mathbf{d}}); \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) dt \\ & + \int_0^T \phi ((\nabla \mathbf{d} - \nabla \tilde{\mathbf{d}}) \otimes (\mathbf{d} - \tilde{\mathbf{d}}); \Theta : \partial_t (\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) dt \end{aligned} \quad (5.10a)$$

$$= \int_0^T \phi (\partial_t \mathbf{d}, -\nabla \cdot (\tilde{\mathbf{d}} \cdot \Theta : (\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) dt \quad (5.10b)$$

$$+ \int_0^T \phi (\nabla \partial_t \tilde{\mathbf{d}} \otimes \mathbf{d}; \Theta : \nabla \mathbf{d} \otimes \mathbf{d}) dt + \int_0^T \phi \langle v_t, \mathbf{S} \otimes \mathbf{h} : \Theta : \mathbf{S} \otimes \partial_t \tilde{\mathbf{d}} \rangle dt \quad (5.10c)$$

$$+ \int_0^T \phi ((\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} - \nabla \mathbf{d} \otimes \mathbf{d}); \Theta : \nabla \partial_t \tilde{\mathbf{d}} \otimes (\mathbf{d} - \tilde{\mathbf{d}})) dt \quad (5.10d)$$

$$+ \int_0^T \phi \langle v_t, (\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} - \mathbf{S} \otimes \mathbf{h}) : \Theta : (\mathbf{S} - \nabla \tilde{\mathbf{d}}) \otimes \partial_t \tilde{\mathbf{d}} \rangle dt \quad (5.10e)$$

We observe that the last two terms of the lines (5.10d) and (5.10e) can be estimated by

$$\begin{aligned} & + \int_0^T \phi ((\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} - \nabla \mathbf{d} \otimes \mathbf{d}); \Theta : \nabla \partial_t \tilde{\mathbf{d}} \otimes (\mathbf{d} - \tilde{\mathbf{d}})) dt \\ & + \int_0^T \phi \langle v_t, (\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} - \mathbf{S} \otimes \mathbf{h}) : \Theta : (\mathbf{S} - \nabla \tilde{\mathbf{d}}) \otimes \partial_t \tilde{\mathbf{d}} \rangle dt \\ & \leq \int_0^T \phi \|\nabla \partial_t \tilde{\mathbf{d}}\|_{L^3} \left(\langle v_t, |\Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})|^2 \rangle + \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2 \right) dt \\ & + \int_0^T \phi \|\partial_t \tilde{\mathbf{d}}\|_{L^\infty} \left(\langle v_t, |\Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})|^2 \rangle + \langle v_t, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 \rangle \right) dt. \end{aligned}$$

The line (5.10a) can be estimated by

$$\begin{aligned} & \int_0^T \phi ((\nabla \mathbf{d} - \nabla \tilde{\mathbf{d}}) \otimes (\mathbf{d} - \tilde{\mathbf{d}}); \Theta : \partial_t (\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) dt \\ & \leq \|\tilde{\mathbf{d}}\|_{L^\infty(L^\infty)} \int_0^T \left[\phi \|\nabla \partial_t \tilde{\mathbf{d}}\|_{L^3} \langle v_t, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 \rangle + \phi \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2 \right] dt \\ & \quad \|\nabla \tilde{\mathbf{d}}\|_{L^\infty(L^3)} \int_0^T \left[\phi \|\partial_t \tilde{\mathbf{d}}\|_{L^\infty} \langle v_t, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 \rangle + \phi \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2 \right] dt. \end{aligned}$$

The remaining terms of the right-hand side of (5.10), in line (5.10b) and in line (5.10c) are considered further on. Due to Lemma 4.2 we see

$$\begin{aligned} & \int_0^T \phi (\partial_t \mathbf{d}, -\nabla \cdot (\tilde{\mathbf{d}} \cdot \Theta : (\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) dt \\ & + \int_0^T \phi (\nabla \partial_t \tilde{\mathbf{d}} \otimes \mathbf{d}; \Theta : \nabla \mathbf{d} \otimes \mathbf{d}) dt + \int_0^T \phi \langle v_t, \mathbf{S} \otimes \mathbf{h} : \Theta : \mathbf{S} \otimes \partial_t \tilde{\mathbf{d}} \rangle dt \\ & = \int_0^T \phi ([\mathbf{d}]_X^T [\mathbf{d}] \partial_t \mathbf{d}, -\nabla \cdot (\tilde{\mathbf{d}} \cdot \Theta : (\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) dt \end{aligned} \quad (5.11a)$$

$$+ \int_0^T \phi (\nabla ([\tilde{\mathbf{d}}]_X^T [\tilde{\mathbf{d}}] \partial_t \tilde{\mathbf{d}}) \otimes \mathbf{d}; \Theta : \nabla \mathbf{d} \otimes \mathbf{d}) dt + \int_0^T \phi \langle v_t, \mathbf{S} \otimes \mathbf{h} : \Theta : \mathbf{S} \otimes [\tilde{\mathbf{d}}]_X^T [\tilde{\mathbf{d}}] \partial_t \tilde{\mathbf{d}} \rangle dt. \quad (5.11b)$$

Considering now the difference of the right-hand side of the above equality and the desired terms due to the integration-by-parts formula (5.4), we get for the terms of the lines (5.4b) and (5.11a)

$$\begin{aligned}
& \int_0^T \phi ([\mathbf{d}]_X^T [\mathbf{d}]_X \partial_t \mathbf{d}, -\nabla \cdot (\tilde{\mathbf{d}} \cdot \Theta : (\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) dt \\
& - \int_0^T \phi ([\mathbf{d}]_X \partial_t \mathbf{d}, [\tilde{\mathbf{d}}]_X (-\nabla \cdot (\tilde{\mathbf{d}} \cdot \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) dt \\
& = \int_0^T \phi ([\mathbf{d}]_X \partial_t \mathbf{d} - [\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}}, ([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X) (-\nabla \cdot (\tilde{\mathbf{d}} \cdot \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) dt \\
& + \int_0^T \phi ([\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}}, ([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X) (-\nabla \cdot (\tilde{\mathbf{d}} \cdot \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) dt,
\end{aligned} \tag{5.12}$$

as well as for the terms of the lines (5.4d), (5.4e) and the first term in line (5.11b)

$$\begin{aligned}
& \int_0^T \phi (\nabla ([\tilde{\mathbf{d}}]_X^T [\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}}) \otimes \mathbf{d}; \Theta : \nabla \mathbf{d} \otimes \mathbf{d}) dt \\
& - \int_0^T \phi \langle \langle v_\tau, \mathbf{S} \otimes \mathbf{h} : \Theta : ([\tilde{\mathbf{d}}]_X^T \partial_t \tilde{\mathbf{d}} \otimes \mathbf{S}) \otimes \mathbf{h} \rangle \rangle dt - \int_0^T \phi (\nabla \mathbf{d} \otimes \mathbf{d} ; \Theta : [\mathbf{d}]_X^T \nabla ([\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}}) \otimes \mathbf{d}) dt \\
& = \int_0^T \phi \langle \langle v_\tau, \mathbf{S} \otimes \mathbf{h} : \Theta : (\Upsilon : ([\tilde{\mathbf{d}}]_X^T \partial_t \tilde{\mathbf{d}} \otimes (\nabla \tilde{\mathbf{d}} - \mathbf{S}))) \otimes \mathbf{h} \rangle \rangle dt \\
& + \int_0^T \phi (\nabla \mathbf{d} \otimes \mathbf{d} ; \Theta : ([\tilde{\mathbf{d}}]_X - [\mathbf{d}]_X)^T \nabla ([\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}}) \otimes \mathbf{d}) dt.
\end{aligned} \tag{5.13}$$

Simimilar, we get for the terms in line (5.4c) subtracted from the second term in line (5.11b) that

$$\begin{aligned}
& \int_0^T \phi \langle \langle v_\tau, \mathbf{S} \otimes \mathbf{h} : \Theta : \mathbf{S} \otimes [\tilde{\mathbf{d}}]_X^T [\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}} \rangle \rangle dt - \int_0^T \phi \langle \langle v_\tau, \mathbf{S} \otimes \mathbf{h} : \Theta : \mathbf{S} \otimes [\mathbf{h}]_X^T [\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}} \rangle \rangle dt \\
& = \int_0^T \phi \langle \langle v_\tau, \mathbf{S} \otimes \mathbf{h} : \Theta : \mathbf{S} \otimes ([\tilde{\mathbf{d}}]_X - [\mathbf{h}]_X)^T [\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}} \rangle \rangle dt.
\end{aligned} \tag{5.14}$$

We recognise that the first term on the right hand side of (5.12) can be estimated by Lemma 4.2. The last term of equation (5.12) can, via an integration-by-parts, be expressed as

$$\begin{aligned}
& \int_0^T \phi ([\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}}, ([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X) (-\nabla \cdot (\tilde{\mathbf{d}} \cdot \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) dt \\
& = \int_0^T \phi \left(\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} ; \Theta : \left(\Upsilon : ([\tilde{\mathbf{d}}]_X^T \partial_t \tilde{\mathbf{d}} \otimes (\nabla \mathbf{d} - \nabla \tilde{\mathbf{d}})) \otimes \tilde{\mathbf{d}} + ([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X)^T \nabla ([\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}}) \otimes \tilde{\mathbf{d}} \right) \right) dt
\end{aligned} \tag{5.15a}$$

$$+ \int_0^T \phi \left(\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} ; \Theta : \nabla \tilde{\mathbf{d}} \otimes ([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X)^T [\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}} \right) dt. \tag{5.15b}$$

First we consider the difference of the right-hand side of (5.13) and the terms of line (5.15a):

$$\begin{aligned}
& \int_0^T \phi \langle \langle v_\tau, \mathbf{S} \otimes \mathbf{h} : \Theta : \Upsilon : ([\tilde{\mathbf{d}}]_X^T \partial_t \tilde{\mathbf{d}} \otimes (\nabla \tilde{\mathbf{d}} - \mathbf{S})) \otimes \mathbf{h} \rangle \rangle dt \\
& + \int_0^T \phi \left(\nabla \mathbf{d} \otimes \mathbf{d} ; \Theta : ([\tilde{\mathbf{d}}]_X - [\mathbf{d}]_X)^T \nabla ([\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}}) \otimes \mathbf{d} \right) dt \\
& + \int_0^T \phi \left(\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} ; \Theta : \left((\Upsilon : ([\tilde{\mathbf{d}}]_X^T \partial_t \tilde{\mathbf{d}} \otimes (\mathbf{S} - \nabla \tilde{\mathbf{d}}))) \otimes \tilde{\mathbf{d}} + ([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X)^T \nabla ([\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}}) \otimes \tilde{\mathbf{d}} \right) \right) dt \\
& = \int_0^T \phi \langle \langle v_\tau, (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) : \Theta : (\Upsilon : ([\tilde{\mathbf{d}}]_X^T \partial_t \tilde{\mathbf{d}} \otimes (\nabla \tilde{\mathbf{d}} - \mathbf{S}))) \otimes (\mathbf{h} - \tilde{\mathbf{d}}) \rangle \rangle dt \\
& + \int_0^T \phi \langle \langle v_\tau, (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) : \Theta : (\Upsilon : ([\tilde{\mathbf{d}}]_X^T \partial_t \tilde{\mathbf{d}} \otimes (\nabla \tilde{\mathbf{d}} - \mathbf{S}))) \otimes \tilde{\mathbf{d}} \rangle \rangle dt \\
& + \int_0^T \phi \left(\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} ; \Theta : (\Upsilon : ([\tilde{\mathbf{d}}]_X^T \partial_t \tilde{\mathbf{d}} \otimes (\mathbf{S} - \nabla \tilde{\mathbf{d}}))) \otimes (\tilde{\mathbf{d}} - \mathbf{d}) \right) dt \\
& + \int_0^T \phi \left((\nabla \mathbf{d} \otimes \mathbf{d} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) ; \Theta : ([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X)^T \nabla ([\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}}) \otimes (\mathbf{d} - \tilde{\mathbf{d}}) \right) dt \\
& + \int_0^T \phi \left((\nabla \mathbf{d} \otimes \mathbf{d} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) ; \Theta : ([\tilde{\mathbf{d}}]_X - [\mathbf{d}]_X)^T \nabla ([\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}}) \otimes \tilde{\mathbf{d}} \right) dt \\
& + \int_0^T \phi \left(\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} ; \Theta : ([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X)^T \nabla ([\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}}) \otimes (\tilde{\mathbf{d}} - \mathbf{d}) \right) dt.
\end{aligned}$$

This can now be estimated by

$$\begin{aligned}
& \int_0^T \phi \langle v_\tau, \mathbf{S} \otimes \mathbf{h} : \Theta : \Upsilon : ([\tilde{\mathbf{d}}]_X^T \partial_t \tilde{\mathbf{d}} \otimes (\nabla \tilde{\mathbf{d}} - \mathbf{S})) \otimes \mathbf{h} \rangle dt \\
& + \int_0^T \phi \left(\nabla \mathbf{d} \otimes \mathbf{d} : \Theta : ([\tilde{\mathbf{d}}]_X - [\mathbf{d}]_X)^T \nabla ([\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}}) \otimes \mathbf{d} \right) dt \\
& + \int_0^T \phi \left(\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} : \Theta : \left(\Upsilon : ([\tilde{\mathbf{d}}]_X^T \partial_t \tilde{\mathbf{d}} \otimes (\mathbf{S} - \nabla \tilde{\mathbf{d}})) \otimes \tilde{\mathbf{d}} + ([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X)^T \nabla ([\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}}) \otimes \tilde{\mathbf{d}} \right) \right) dt \\
& \leq \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)} \int_0^T \phi \|\partial_t \tilde{\mathbf{d}}\|_{L^\infty} \langle v_t, |\Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})|^2 + |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 |\mathbf{h} - \tilde{\mathbf{d}}|^2 \rangle dt \\
& + \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^2 \int_0^T \phi \|\partial_t \tilde{\mathbf{d}}\|_{L^\infty} \langle v_t, |\Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})|^2 + |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 \rangle dt \\
& + \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{W}^{1,3})} \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^2 \int_0^T \phi \|\partial_t \tilde{\mathbf{d}}\|_{L^\infty} (\langle v_t, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 \rangle + \|\mathbf{d} - \tilde{\mathbf{d}}\|_{\mathbf{L}^6}^2) dt \\
& + \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{W}^{1,3})} \int_0^T \phi \|\partial_t \tilde{\mathbf{d}}\|_{L^\infty} \left(\langle v_t, |\Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})|^2 \rangle + \|\mathbf{d} - \tilde{\mathbf{d}}\|_{\mathbf{L}^{12}}^4 \right) dt \\
& + \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)} \int_0^T \phi \|\partial_t \tilde{\mathbf{d}}\|_{\mathbf{W}^{1,3}} \left(\langle v_t, |\Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})|^2 \rangle + \|\mathbf{d} - \tilde{\mathbf{d}}\|_{\mathbf{L}^{12}}^4 \right) dt \\
& + \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{W}^{1,3})} \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)} \int_0^T \phi \|\partial_t \tilde{\mathbf{d}}\|_{L^\infty} \left(\phi \langle v_t, |\Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})|^2 \rangle + \phi \|\mathbf{d} - \tilde{\mathbf{d}}\|_{\mathbf{L}^6}^2 \right) dt \\
& + \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^2 \int_0^T \phi \|\partial_t \tilde{\mathbf{d}}\|_{\mathbf{W}^{1,3}} \left(\langle v_t, |\Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})|^2 \rangle + \|\mathbf{d} - \tilde{\mathbf{d}}\|_{\mathbf{L}^6}^2 \right) dt \\
& + \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{W}^{1,3})} (\|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)} + \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{W}^{1,3})}) \int_0^T \phi (\|\partial_t \tilde{\mathbf{d}}\|_{\mathbf{W}^{1,3}} + \|\partial_t \tilde{\mathbf{d}}\|_{L^\infty}) \|\mathbf{d} - \tilde{\mathbf{d}}\|_{\mathbf{L}^6}^2 dt.
\end{aligned}$$

All the terms on the right-hand side can be bounded due to Lemma 4.1.

The last line of (5.14) and (5.15b) can be expressed as

$$\begin{aligned}
& \int_0^T \phi \langle v_\tau, \mathbf{S} \otimes \mathbf{h} : \Theta : \mathbf{S} \otimes ([\tilde{\mathbf{d}}]_X - [\mathbf{h}]_X)^T [\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}} \rangle dt + \int_0^T \phi \left(\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes ([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X)^T [\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}} \right) dt \\
& = \int_0^T \phi \langle v_\tau, (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) : \Theta : (\mathbf{S} - \nabla \tilde{\mathbf{d}}) \otimes ([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X)^T [\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}} \rangle dt \\
& + \int_0^T \phi \langle v_\tau, (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) : \Theta : \nabla \tilde{\mathbf{d}} \otimes ([\tilde{\mathbf{d}}]_X - [\mathbf{h}]_X)^T [\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}} \rangle dt \\
& + \int_0^T \phi \left(\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} : \Theta : (\nabla \tilde{\mathbf{d}} - \mathbf{S}) \otimes ([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X)^T [\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}} \right) dt.
\end{aligned}$$

Estimating the right hand side gives

$$\begin{aligned}
& \int_0^T \phi \langle v_\tau, \mathbf{S} \otimes \mathbf{h} : \Theta : \mathbf{S} \otimes ([\tilde{\mathbf{d}}]_X - [\mathbf{h}]_X)^T [\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}} \rangle dt + \int_0^T \phi \left(\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes ([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X)^T [\tilde{\mathbf{d}}]_X \partial_t \tilde{\mathbf{d}} \right) dt \\
& \leq \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)} \int_0^T \phi \|\partial_t \tilde{\mathbf{d}}\|_{L^\infty} \langle v_t, |\Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})|^2 + |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 |\mathbf{h} - \tilde{\mathbf{d}}|^2 \rangle dt \\
& + \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{W}^{1,3})} \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)} \int_0^T \phi \|\partial_t \tilde{\mathbf{d}}\|_{L^\infty} \left(\langle v_t, |\Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})|^2 \rangle + \|\mathbf{d} - \tilde{\mathbf{d}}\|_{\mathbf{L}^6}^2 \right) dt \\
& + \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{W}^{1,3})} \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^2 \int_0^T \phi \|\partial_t \tilde{\mathbf{d}}\|_{L^\infty} (\langle v_t, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 \rangle + \|\mathbf{d} - \tilde{\mathbf{d}}\|_{\mathbf{L}^6}^2) dt,
\end{aligned}$$

which can be bounded due to Lemma 4.1 by the relative Energy.

Putting all the estimates back into (5.10) and using Lemma 3.1 gives the asserted integration-by-parts formula (5.4). \square

Corollary 5.1. Let (v, \mathbf{d}) be a measure-valued solution and $(\tilde{v}, \tilde{\mathbf{d}})$ be a strong solution. Then there exists for every $\delta > 0$

a possibly large constant C_δ such that

$$\begin{aligned}
& -(\mathbf{v}(t), \tilde{\mathbf{v}}(t)) - (\nabla \mathbf{d}(t); \boldsymbol{\Lambda} : \nabla \tilde{\mathbf{d}}(t)) - (\nabla \mathbf{d}(t) \otimes \mathbf{d}(t); \boldsymbol{\Theta} : \nabla \tilde{\mathbf{d}}(t) \otimes \tilde{\mathbf{d}}(t)) \\
& \leq - \int_0^t [(\mathbf{v}, \partial_t \tilde{\mathbf{v}}) + (\partial_t \mathbf{v}, \tilde{\mathbf{v}})] \, ds - (\mathbf{v}(0), \tilde{\mathbf{v}}(0)) - \int_0^t [(\mathbf{d} \times \partial_t \mathbf{d}, \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}) + (\mathbf{d} \times \mathbf{q}, \tilde{\mathbf{d}} \times \partial_t \tilde{\mathbf{d}})] \, ds \\
& \quad - (\nabla \mathbf{d}(0); \boldsymbol{\Lambda} : \nabla \tilde{\mathbf{d}}(0)) - (\nabla \mathbf{d}(0) \otimes \mathbf{d}(0); \boldsymbol{\Theta} : \nabla \tilde{\mathbf{d}}(0) \otimes \tilde{\mathbf{d}}(0)) \\
& \quad + \frac{1}{2} \mathcal{E}(t) + \delta \int_0^t \mathcal{W}(s) \, ds + C_\delta \int_0^t \mathcal{E}(s) \, ds \\
& \quad + ((\nabla \mathbf{d}(0) - \nabla \tilde{\mathbf{d}}(0)) \otimes (\mathbf{d}(0) - \tilde{\mathbf{d}}(0))) : \boldsymbol{\Theta} : (\nabla \tilde{\mathbf{d}}(0) \otimes \tilde{\mathbf{d}}(0)) + c \|\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^2 \|\mathbf{d}(0) - \tilde{\mathbf{d}}(0)\|_{\mathbf{L}^2}^2
\end{aligned} \tag{5.16}$$

holds for a.e. $t \in (0, T)$.

Proof. First, we observe that the integration-by-parts formulae of Lemma 5.1 holds true for $s = 0$. Observe that $\mathbf{v} \in L^\infty(0, T; \mathbf{L}_\sigma^2) \cap W^{1,2}(0, T; (\mathbf{W}_{0,\sigma}^{1,3})^*) \hookrightarrow \mathcal{C}_w([0, T]; \mathbf{L}_\sigma^2)$ and $\tilde{\mathbf{v}} \in \mathcal{C}([0, T]; \mathbf{L}_\sigma^2)$ such that the limit $s \rightarrow 0$ in (5.1) is well defined. Since

$$\mathbf{d} \in L^\infty(0, T; \mathbf{H}^1) \cap \mathcal{AC}([0, T]; \mathbf{L}^{3/2}) \quad \text{and} \quad \tilde{\mathbf{d}} \in L^\infty(0, T; \mathbf{W}^{2,3}) \cap \mathcal{AC}([0, T]; \mathbf{W}^{1,3}).$$

A standard lemma (see Lions and Magenes [35, p. 297] shows that $\mathbf{d} \in \mathcal{C}_w([0, T]; \mathbf{H}^1)$. With the compact Sobolev embedding $\mathbf{H}^1 \hookrightarrow^c \mathbf{L}^5$ we see $\mathbf{d} \in \mathcal{C}([0, T]; \mathbf{L}^5)$. Similarly, we see $\mathbf{d} \in \mathcal{C}([0, T]; \mathbf{W}^{1,10/3})$. Together, we can estimate

$$\begin{aligned}
& |(\nabla \mathbf{d}(0) \otimes \mathbf{d}(0)) : \boldsymbol{\Theta} : (\nabla \tilde{\mathbf{d}}(0) \otimes \tilde{\mathbf{d}}(0)) - (\nabla \mathbf{d}(s) \otimes \mathbf{d}(s)) : \boldsymbol{\Theta} : (\nabla \tilde{\mathbf{d}}(s) \otimes \tilde{\mathbf{d}}(s))| \\
& \leq |((\nabla \mathbf{d}(0) - \nabla \mathbf{d}(s)) \otimes \mathbf{d}(0)) : \boldsymbol{\Theta} : (\nabla \tilde{\mathbf{d}}(0) \otimes \tilde{\mathbf{d}}(0))| \\
& \quad + \|\nabla \mathbf{d}(s)\|_{\mathbf{L}^2}^2 \|\mathbf{d}(0) - \mathbf{d}(s)\|_{\mathbf{L}^5}^2 \|\nabla \tilde{\mathbf{d}}(0)\|_{\mathbf{L}^{10/3}} \|\tilde{\mathbf{d}}(0)\|_{\mathbf{L}^\infty} \\
& \quad + \|\nabla \mathbf{d}(s)\|_{\mathbf{L}^2}^2 \|\mathbf{d}(s)\|_{\mathbf{L}^6}^2 \|\nabla \tilde{\mathbf{d}}(0) - \nabla \tilde{\mathbf{d}}(s)\|_{\mathbf{L}^{10/3}} \|\tilde{\mathbf{d}}(0)\|_{\mathbf{L}^{15/2}} \\
& \quad + \|\nabla \mathbf{d}(s)\|_{\mathbf{L}^2}^2 \|\mathbf{d}(s)\|_{\mathbf{L}^6}^2 \|\nabla \tilde{\mathbf{d}}(s)\|_{\mathbf{L}^{10/3}} \|\tilde{\mathbf{d}}(0) - \tilde{\mathbf{d}}(s)\|_{\mathbf{L}^{15/2}},
\end{aligned}$$

and the limit $s \rightarrow 0$ vanishes due to the weak convergence of the gradient of \mathbf{d} and the strong convergence of the other terms. The additional terms depending on s in (5.2), (5.3) and (5.4) can be handled similar, even somehow simpler since they are more regular. For the term (5.4a) depending on t we observe that we can represent it due to (2.12) almost everywhere via the generalized gradient Young measure (see Definition (2)) and apply Youngs inequality

$$\begin{aligned}
& ((\nabla \mathbf{d}(t) - \nabla \tilde{\mathbf{d}}(t)) \otimes (\mathbf{d}(t) - \tilde{\mathbf{d}}(t))) : \boldsymbol{\Theta} : \nabla \tilde{\mathbf{d}}(t) \otimes \tilde{\mathbf{d}}(t) \\
& = \int_{\Omega} \left\langle v_{(\mathbf{x}, t)}^o, ((\mathbf{S} - \nabla \tilde{\mathbf{d}}(t)) \otimes (\mathbf{d}(t) - \tilde{\mathbf{d}}(t))) : \boldsymbol{\Theta} : (\nabla \tilde{\mathbf{d}}(t) \otimes \tilde{\mathbf{d}}(t)) \right\rangle \, d\mathbf{x} \\
& \leq \frac{k}{4} \int_{\Omega} \left\langle v_{(\mathbf{x}, t)}^o, |\mathbf{S} - \nabla \tilde{\mathbf{d}}(t)|^2 \right\rangle \, d\mathbf{x} + c \|\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^2 \int_{\Omega} |\mathbf{d}(t) - \tilde{\mathbf{d}}(t)|^2 \, d\mathbf{x} \\
& \leq \frac{k}{4} \langle v_t, |\mathbf{S} - \nabla \tilde{\mathbf{d}}(t)|^2 \rangle + c \|\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^2 \|\mathbf{d}(t) - \tilde{\mathbf{d}}(t)\|_{\mathbf{L}^2}^2
\end{aligned} \tag{5.17}$$

Remark that the defect measure m_t is a positive measure, $m_t \in \mathcal{M}^+(\overline{\Omega})$. The constant k is chosen small enough, such that

$$k \langle v_t, |\mathbf{S} - \nabla \tilde{\mathbf{d}}(t)|^2 \rangle \leq \langle v_t, (\mathbf{S} - \nabla \tilde{\mathbf{d}}(t)) : \boldsymbol{\Lambda} : (\mathbf{S} - \nabla \tilde{\mathbf{d}}(t)) \rangle \leq 2\mathcal{E}(t)$$

The second term on the right-hand side of (5.17) can be estimated by formula (5.2). Adding the integration-by-parts formulae (5.1), (5.2) and (5.4) for $s = 0$ and estimating the term in line (5.4a) by (5.17) implies the assertion (5.16). The definitions (2.9) and (2.11c) of the variational derivatives are inserted for the strong solution and the measure-valued solution, respectively.

□

6 Proof of the main result

The following lemma implies the main result Theorem 2.1 and is proven in this Section.

Proposition 6.1. Let (\mathbf{v}, \mathbf{d}) and $(\mathbf{v}, m, \mathbf{v}^\infty)$ be a suitable measure-valued solution for initial values $(\mathbf{v}_0, \mathbf{d}_0)$ and let $(\tilde{\mathbf{v}}, \tilde{\mathbf{d}})$ be a strong solution for the initial values $(\tilde{\mathbf{v}}_0, \tilde{\mathbf{d}}_0)$ fulfilling the same boundary conditions. Then there exists a possible large constant c such that

$$\frac{1}{2}\mathcal{E}(t) \leq \left(\mathcal{E}(0) + ((\nabla \mathbf{d}(0) - \nabla \tilde{\mathbf{d}}(0)) \otimes (\mathbf{d}(0) - \tilde{\mathbf{d}}(0)) ; \Theta : \nabla \tilde{\mathbf{d}}(0) \otimes \tilde{\mathbf{d}}(0)) \right. \quad (6.1)$$

$$\left. + c \|\mathbf{d}(0) - \tilde{\mathbf{d}}(0)\|_{L^2}^2 \right) e^{\int_0^t \mathcal{K}(s) ds}, \quad (6.2)$$

where \mathcal{K} is bounded in $L^1(0, T)$.

Remark 8. The result of Proposition 6.1 implies a continuous dependence of the relative energy (4.1) on the difference of the initial values.

Proof. Considering the relative energy, we observe

$$\begin{aligned} \mathcal{E}(t) = & \langle \mu_t, 1 \rangle + \langle \mathbf{v}_t, F \rangle + \frac{1}{2} \|\mathbf{v}(t)\|_{L^2}^2 + \mathcal{F}(\tilde{\mathbf{d}}(t), \nabla \tilde{\mathbf{d}}(t)) + \frac{1}{2} \|\tilde{\mathbf{v}}(t)\|_{L^2}^2 \\ & - (\mathbf{v}(t), \tilde{\mathbf{v}}(t)) - (\nabla \mathbf{d}(t); \Lambda : \nabla \tilde{\mathbf{d}}(t)) - (\nabla \mathbf{d}(t) \otimes \mathbf{d}(t) ; \Theta : \nabla \tilde{\mathbf{d}}(t) \otimes \tilde{\mathbf{d}}(t)). \end{aligned}$$

We insert the energy inequality for the measure-valued solution and the energy equality for the strong solution. Adding the integral over the relative dissipation gives

$$\begin{aligned} \mathcal{E}(t) + \int_0^t \mathcal{W}(s) ds \leq & \mathcal{F}(\mathbf{d}_0) + \frac{1}{2} \|\mathbf{v}_0\|_{L^2}^2 + \mathcal{F}(\tilde{\mathbf{d}}_0) + \frac{1}{2} \|\tilde{\mathbf{v}}_0\|_{L^2}^2 + \int_0^t \langle \mathbf{g}, \mathbf{v} + \tilde{\mathbf{v}} \rangle ds \\ & - 2(\mu_1 + \lambda(\mu_2 + \mu_3)) \int_0^t (\mathbf{d} \cdot (\nabla \mathbf{v})_{\text{sym}} \mathbf{d}, \tilde{\mathbf{d}} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}) ds - 2\mu_4 \int_0^t ((\nabla \mathbf{v})_{\text{sym}}; (\nabla \tilde{\mathbf{v}})_{\text{sym}}) ds \\ & - 2(\mu_5 + \mu_6 - \lambda(\mu_2 + \mu_3)) \int_0^t ((\nabla \mathbf{v})_{\text{sym}} \mathbf{d}, (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}) ds - 2 \int_0^t (\mathbf{d} \times \mathbf{q}, \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}) ds \\ & + ((\mu_2 + \mu_3) - \lambda) \int_0^t [(\mathbf{d} \times \mathbf{q}, \mathbf{d} \times (\nabla \mathbf{v})_{\text{sym}} \mathbf{d}) + (\tilde{\mathbf{d}} \times \tilde{\mathbf{q}}, \tilde{\mathbf{d}} \times (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}})] ds \\ & - (\mathbf{v}(t), \tilde{\mathbf{v}}(t)) - (\nabla \mathbf{d}(t); \Lambda : \nabla \tilde{\mathbf{d}}(t)) - (\nabla \mathbf{d}(t) \otimes \mathbf{d}(t) ; \Theta : \nabla \tilde{\mathbf{d}}(t) \otimes \tilde{\mathbf{d}}(t)). \end{aligned}$$

Using Corollary 5.1 gives

$$\begin{aligned} \mathcal{E}(t) + \int_0^t \mathcal{W}(s) ds \leq & \mathcal{F}(\mathbf{d}_0) + \frac{1}{2} \|\mathbf{v}_0\|_{L^2}^2 + \mathcal{F}(\tilde{\mathbf{d}}_0) + \frac{1}{2} \|\tilde{\mathbf{v}}_0\|_{L^2}^2 \\ & - 2(\mu_1 + \lambda(\mu_2 + \mu_3)) \int_0^t (\mathbf{d} \cdot (\nabla \mathbf{v})_{\text{sym}} \mathbf{d}, \tilde{\mathbf{d}} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}) ds \\ & - 2(\mu_5 + \mu_6 - \lambda(\mu_2 + \mu_3)) \int_0^t ((\nabla \mathbf{v})_{\text{sym}} \mathbf{d}, (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}) ds \\ & - 2\mu_4 \int_0^t ((\nabla \mathbf{v})_{\text{sym}}; (\nabla \tilde{\mathbf{v}})_{\text{sym}}) ds - 2 \int_0^t (\mathbf{d} \times \mathbf{q}, \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}) ds + \int_0^t \langle \mathbf{g}, \mathbf{v} + \tilde{\mathbf{v}} \rangle ds \\ & + ((\mu_2 + \mu_3) - \lambda) \int_0^t [(\mathbf{d} \times \mathbf{q}, \mathbf{d} \times (\nabla \mathbf{v})_{\text{sym}} \mathbf{d}) + (\tilde{\mathbf{d}} \times \tilde{\mathbf{q}}, \tilde{\mathbf{d}} \times (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}})] ds - (\mathbf{v}(0), \tilde{\mathbf{v}}(0)) \\ & - \int_0^t [(\mathbf{v}, \partial_t \tilde{\mathbf{v}}) + (\partial_t \mathbf{v}, \tilde{\mathbf{v}})] ds - (\nabla \mathbf{d}(0); \Lambda : \nabla \tilde{\mathbf{d}}(0)) - (\nabla \mathbf{d}(0) \otimes \mathbf{d}(0) ; \Theta : \nabla \tilde{\mathbf{d}}(0) \otimes \tilde{\mathbf{d}}(0)) \\ & - \int_0^t [(\mathbf{d} \times \partial_t \mathbf{d}, \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}) + (\mathbf{d} \times \mathbf{q}, \tilde{\mathbf{d}} \times \partial_t \tilde{\mathbf{d}})] ds + \frac{1}{2} \mathcal{E}(t) + \delta \int_0^t \mathcal{W}(s) ds + C_\delta \int_0^t \mathcal{E}(s) ds \\ & + ((\nabla \mathbf{d}(0) - \nabla \tilde{\mathbf{d}}(0)) \otimes (\mathbf{d}(0) - \tilde{\mathbf{d}}(0))) : \Theta : (\nabla \tilde{\mathbf{d}}(0) \otimes \tilde{\mathbf{d}}(0)) \\ & + c \|\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}\|_{L^\infty(L^\infty)}^2 \|\mathbf{d}(0) - \tilde{\mathbf{d}}(0)\|_{L^2}^2 \end{aligned} \quad (6.3)$$

Regarding the terms incorporating the initial values, we see

$$\begin{aligned} \mathcal{F}(\mathbf{d}_0) + \frac{1}{2} \|\mathbf{v}_0\|_{L^2}^2 + \mathcal{F}(\tilde{\mathbf{d}}_0) + \frac{1}{2} \|\tilde{\mathbf{v}}_0\|_{L^2}^2 - (\mathbf{v}(0), \tilde{\mathbf{v}}(0)) - (\nabla \mathbf{d}(0); \Lambda : \nabla \tilde{\mathbf{d}}(0)) \\ - (\nabla \mathbf{d}(0) \otimes \mathbf{d}(0) ; \Theta : \nabla \tilde{\mathbf{d}}(0) \otimes \tilde{\mathbf{d}}(0)) = \mathcal{E}(0). \end{aligned}$$

In the next step, we use that (\mathbf{v}, \mathbf{d}) and $(\tilde{\mathbf{v}}, \tilde{\mathbf{d}})$ are measure-valued and strong solutions, respectively. For the measure-valued solution, we consider equation (2.11a) tested with $\tilde{\mathbf{v}}$ and the equation (2.11b) tested with $\tilde{\mathbf{d}} \times \tilde{\mathbf{q}}$ and added up:

$$\begin{aligned} & - \int_0^T [(\partial_t \mathbf{v}, \tilde{\mathbf{v}}) + (\mathbf{d} \times \partial_t \mathbf{d}, \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}) + \mu_4 ((\nabla \mathbf{v})_{\text{sym}}, (\nabla \tilde{\mathbf{v}})_{\text{sym}}) + (\mathbf{d} \times \mathbf{q}, \tilde{\mathbf{d}} \times \tilde{\mathbf{q}})] dt \\ &= \int_0^T [((\mathbf{v} \cdot \nabla) \mathbf{v}, \tilde{\mathbf{v}}) + \mu_1 (\mathbf{d} \cdot (\nabla \mathbf{v})_{\text{sym}}, \mathbf{d} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \mathbf{d}) + (\mu_5 + \mu_6) ((\nabla \mathbf{v})_{\text{sym}} \mathbf{d}, (\nabla \tilde{\mathbf{v}})_{\text{sym}} \mathbf{d})] dt \\ &\quad + \int_0^T [(\mu_2 + \mu_3) ((\nabla \tilde{\mathbf{v}})_{\text{sym}} \mathbf{d}, \mathbf{e}) + \lambda ((\nabla \mathbf{v})_{\text{sym}} \mathbf{d}, (\nabla \tilde{\mathbf{v}})_{\text{skw}} \mathbf{d}) + (\mathbf{e}, (\nabla \tilde{\mathbf{v}})_{\text{skw}} \mathbf{d})] dt \\ &\quad - \int_0^T [\langle\langle v_t, \mathbf{S}^T F_{\mathbf{S}}(\mathbf{h}, \mathbf{S}) : \nabla \tilde{\mathbf{v}} \rangle\rangle + 2 \langle\langle \mu_t, \boldsymbol{\Gamma} : (\boldsymbol{\Gamma} \cdot \nabla \tilde{\mathbf{v}}) \rangle\rangle] dt \\ &\quad + \int_0^T [([\mathbf{d}]_{\mathbf{X}} (\mathbf{v} \cdot \nabla) \mathbf{d}, [\tilde{\mathbf{d}}]_{\mathbf{X}} \tilde{\mathbf{q}}) - ([\mathbf{d}]_{\mathbf{X}} (\nabla \mathbf{v})_{\text{skw}} \mathbf{d}, [\tilde{\mathbf{d}}]_{\mathbf{X}} \tilde{\mathbf{q}}) + \lambda ([\mathbf{d}]_{\mathbf{X}} (\nabla \mathbf{v})_{\text{sym}} \mathbf{d}, [\tilde{\mathbf{d}}]_{\mathbf{X}} \tilde{\mathbf{q}})] dt. \end{aligned}$$

Using Corollary 3.1 and equation (2.11b) to replace the term \mathbf{e} shows

$$\begin{aligned} & - \int_0^T [(\partial_t \mathbf{v}, \tilde{\mathbf{v}}) + (\mathbf{d} \times \partial_t \mathbf{d}, \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}) + \mu_4 ((\nabla \mathbf{v})_{\text{sym}}, (\nabla \tilde{\mathbf{v}})_{\text{sym}}) + (\mathbf{d} \times \mathbf{q}, \tilde{\mathbf{d}} \times \tilde{\mathbf{q}})] dt \\ &= \int_0^T [((\mathbf{v} \cdot \nabla) \mathbf{v}, \tilde{\mathbf{v}}) + \mu_1 (\mathbf{d} \cdot (\nabla \mathbf{v})_{\text{sym}}, \mathbf{d} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \mathbf{d}) + (\mu_5 + \mu_6) ((\nabla \mathbf{v})_{\text{sym}} \mathbf{d}, (\nabla \tilde{\mathbf{v}})_{\text{sym}} \mathbf{d})] dt \\ &\quad - \int_0^T [(\mu_2 + \mu_3) (\mathbf{d} \times (\nabla \tilde{\mathbf{v}})_{\text{sym}} \mathbf{d}, \mathbf{d} \times \mathbf{q}) + \lambda (\mu_2 + \mu_3) (\mathbf{d} \times (\nabla \tilde{\mathbf{v}})_{\text{sym}} \mathbf{d}, \mathbf{d} \times (\nabla \mathbf{v})_{\text{sym}}) + (\mathbf{d} \times \mathbf{q}, \mathbf{d} \times (\nabla \tilde{\mathbf{v}})_{\text{skw}} \mathbf{d})] dt \\ &\quad - \int_0^T [\langle\langle v_t, \mathbf{S}^T F_{\mathbf{S}}(\mathbf{h}, \mathbf{S}) : \nabla \tilde{\mathbf{v}} \rangle\rangle + 2 \langle\langle \mu_t, \boldsymbol{\Gamma} : (\boldsymbol{\Gamma} \cdot \nabla \tilde{\mathbf{v}}) \rangle\rangle] dt \\ &\quad + \int_0^T [([\mathbf{d}]_{\mathbf{X}} (\mathbf{v} \cdot \nabla) \mathbf{d}, [\tilde{\mathbf{d}}]_{\mathbf{X}} \tilde{\mathbf{q}}) - ([\mathbf{d}]_{\mathbf{X}} (\nabla \mathbf{v})_{\text{skw}} \mathbf{d}, [\tilde{\mathbf{d}}]_{\mathbf{X}} \tilde{\mathbf{q}}) + \lambda ([\mathbf{d}]_{\mathbf{X}} (\nabla \mathbf{v})_{\text{sym}} \mathbf{d}, [\tilde{\mathbf{d}}]_{\mathbf{X}} \tilde{\mathbf{q}})] dt. \end{aligned} \tag{6.4}$$

Remark that $\mathbf{d} \cdot (\nabla \tilde{\mathbf{v}})_{\text{skw}} \mathbf{d} = 0$. The strong solution is treated similar, by testing equation (2.1a) with \mathbf{v} and equation (2.1b) with $\mathbf{d} \times \mathbf{q}$ and adding them up:

$$\begin{aligned} & - \int_0^T [(\partial_t \tilde{\mathbf{v}}, \mathbf{v}) + (\tilde{\mathbf{d}} \times \partial_t \tilde{\mathbf{d}}, \mathbf{d} \times \mathbf{q}) + \mu_4 ((\nabla \tilde{\mathbf{v}})_{\text{sym}}, (\nabla \mathbf{v})_{\text{sym}}) + (\tilde{\mathbf{d}} \times \tilde{\mathbf{q}}, \mathbf{d} \times \mathbf{q})] dt \\ &= \int_0^T [((\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{v}}, \mathbf{v}) + \mu_1 (\tilde{\mathbf{d}} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, \tilde{\mathbf{d}} \cdot (\nabla \mathbf{v})_{\text{sym}} \tilde{\mathbf{d}}) + (\mu_5 + \mu_6) ((\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, (\nabla \mathbf{v})_{\text{sym}} \tilde{\mathbf{d}})] dt \\ &\quad + \int_0^T [(\mu_2 + \mu_3) (\tilde{\mathbf{e}}, (\nabla \mathbf{v})_{\text{sym}} \tilde{\mathbf{d}}) + \lambda ((\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, (\nabla \mathbf{v})_{\text{skw}} \tilde{\mathbf{d}}) + (\tilde{\mathbf{e}}, (\nabla \mathbf{v})_{\text{skw}} \tilde{\mathbf{d}})] dt \\ &\quad - \int_0^T [([\tilde{\mathbf{d}}]_{\mathbf{X}} ((\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} - (\nabla \tilde{\mathbf{v}})_{\text{skw}} \tilde{\mathbf{d}} + \lambda (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}), [\tilde{\mathbf{d}}]_{\mathbf{X}} \mathbf{q}) + (\nabla \tilde{\mathbf{d}}^T F_{\mathbf{S}}(\tilde{\mathbf{d}}, \nabla \tilde{\mathbf{d}}); \nabla \mathbf{v})] dt. \end{aligned}$$

Using Corollary 3.1 and equation (2.1b) to replace the term $\tilde{\mathbf{e}}$ shows

$$\begin{aligned} & - \int_0^T [(\partial_t \tilde{\mathbf{v}}, \mathbf{v}) + (\tilde{\mathbf{d}} \times \partial_t \tilde{\mathbf{d}}, \mathbf{d} \times \mathbf{q}) + \mu_4 ((\nabla \tilde{\mathbf{v}})_{\text{sym}}, (\nabla \mathbf{v})_{\text{sym}}) + (\tilde{\mathbf{d}} \times \tilde{\mathbf{q}}, \mathbf{d} \times \mathbf{q})] dt \\ &= \int_0^T [((\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{v}}, \mathbf{v}) + \mu_1 (\tilde{\mathbf{d}} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, \tilde{\mathbf{d}} \cdot (\nabla \mathbf{v})_{\text{sym}} \tilde{\mathbf{d}}) + (\mu_5 + \mu_6) ((\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, (\nabla \mathbf{v})_{\text{sym}} \tilde{\mathbf{d}})] dt \\ &\quad - \int_0^T [(\mu_2 + \mu_3) (\tilde{\mathbf{d}} \times \tilde{\mathbf{q}}, \tilde{\mathbf{d}} \times (\nabla \mathbf{v})_{\text{sym}} \tilde{\mathbf{d}}) + \lambda (\mu_2 + \mu_3) (\tilde{\mathbf{d}} \times (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, \tilde{\mathbf{d}} \times (\nabla \mathbf{v})_{\text{sym}} \tilde{\mathbf{d}})] dt \\ &\quad - \int_0^T [(\tilde{\mathbf{d}} \times \tilde{\mathbf{q}}, \tilde{\mathbf{d}} \times (\nabla \mathbf{v})_{\text{skw}} \tilde{\mathbf{d}}) + ([\tilde{\mathbf{d}}]_{\mathbf{X}} ((\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} - (\nabla \tilde{\mathbf{v}})_{\text{skw}} \tilde{\mathbf{d}} + \lambda (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}), [\tilde{\mathbf{d}}]_{\mathbf{X}} \mathbf{q}) + (\nabla \tilde{\mathbf{d}}^T F_{\mathbf{S}}(\tilde{\mathbf{d}}, \nabla \tilde{\mathbf{d}}); \nabla \mathbf{v})] dt. \end{aligned} \tag{6.5}$$

Remark that $\tilde{\mathbf{d}} \cdot (\nabla \mathbf{v})_{\text{skw}} \tilde{\mathbf{d}} = 0$. Inserting the equations (6.4) and (6.5) in (6.3) gives

$$\begin{aligned}
& \frac{1}{2} \mathcal{E}(t) + \int_0^t \mathcal{W}(s) \, ds - \mathcal{E}(0) - ((\nabla \mathbf{d}(0) - \nabla \tilde{\mathbf{d}}(0)) \otimes (\mathbf{d}(0) - \tilde{\mathbf{d}}(0))) : \Theta : (\nabla \tilde{\mathbf{d}}(0) \otimes \tilde{\mathbf{d}}(0)) \\
& \quad - c \|\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^2 \|\mathbf{d}(0) - \tilde{\mathbf{d}}(0)\|_{L^2}^2 \\
& \leq \delta \int_0^t \mathcal{W}(s) \, ds + C_\delta \int_0^t \mathcal{E}(s) \, ds + \int_0^t [((\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{v}}, \mathbf{v}) + ((\mathbf{v} \cdot \nabla) \mathbf{v}, \tilde{\mathbf{v}})] \, ds \\
& \quad + \mu_1 \int_0^t [(\tilde{\mathbf{d}} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, (\nabla \mathbf{v})_{\text{sym}} : (\tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} - \mathbf{d} \otimes \mathbf{d})) + (\mathbf{d} \cdot (\nabla \mathbf{v})_{\text{sym}} \mathbf{d}, (\nabla \tilde{\mathbf{v}})_{\text{sym}} : (\mathbf{d} \otimes \mathbf{d} - \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}))] \, ds \\
& \quad + (\mu_5 + \mu_6) \int_0^t [((\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, (\nabla \mathbf{v})_{\text{sym}} (\tilde{\mathbf{d}} - \mathbf{d})) + ((\nabla \mathbf{v})_{\text{sym}} \mathbf{d}, (\nabla \tilde{\mathbf{v}})_{\text{sym}} (\mathbf{d} - \tilde{\mathbf{d}}))] \, dt \\
& \quad + \int_0^t [(\tilde{\mathbf{d}} \times (\nabla \tilde{\mathbf{v}})_{\text{skw}} \tilde{\mathbf{d}} - \mathbf{d} \times (\nabla \tilde{\mathbf{v}})_{\text{skw}} \mathbf{d}, \mathbf{d} \times \mathbf{q}) + (\mathbf{d} \times (\nabla \mathbf{v})_{\text{skw}} \mathbf{d} - \tilde{\mathbf{d}} \times (\nabla \mathbf{v})_{\text{skw}} \tilde{\mathbf{d}}, \tilde{\mathbf{d}} \times \tilde{\mathbf{q}})] \, ds \\
& \quad - (\mu_2 + \mu_3) \int_0^t [(\tilde{\mathbf{d}} \times ((\nabla \mathbf{v})_{\text{sym}} - (\nabla \tilde{\mathbf{v}})_{\text{sym}}) \tilde{\mathbf{d}}, \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}) + (\mathbf{d} \times ((\nabla \tilde{\mathbf{v}})_{\text{sym}} - (\nabla \mathbf{v})_{\text{sym}}) \mathbf{d}, \mathbf{d} \times \mathbf{q})] \, ds \\
& \quad + \lambda \int_0^t [(\mathbf{d} \times (\nabla \mathbf{v})_{\text{sym}} \mathbf{d} - \tilde{\mathbf{d}} \times (\nabla \mathbf{v})_{\text{sym}} \tilde{\mathbf{d}}, \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}) + (\tilde{\mathbf{d}} \times (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}} - \mathbf{d} \times (\nabla \mathbf{v})_{\text{sym}} \mathbf{d}, \mathbf{d} \times \mathbf{q})] \, ds \\
& \quad - 2\lambda(\mu_2 + \mu_3) \int_0^t [(\tilde{\mathbf{d}} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, \mathbf{d} \cdot (\nabla \mathbf{v})_{\text{sym}} \mathbf{d}) - ((\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, (\nabla \mathbf{v})_{\text{sym}} \mathbf{d})] \, ds \\
& \quad - \lambda(\mu_2 + \mu_3) \int_0^t [(\mathbf{d} \times (\nabla \mathbf{v})_{\text{sym}} \mathbf{d}, \mathbf{d} \times (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}) + (\tilde{\mathbf{d}} \times (\nabla \mathbf{v})_{\text{sym}} \tilde{\mathbf{d}}, \tilde{\mathbf{d}} \times (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}})] \, ds \\
& \quad + \int_0^t [(\tilde{\mathbf{d}} \times (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}}, \mathbf{d} \times \mathbf{q}) + (\mathbf{d} \times (\mathbf{v} \cdot \nabla) \mathbf{d}, \tilde{\mathbf{d}} \times \tilde{\mathbf{q}})] \, ds \\
& \quad - \int_0^t [(\nabla \tilde{\mathbf{d}}^T F_{\mathbf{S}}(\tilde{\mathbf{d}}, \nabla \tilde{\mathbf{d}}); \nabla \mathbf{v}) + \langle\langle \mathbf{v}_s, \mathbf{S}^T F_{\mathbf{S}}(\mathbf{h}, \mathbf{S}) : \nabla \tilde{\mathbf{v}} \rangle\rangle] \, ds - 2 \int_0^t \langle\langle \mu_s, \boldsymbol{\Gamma} : (\boldsymbol{\Gamma} \cdot \nabla \tilde{\mathbf{v}}) \rangle\rangle \, ds \\
& = \delta \int_0^t \mathcal{W}(s) \, ds + C_\delta \int_0^t \mathcal{E}(s) \, ds + I_1 + \mu_1 I_2 + (\mu_5 + \mu_6) I_3 + I_4 - (\mu_2 + \mu_3) I_5 + \lambda I_6 \\
& \quad - 2\lambda(\mu_2 + \mu_3) I_7 - \lambda(\mu_2 + \mu_3) I_8 + I_9 - 2I_{10}.
\end{aligned}$$

The terms are estimated individually by integrals over \mathcal{E} and \mathcal{W} . We start with I_1 . After rearranging one can estimate

$$\begin{aligned}
I_1 &= \int_0^t [((\mathbf{v} \cdot \nabla(\mathbf{v} - \tilde{\mathbf{v}}), \tilde{\mathbf{v}} - \mathbf{v}) + (((\mathbf{v} - \tilde{\mathbf{v}}) \cdot \nabla) \tilde{\mathbf{v}}, \tilde{\mathbf{v}} - \mathbf{v})] \, ds \\
&= \int_0^t ((\mathbf{v} - \tilde{\mathbf{v}}) \otimes (\tilde{\mathbf{v}} - \mathbf{v}); (\nabla \tilde{\mathbf{v}})_{\text{sym}}) \, ds \\
&\leq C_\delta \int_0^t \|(\nabla \tilde{\mathbf{v}})_{\text{sym}}\|_{L^3}^2 \|\mathbf{v} - \tilde{\mathbf{v}}\|_{L^2}^2 \, ds + \delta \int_0^t \|\mathbf{v} - \tilde{\mathbf{v}}\|_{L^6}^2 \, ds.
\end{aligned}$$

The term I_2 can be rearranged as

$$\begin{aligned}
I_2 &= \int_0^T (\tilde{\mathbf{d}} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, (\nabla \mathbf{v})_{\text{sym}} : (\tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} - \mathbf{d} \otimes \mathbf{d})) \, dt + \int_0^t (\mathbf{d} \cdot (\nabla \mathbf{v})_{\text{sym}} \mathbf{d}, (\nabla \tilde{\mathbf{v}})_{\text{sym}} : (\mathbf{d} \otimes \mathbf{d} - \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) \, ds \\
&= \int_0^t (\tilde{\mathbf{d}} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, ((\nabla \mathbf{v})_{\text{sym}} - (\nabla \tilde{\mathbf{v}})_{\text{sym}}) : (\tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} - \mathbf{d} \otimes \mathbf{d})) \, ds \\
&\quad + \int_0^t (\mathbf{d} \cdot (\nabla \mathbf{v})_{\text{sym}} \mathbf{d} - \tilde{\mathbf{d}} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, (\nabla \tilde{\mathbf{v}})_{\text{sym}} : (\mathbf{d} \otimes \mathbf{d} - \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) \, ds \\
&= \int_0^t (\tilde{\mathbf{d}} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, ((\nabla \mathbf{v})_{\text{sym}} - (\nabla \tilde{\mathbf{v}})_{\text{sym}}) : (\tilde{\mathbf{d}} \otimes (\tilde{\mathbf{d}} - \mathbf{d}))) \, ds \\
&\quad + \int_0^t (\tilde{\mathbf{d}} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, ((\nabla \mathbf{v})_{\text{sym}} \mathbf{d} - (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}) \cdot (\tilde{\mathbf{d}} - \mathbf{d})) \, ds + \int_0^t (\tilde{\mathbf{d}} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, (\nabla \tilde{\mathbf{v}})_{\text{sym}} : (\tilde{\mathbf{d}} - \mathbf{d}) \otimes (\tilde{\mathbf{d}} - \mathbf{d})) \, ds \\
&\quad + 2 \int_0^t (\mathbf{d} \cdot (\nabla \mathbf{v})_{\text{sym}} \mathbf{d} - \tilde{\mathbf{d}} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, (\nabla \tilde{\mathbf{v}})_{\text{sym}} : (\mathbf{d} - \tilde{\mathbf{d}}) \otimes \tilde{\mathbf{d}}) \, ds \\
&\quad + \int_0^t (\mathbf{d} \cdot (\nabla \mathbf{v})_{\text{sym}} \mathbf{d} - \tilde{\mathbf{d}} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, (\nabla \tilde{\mathbf{v}})_{\text{sym}} : (\mathbf{d} - \tilde{\mathbf{d}}) \otimes (\mathbf{d} - \tilde{\mathbf{d}})) \, ds.
\end{aligned}$$

It is thus possible to estimate I_2 by

$$\begin{aligned} I_2 &\leq \delta \int_0^t \|(\nabla \mathbf{v})_{\text{sym}} - (\nabla \tilde{\mathbf{v}})_{\text{sym}}\|_{L^2}^2 ds + C_\delta \|\tilde{\mathbf{d}}\|_{L^\infty(L^\infty)}^2 \int_0^t \|\tilde{\mathbf{d}} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}\|_{L^3}^2 \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2 ds \\ &\quad + \delta \int_0^t \|(\nabla \mathbf{v})_{\text{sym}} \mathbf{d} - (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}\|_{L^2}^2 ds + C_\delta \int_0^t \|\tilde{\mathbf{d}} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}\|_{L^3}^2 \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2 ds \\ &\quad + \int_0^t \|\tilde{\mathbf{d}} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}\|_{L^3} \|(\nabla \tilde{\mathbf{v}})_{\text{sym}}\|_{L^3} \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2 ds \\ &\quad + \delta \int_0^t \|\mathbf{d} \cdot (\nabla \mathbf{v})_{\text{sym}} \mathbf{d} - \tilde{\mathbf{d}} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}\|_{L^2}^2 ds + C_\delta \|\tilde{\mathbf{d}}\|_{L^\infty(L^\infty)}^2 \int_0^t \|(\nabla \tilde{\mathbf{v}})_{\text{sym}}\|_{L^3}^2 \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2 ds \\ &\quad + \delta \int_0^t \|\mathbf{d} \cdot (\nabla \mathbf{v})_{\text{sym}} \mathbf{d} - \tilde{\mathbf{d}} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}\|_{L^2}^2 ds + C_\delta \int_0^t \|(\nabla \tilde{\mathbf{v}})_{\text{sym}}\|_{L^3}^2 \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^{12}}^4 ds \\ &\leq \delta \int_0^t \mathcal{W}(s) ds + \int_0^t \mathcal{K}(s) \mathcal{E}(s) ds \end{aligned}$$

Regarding the term I_3 we observe

$$\begin{aligned} I_3 &= \int_0^t [((\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, (\nabla \mathbf{v})_{\text{sym}} (\tilde{\mathbf{d}} - \mathbf{d})) + ((\nabla \mathbf{v})_{\text{sym}} \mathbf{d}, (\nabla \tilde{\mathbf{v}})_{\text{sym}} (\mathbf{d} - \tilde{\mathbf{d}})))] ds \\ &= \int_0^t [(((\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, ((\nabla \mathbf{v})_{\text{sym}} - (\nabla \tilde{\mathbf{v}})_{\text{sym}}) (\tilde{\mathbf{d}} - \mathbf{d})) + ((\nabla \mathbf{v})_{\text{sym}} \mathbf{d} - (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, (\nabla \tilde{\mathbf{v}})_{\text{sym}} (\mathbf{d} - \tilde{\mathbf{d}})))] ds \\ &\leq \delta \int_0^t \|(\nabla \mathbf{v})_{\text{sym}} - (\nabla \tilde{\mathbf{v}})_{\text{sym}}\|_{L^2}^2 ds + C_\delta \int_0^t \|(\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}\|_{L^3}^2 \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2 ds \\ &\quad + \delta \int_0^t \|(\nabla \mathbf{v})_{\text{sym}} \mathbf{d} - (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}\|_{L^2}^2 ds + C_\delta \int_0^t \|(\nabla \tilde{\mathbf{v}})_{\text{sym}}\|_{L^3}^2 \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2 ds \\ &\leq \delta \int_0^t \mathcal{W}(s) ds + \int_0^t \mathcal{K}(s) \mathcal{E}(s) ds \end{aligned}$$

Similar, the term I_4 can be rearranged

$$\begin{aligned} I_4 &= \int_0^t [(\mathbf{d} \times \mathbf{q}, (\tilde{\mathbf{d}} - \mathbf{d}) \times (\nabla \tilde{\mathbf{v}})_{\text{skw}} \mathbf{d}) + (\tilde{\mathbf{d}} \times \tilde{\mathbf{q}}, (\mathbf{d} - \tilde{\mathbf{d}}) \times (\nabla \mathbf{v})_{\text{skw}} \tilde{\mathbf{d}})] ds \\ &\quad + \int_0^t [(\tilde{\mathbf{d}} \times (\nabla \tilde{\mathbf{v}})_{\text{skw}} (\tilde{\mathbf{d}} - \mathbf{d}), \mathbf{d} \times \mathbf{q}) + (\mathbf{d} \times (\nabla \mathbf{v})_{\text{skw}} (\mathbf{d} - \tilde{\mathbf{d}}), \tilde{\mathbf{d}} \times \tilde{\mathbf{q}})] ds \\ &= \int_0^t [(\mathbf{d} \times \mathbf{q} - \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}, (\tilde{\mathbf{d}} - \mathbf{d}) \times (\nabla \tilde{\mathbf{v}})_{\text{skw}} \mathbf{d}) + (\tilde{\mathbf{d}} \times \tilde{\mathbf{q}}, (\mathbf{d} - \tilde{\mathbf{d}}) \times ((\nabla \tilde{\mathbf{v}})_{\text{skw}} (\mathbf{d} - \tilde{\mathbf{d}}) + ((\nabla \mathbf{v})_{\text{skw}} - (\nabla \tilde{\mathbf{v}})_{\text{skw}}) \tilde{\mathbf{d}}))] ds \\ &\quad + \int_0^t (\tilde{\mathbf{d}} \times (\nabla \tilde{\mathbf{v}})_{\text{skw}} (\tilde{\mathbf{d}} - \mathbf{d}), \mathbf{d} \times \mathbf{q} - \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}) ds \\ &\quad + \int_0^t (((\tilde{\mathbf{d}} - \mathbf{d}) \times (\nabla \tilde{\mathbf{v}})_{\text{skw}} + \mathbf{d} \times ((\nabla \tilde{\mathbf{v}})_{\text{skw}} - (\nabla \mathbf{v})_{\text{skw}})) (\tilde{\mathbf{d}} - \mathbf{d}), \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}) ds \\ &= \int_0^t [(\mathbf{d} \times \mathbf{q} - \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}, (\tilde{\mathbf{d}} - \mathbf{d}) \times (\nabla \tilde{\mathbf{v}})_{\text{skw}} (\mathbf{d} - \tilde{\mathbf{d}})) + (\mathbf{d} \times \mathbf{q} - \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}, (\mathbf{d} - \tilde{\mathbf{d}}) \times (\nabla \tilde{\mathbf{v}})_{\text{skw}} \tilde{\mathbf{d}})] ds \\ &\quad + \int_0^t [(\tilde{\mathbf{d}} \times \tilde{\mathbf{q}}, (\mathbf{d} - \tilde{\mathbf{d}}) \times ((\nabla \mathbf{v})_{\text{skw}} - (\nabla \tilde{\mathbf{v}})_{\text{skw}}) \tilde{\mathbf{d}}) + (\tilde{\mathbf{d}} \times (\nabla \tilde{\mathbf{v}})_{\text{skw}} (\mathbf{d} - \tilde{\mathbf{d}}), \mathbf{d} \times \mathbf{q} - \tilde{\mathbf{d}} \times \tilde{\mathbf{q}})] ds \\ &\quad + \int_0^t [2((\mathbf{d} - \tilde{\mathbf{d}}) \times (\nabla \tilde{\mathbf{v}})_{\text{skw}} (\mathbf{d} - \tilde{\mathbf{d}}), \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}) + ((\mathbf{d} - \tilde{\mathbf{d}}) \times ((\nabla \mathbf{v})_{\text{skw}} - (\nabla \tilde{\mathbf{v}})_{\text{skw}}) (\mathbf{d} - \tilde{\mathbf{d}}), \tilde{\mathbf{d}} \times \tilde{\mathbf{q}})] ds \\ &\quad + \int_0^t (\tilde{\mathbf{d}} \times ((\nabla \mathbf{v})_{\text{skw}} - (\nabla \tilde{\mathbf{v}})_{\text{skw}}) (\mathbf{d} - \tilde{\mathbf{d}}), \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}) ds \end{aligned}$$

This can now be estimated by

$$\begin{aligned} I_4 &\leq \delta \int_0^t \|\mathbf{d} \times \mathbf{q} - \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}\|_{L^2}^2 ds + C_\delta \int_0^t \|(\nabla \tilde{\mathbf{v}})_{\text{skw}}\|_{L^3}^2 \left(\|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^{12}}^4 + \|\tilde{\mathbf{d}}\|_{L^\infty(L^\infty)}^2 \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2 \right) ds \\ &\quad + \delta \int_0^t \|(\nabla \mathbf{v})_{\text{skw}} - (\nabla \tilde{\mathbf{v}})_{\text{skw}}\|_{L^2}^2 ds + C_\delta \|\tilde{\mathbf{d}}\|_{L^\infty(L^\infty)}^2 \int_0^t \|\tilde{\mathbf{q}}\|_{L^3}^2 \left(\|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^{12}}^4 + \|\tilde{\mathbf{d}}\|_{L^\infty(L^\infty)}^2 \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2 \right) ds \\ &\quad + 2\|\tilde{\mathbf{d}}\|_{L^\infty(L^\infty)} \int_0^t \|(\nabla \tilde{\mathbf{v}})_{\text{skw}}\|_{L^3} \|\tilde{\mathbf{q}}\|_{L^3} \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2 ds \\ &\leq \delta \int_0^t \mathcal{W} ds + \int_0^t \mathcal{K}(s) \mathcal{E}(s) ds. \end{aligned}$$

The term I_5 can be rearranged as

$$\begin{aligned} I_5 &= \int_0^t [(\tilde{\mathbf{d}} \times ((\nabla \mathbf{v})_{\text{sym}} - (\nabla \tilde{\mathbf{v}})_{\text{sym}}) \tilde{\mathbf{d}}, \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}) + (\mathbf{d} \times ((\nabla \tilde{\mathbf{v}})_{\text{sym}} - (\nabla \mathbf{v})_{\text{sym}}) \mathbf{d}, \mathbf{d} \times \mathbf{q})] \, ds \\ &= \int_0^t ((\tilde{\mathbf{d}} \times (\nabla \mathbf{v})_{\text{sym}} \tilde{\mathbf{d}} - \mathbf{d} \times (\nabla \mathbf{v})_{\text{sym}} \mathbf{d}, \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}) + (\mathbf{d} \times (\nabla \tilde{\mathbf{v}})_{\text{sym}} \mathbf{d} - \tilde{\mathbf{d}} \times (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, \mathbf{d} \times \mathbf{q})) \, ds \\ &\quad - \int_0^t (\mathbf{d} \times (\nabla \mathbf{v})_{\text{sym}} \mathbf{d} - \tilde{\mathbf{d}} \times (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, \mathbf{d} \times \mathbf{q} - \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}) \, ds. \end{aligned} \quad (6.6)$$

Using the simple reformulation

$$\mathbf{d} \times (\nabla \mathbf{v})_{\text{sym}} \mathbf{d} - \tilde{\mathbf{d}} \times (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}} = (\mathbf{d} - \tilde{\mathbf{d}}) \times (\nabla \mathbf{v})_{\text{sym}} (\mathbf{d} - \tilde{\mathbf{d}}) + (\mathbf{d} - \tilde{\mathbf{d}}) \times (\nabla \mathbf{v})_{\text{sym}} \tilde{\mathbf{d}} + \tilde{\mathbf{d}} \times (\nabla \mathbf{v})_{\text{sym}} (\mathbf{d} - \tilde{\mathbf{d}}),$$

the first line of the right-hand side of (6.6) can be rearranged to

$$\begin{aligned} &\int_0^t ((\tilde{\mathbf{d}} \times (\nabla \mathbf{v})_{\text{sym}} \tilde{\mathbf{d}} - \mathbf{d} \times (\nabla \mathbf{v})_{\text{sym}} \mathbf{d}, \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}) + (\mathbf{d} \times (\nabla \tilde{\mathbf{v}})_{\text{sym}} \mathbf{d} - \tilde{\mathbf{d}} \times (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, \mathbf{d} \times \mathbf{q})) \, ds \\ &= \int_0^t ((\mathbf{d} - \tilde{\mathbf{d}}) \times ((\nabla \mathbf{v})_{\text{sym}} - (\nabla \tilde{\mathbf{v}})_{\text{sym}}) (\mathbf{d} - \tilde{\mathbf{d}}) + (\mathbf{d} - \tilde{\mathbf{d}}) \times ((\nabla \mathbf{v})_{\text{sym}} - (\nabla \tilde{\mathbf{v}})_{\text{sym}}) \tilde{\mathbf{d}}, \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}) \, ds \\ &\quad + \int_0^t (\tilde{\mathbf{d}} \times ((\nabla \mathbf{v})_{\text{sym}} - (\nabla \tilde{\mathbf{v}})_{\text{sym}}) (\mathbf{d} - \tilde{\mathbf{d}}), \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}) \, ds \\ &\quad + \int_0^t ((\mathbf{d} - \tilde{\mathbf{d}}) \times (\nabla \tilde{\mathbf{v}})_{\text{sym}} (\mathbf{d} - \tilde{\mathbf{d}}) + (\mathbf{d} - \tilde{\mathbf{d}}) \times (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}} + \tilde{\mathbf{d}} \times (\nabla \mathbf{v})_{\text{sym}} (\mathbf{d} - \tilde{\mathbf{d}}), \mathbf{d} \times \mathbf{q} - \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}) \, ds \end{aligned}$$

and estimated by

$$\begin{aligned} &\int_0^t [(\tilde{\mathbf{d}} \times (\nabla \mathbf{v})_{\text{sym}} \tilde{\mathbf{d}} - \mathbf{d} \times (\nabla \mathbf{v})_{\text{sym}} \mathbf{d}, \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}) + (\mathbf{d} \times (\nabla \tilde{\mathbf{v}})_{\text{sym}} \mathbf{d} - \tilde{\mathbf{d}} \times (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}, \mathbf{d} \times \mathbf{q})] \, ds \\ &\leq \delta \int_0^t \|(\nabla \mathbf{v})_{\text{sym}} - (\nabla \tilde{\mathbf{v}})_{\text{sym}}\|_{L^2}^2 \, ds + \delta \int_0^t \|\mathbf{d} \times \mathbf{q} - \tilde{\mathbf{d}} \times \tilde{\mathbf{q}}\|_{L^2}^2 \, ds \\ &\quad + C_\delta \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^2 \int_0^t (\|\tilde{\mathbf{q}}\|_{L^3}^2 + \|(\nabla \tilde{\mathbf{v}})_{\text{sym}}\|_{L^3}^2) \left(\|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^{12}}^4 + \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^2 \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2 \right) \, ds. \end{aligned}$$

For I_5 and I_6 together we have

$$\begin{aligned} -(\mu_2 + \mu_3)I_5 + \lambda I_6 &\leq ((\mu_2 + \mu_3) - \lambda) \int_0^t (\tilde{\mathbf{d}} \times (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}} - \mathbf{d} \times (\nabla \mathbf{v})_{\text{sym}} \mathbf{d}, \tilde{\mathbf{d}} \times \tilde{\mathbf{q}} - \mathbf{d} \times \mathbf{q}) \, ds \\ &\quad + \delta \int_0^t \mathcal{W}(s) \, ds + \int_0^t \mathcal{K}(s) \mathcal{E}(s) \, ds. \end{aligned}$$

For the term I_8 we observe (see Section 1)

$$\begin{aligned} &\int_0^t [(\mathbf{d} \times (\nabla \mathbf{v})_{\text{sym}} \mathbf{d}, \mathbf{d} \times (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}) + (\tilde{\mathbf{d}} \times (\nabla \mathbf{v})_{\text{sym}} \tilde{\mathbf{d}}, \tilde{\mathbf{d}} \times (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}})] \, ds \\ &= \int_0^t [((\nabla \mathbf{v})_{\text{sym}} \mathbf{d}, (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}) - (\mathbf{d} \cdot (\nabla \mathbf{v})_{\text{sym}} \mathbf{d}, \mathbf{d} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}})] \, ds \\ &\quad + \int_0^t [((\nabla \mathbf{v})_{\text{sym}} \tilde{\mathbf{d}}, (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}}) - (\tilde{\mathbf{d}} \cdot (\nabla \mathbf{v})_{\text{sym}} \tilde{\mathbf{d}}, \tilde{\mathbf{d}} \cdot (\nabla \tilde{\mathbf{v}})_{\text{sym}} \tilde{\mathbf{d}})] \, ds \end{aligned}$$

Comparing the terms I_7 and I_8 with I_2 and I_3 , we see

$$-2I_7 - I_8 = (I_2 - I_3).$$

The appearing terms I_7 and I_8 can, hence, be estimated as in the cases of I_2 and I_3 .

The definition of the variational derivative \mathbf{q} (2.9) and the derivative $\partial F / \partial \nabla \mathbf{d}$ of the free energy potential inserted in the term

I_9 gives

$$\begin{aligned}
I_9 &= \int_0^t [\langle\langle v_s, S : \Lambda : (\Upsilon : ([\tilde{\mathbf{d}}]_X(\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} \otimes S)) \rangle\rangle + (([\mathbf{d}]_X^T \nabla([\tilde{\mathbf{d}}]_X^T(\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}})) ; \Lambda : \nabla \mathbf{d})] ds \\
&\quad + \int_0^t [(\mathbf{d} \times (\mathbf{v} \cdot \nabla) \mathbf{d}, -\tilde{\mathbf{d}} \times \Delta_\Lambda \tilde{\mathbf{d}}) - (\nabla \mathbf{v}; \nabla \tilde{\mathbf{d}}^T \Lambda : \nabla \tilde{\mathbf{d}}) - \langle\langle v_s, (S \nabla \tilde{\mathbf{v}}) : \Lambda : S \rangle\rangle] ds \\
&\quad + \int_0^t [([\mathbf{d}]_X^T \nabla([\tilde{\mathbf{d}}]_X(\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}}) \otimes \mathbf{d} ; \Theta : \nabla \mathbf{d} \otimes \mathbf{d}) + ([\mathbf{d}]_X(\mathbf{v} \cdot \nabla) \mathbf{d}, -[\tilde{\mathbf{d}}]_X \nabla \cdot (\tilde{\mathbf{d}} \cdot \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}))] ds \\
&\quad + \int_0^t \langle\langle v_s, \Upsilon : ([\mathbf{d}]_X(\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} \otimes S) \otimes \mathbf{h} : \Theta : S \otimes \mathbf{h} \rangle\rangle ds \\
&\quad + \int_0^t [(\nabla \tilde{\mathbf{d}} \otimes ([\tilde{\mathbf{d}}]_X^T \mathbf{d}(\mathbf{v} \cdot \nabla) \mathbf{d}) ; \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \langle\langle v_s, S \otimes ([\mathbf{h}]_X^T \tilde{\mathbf{d}}_X(\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}}) : \Theta : S \otimes \mathbf{h} \rangle\rangle] ds \\
&\quad - \int_0^t \langle\langle v_s, (S \nabla \tilde{\mathbf{v}}) \otimes \mathbf{h} : \Theta : S \otimes \mathbf{h} \rangle\rangle ds - \int_0^t (\nabla \mathbf{v}; \nabla \tilde{\mathbf{d}}^T (\tilde{\mathbf{d}} \cdot \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) ds \\
&= J_1 + J_2.
\end{aligned}$$

All terms involving the tensor Λ , i.e. the first two lines, are denoted by the abbreviation J_1 . The terms involving the tensor Θ are denoted by the abbreviation J_2 . In the following, a vector identity is of need. Therefore, we consider a sequence $\{\mathbf{d}_n\} \subset L^2(0, T; \mathbf{H}^2)$ approximating the function \mathbf{d} such that $\mathbf{d}_n \rightharpoonup \mathbf{d}$ in $L^2(0, T; \mathbf{H}^1)$. The sequence $\{\mathbf{d}_n\}$ can be chosen in accordance to Definition 2. Since $\tilde{\mathbf{v}}$ is solenoidal, we see

$$\begin{aligned}
0 &= (\nabla \cdot \tilde{\mathbf{v}}, (\nabla \tilde{\mathbf{d}} : \Lambda : \nabla \mathbf{d}_n)) = -((\tilde{\mathbf{v}} \cdot \nabla) \nabla \tilde{\mathbf{d}}; \Lambda \nabla \mathbf{d}) - ((\tilde{\mathbf{v}} \cdot \nabla) \nabla \mathbf{d}; \Lambda \nabla \tilde{\mathbf{d}}) \\
&= (\nabla \tilde{\mathbf{v}}; \nabla \tilde{\mathbf{d}}^T \Lambda : \nabla \mathbf{d}_n) + (\nabla \tilde{\mathbf{v}}; \nabla \mathbf{d}_n^T \Lambda : \nabla \tilde{\mathbf{d}}) \\
&\quad - ((\tilde{\mathbf{v}} \cdot \nabla) \mathbf{d}_n, -\Delta_\Lambda \tilde{\mathbf{d}}) - (\Lambda : \nabla((\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}}); \nabla \mathbf{d}_n) \\
&= (\nabla \tilde{\mathbf{v}}; \nabla \tilde{\mathbf{d}}^T \Lambda : \nabla \mathbf{d}_n) + (\nabla \tilde{\mathbf{v}}; (\nabla \mathbf{d}_n - \nabla \tilde{\mathbf{d}})^T \Lambda : \nabla \tilde{\mathbf{d}}) + (\nabla \tilde{\mathbf{v}}; \nabla \tilde{\mathbf{d}}^T \Lambda : \nabla \tilde{\mathbf{d}}) \\
&\quad - ((\tilde{\mathbf{v}} \cdot \nabla) \mathbf{d}_n, -\Delta_\Lambda \tilde{\mathbf{d}}) - (\Lambda : \nabla((\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}}); \nabla \mathbf{d}_n).
\end{aligned} \tag{6.7}$$

The integration in time is omitted for brevity. Going to the limit in n and using Lemma 3.3 shows

$$\begin{aligned}
0 &= (\nabla \tilde{\mathbf{v}}; \nabla \tilde{\mathbf{d}}^T \Lambda : \nabla \mathbf{d}) + (\nabla \tilde{\mathbf{v}}; (\nabla \mathbf{d} - \nabla \tilde{\mathbf{d}})^T \Lambda : \nabla \tilde{\mathbf{d}}) + (\nabla \tilde{\mathbf{v}}; \nabla \tilde{\mathbf{d}}^T \Lambda : \nabla \tilde{\mathbf{d}}) \\
&\quad - (\mathbf{d} \times (\tilde{\mathbf{v}} \cdot \nabla) \mathbf{d}, -\mathbf{d} \times \Delta_\Lambda \tilde{\mathbf{d}}) - (\nabla([\tilde{\mathbf{d}}]_X^T \tilde{\mathbf{d}}_X(\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}}); \Lambda : \nabla \mathbf{d}).
\end{aligned}$$

Adding this zero to J_1 , we get

$$\begin{aligned}
J_1 + 0 &= - \int_0^t \langle\langle v_s, ((S - \nabla \tilde{\mathbf{d}})^T \Lambda : (S - \nabla \tilde{\mathbf{d}})) : \nabla \tilde{\mathbf{v}} \rangle\rangle ds - \int_0^t (\nabla \mathbf{v} - \nabla \tilde{\mathbf{v}}, \nabla \tilde{\mathbf{d}}^T \Lambda : \nabla \tilde{\mathbf{d}}) ds \\
&\quad + \int_0^t [(\mathbf{d} \times (\mathbf{v} \cdot \nabla) \mathbf{d}, -\tilde{\mathbf{d}} \times \Delta_\Lambda \tilde{\mathbf{d}}) - (\mathbf{d} \times (\tilde{\mathbf{v}} \cdot \nabla) \mathbf{d}, -\mathbf{d} \times \Delta_\Lambda \tilde{\mathbf{d}})] ds \\
&\quad + \int_0^t [\langle\langle v_t, S : \Lambda : (\Upsilon ([\tilde{\mathbf{d}}]_X(\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} \otimes (S - \nabla \tilde{\mathbf{d}}))) \rangle\rangle + ([\mathbf{d} - \tilde{\mathbf{d}}]_X^T \nabla([\tilde{\mathbf{d}}]_X^T(\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}}); \Lambda : \nabla \mathbf{d})] ds.
\end{aligned}$$

Since \mathbf{v} and $\tilde{\mathbf{v}}$ are solenoidal, an integration-by-parts and Lemma 3.3 gives

$$\begin{aligned}
(\nabla \mathbf{v} - \nabla \tilde{\mathbf{v}}, \nabla \tilde{\mathbf{d}}^T \Lambda : \nabla \tilde{\mathbf{d}}) &= -(\mathbf{v} - \tilde{\mathbf{v}}, \nabla \tilde{\mathbf{d}}^T \Delta_\Lambda \tilde{\mathbf{d}}) - \frac{1}{2} (\mathbf{v} - \tilde{\mathbf{v}}, \nabla (\nabla \tilde{\mathbf{d}} : \Lambda : \nabla \tilde{\mathbf{d}})) \\
&= -([\tilde{\mathbf{d}}]_X^T \tilde{\mathbf{d}}_X ((\mathbf{v} - \tilde{\mathbf{v}}) \cdot \nabla) \tilde{\mathbf{d}}, -\Delta_\Lambda \tilde{\mathbf{d}}).
\end{aligned}$$

After some additional rearrangements, we see

$$\begin{aligned}
J_1 + 0 = & - \int_0^t \langle\langle v_s, (\mathbf{S} - \nabla \tilde{\mathbf{d}})^T \boldsymbol{\Lambda} : (\mathbf{S} - \nabla \tilde{\mathbf{d}}) : \nabla \tilde{\mathbf{v}} \rangle\rangle \, ds \\
& + \int_0^t ((\mathbf{d} - \tilde{\mathbf{d}}) \times \nabla \tilde{\mathbf{d}} (\mathbf{v} - \tilde{\mathbf{v}}), -\tilde{\mathbf{d}} \times \Delta_{\boldsymbol{\Lambda}} \tilde{\mathbf{d}}) \, ds \\
& + \int_0^t ((\mathbf{d} - \tilde{\mathbf{d}}) \times ((\nabla \mathbf{d} - \nabla \tilde{\mathbf{d}})(\mathbf{v} - \tilde{\mathbf{v}})), -\tilde{\mathbf{d}} \times \Delta_{\boldsymbol{\Lambda}} \tilde{\mathbf{d}}) \, ds \\
& + \int_0^t (\tilde{\mathbf{d}} \times ((\nabla \mathbf{d} - \nabla \tilde{\mathbf{d}})(\mathbf{v} - \tilde{\mathbf{v}})), -\tilde{\mathbf{d}} \times \Delta_{\boldsymbol{\Lambda}} \tilde{\mathbf{d}}) \, ds \\
& + \int_0^t \langle\langle v_s, (\mathbf{S} - \nabla \tilde{\mathbf{d}}) : \boldsymbol{\Lambda} : (\mathbf{r} : ([\tilde{\mathbf{d}}]_X (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} \otimes (\mathbf{S} - \nabla \tilde{\mathbf{d}}))) \rangle\rangle \, ds \\
& + \int_0^t ([\mathbf{d} - \tilde{\mathbf{d}}]_X^T \nabla ([\tilde{\mathbf{d}}]_X^T (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}}); \boldsymbol{\Lambda} : (\nabla \mathbf{d} - \nabla \tilde{\mathbf{d}})) \, ds \\
& + \int_0^t ((\mathbf{d} - \tilde{\mathbf{d}}) \times (\nabla \mathbf{d} - \nabla \tilde{\mathbf{d}}) \tilde{\mathbf{v}}, (\mathbf{d} - \tilde{\mathbf{d}}) \times \Delta_{\boldsymbol{\Lambda}} \tilde{\mathbf{d}}) \, ds \\
& + \int_0^t [(\tilde{\mathbf{d}} \times (\nabla \mathbf{d} - \nabla \tilde{\mathbf{d}}) \tilde{\mathbf{v}}, (\mathbf{d} - \tilde{\mathbf{d}}) \times \Delta_{\boldsymbol{\Lambda}} \tilde{\mathbf{d}}) - ((\mathbf{d} - \tilde{\mathbf{d}}) \times (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}}, (\tilde{\mathbf{d}} - \mathbf{d}) \times \Delta_{\boldsymbol{\Lambda}} \tilde{\mathbf{d}})] \, ds
\end{aligned}$$

Thus, the term J_1 can be estimated by

$$\begin{aligned}
J_1 \leq & c \int_0^t \|\nabla \tilde{\mathbf{v}}\|_{L^\infty} \langle\langle v_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 \rangle\rangle \, ds + \delta \int_0^t \|\mathbf{v} - \tilde{\mathbf{v}}\|_{L^6}^2 \, ds \\
& + C_\delta \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^2 \int_0^t \|\tilde{\mathbf{d}}\|_{W^{2,3}}^2 \left(\|\tilde{\mathbf{d}}\|_{L^\infty(W^{1,3})}^2 \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2 + \langle\langle v_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 |h - \tilde{\mathbf{d}}|^2 \rangle\rangle \right) \, ds \\
& + C_\delta \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^4 \int_0^t \|\tilde{\mathbf{d}}\|_{W^{2,3}}^2 \langle\langle v_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 \rangle\rangle \, ds \\
& + c \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)} \int_0^t \|\tilde{\mathbf{v}}\|_{L^\infty} \|\tilde{\mathbf{d}}\|_{W^{1,\infty}} \langle\langle v_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 \rangle\rangle \, ds \\
& + c \int_0^t \|\tilde{\mathbf{v}}\|_{L^\infty} (\|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)} \|\tilde{\mathbf{d}}\|_{W^{2,3}} + \|\tilde{\mathbf{d}}\|_{W^{1,6}}^2) (\langle\langle v_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 \rangle\rangle + \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2) \, ds \\
& + c \int_0^t \|\tilde{\mathbf{v}}\|_{W^{1,3}} \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)} \|\tilde{\mathbf{d}}\|_{W^{1,\infty}} (\langle\langle v_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 \rangle\rangle + \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2) \, ds \\
& + c \int_0^t \|\tilde{\mathbf{v}}\|_{L^\infty} \|\tilde{\mathbf{d}}\|_{W^{2,3}}^4 (\|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^{12}}^4 + \langle\langle v_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 \rangle\rangle) \, ds \\
& + c \int_0^t \|\tilde{\mathbf{v}}\|_{L^\infty} \|\tilde{\mathbf{d}}\|_{W^{2,3}} \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)} (\langle\langle v_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 \rangle\rangle + \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2) \, ds \\
& + c \int_0^t \|\tilde{\mathbf{v}}\|_{L^\infty} \|\tilde{\mathbf{d}}\|_{W^{2,3}} \|\tilde{\mathbf{d}}\|_{L^\infty(W^{1,3})} \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2 \, ds.
\end{aligned}$$

In order to estimate the term J_2 , there are additional vector identities of need. The function \mathbf{d} associated to the generalized gradient Young measure (V^o, m_t, V^∞) can be approximated by a sequence of functions $\{\mathbf{d}_n\} \subset L^\infty(0, T; \mathbf{H}^2)$ in the sense of generalized Young measures (see Definition 2). To show the identity, the integration is again left out.

Since $\tilde{\mathbf{v}}$ is solenoidal, there holds

$$\begin{aligned}
0 &= (\nabla \cdot \tilde{\mathbf{v}}, \nabla \mathbf{d}_n \otimes \mathbf{d}_n : \boldsymbol{\Theta} : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \\
&= -((\tilde{\mathbf{v}} \cdot \nabla) \nabla \mathbf{d}_n \otimes \mathbf{d}_n : \boldsymbol{\Theta} : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) - (\nabla \mathbf{d}_n \otimes (\tilde{\mathbf{v}} \cdot \nabla) \mathbf{d}_n : \boldsymbol{\Theta} : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \\
&\quad - (\nabla \mathbf{d}_n \otimes \mathbf{d}_n : \boldsymbol{\Theta} : (\tilde{\mathbf{v}} \cdot \nabla) \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) - (\nabla \mathbf{d}_n \otimes \mathbf{d}_n : \boldsymbol{\Theta} : \nabla \tilde{\mathbf{d}} \otimes (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}}) \\
&= -((\tilde{\mathbf{v}} \cdot \nabla) \mathbf{d}_n, -\nabla \cdot (\mathbf{d}_n \cdot \boldsymbol{\Theta} : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \mathbf{d}_n : \boldsymbol{\Theta} : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \\
&\quad - (\nabla ((\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}}) \otimes \tilde{\mathbf{d}} : \boldsymbol{\Theta} : \nabla \mathbf{d}_n \otimes \mathbf{d}_n) - (\nabla \tilde{\mathbf{d}} \otimes (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} : \boldsymbol{\Theta} : \nabla \mathbf{d}_n \otimes \mathbf{d}_n) \\
&\quad + (\nabla \tilde{\mathbf{v}}; \nabla \mathbf{d}_n^T (\mathbf{d}_n \cdot \boldsymbol{\Theta} : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) + (\nabla \tilde{\mathbf{v}}; \nabla \tilde{\mathbf{d}}^T (\tilde{\mathbf{d}} \cdot \boldsymbol{\Theta} : \nabla \mathbf{d}_n \otimes \mathbf{d}_n)) \\
&= K_{1,n} + K_{2,n} + K_{3,n}.
\end{aligned}$$

The three terms abbreviate the three lines. For the first two lines, going to the limit in the sense of generalized Young

measures shows

$$K_{1,\infty} + K_{2,\infty} = + \langle\langle v_s, (\mathbf{S}\tilde{\mathbf{v}} \otimes \mathbf{S}^T - \mathbf{S} \otimes (\mathbf{S}\tilde{\mathbf{v}})) : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} \rangle\rangle + (\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}; \nabla \cdot (\Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) \\ - (\nabla ((\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}}) \otimes \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) - (\nabla \tilde{\mathbf{d}} \otimes (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}).$$

The term $\nabla \cdot (\Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})$ has to be understood as $\sum_{j=1}^3 \partial_{x_j} (\Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})_{ijk}$. Adding and subtracting the terms

$$((\tilde{\mathbf{v}} \cdot \nabla) \mathbf{d}, -\nabla \cdot (\tilde{\mathbf{d}} \cdot \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + (\nabla ((\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}}); \mathbf{d} \cdot \Theta : \nabla \mathbf{d} \otimes \mathbf{d}) + \langle\langle v_s, \mathbf{S} \otimes (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} : \Theta : \mathbf{S} \otimes \mathbf{h} \rangle\rangle,$$

and some additional rearrangements give

$$K_{1,\infty} + K_{2,\infty} = - ((\tilde{\mathbf{v}} \cdot \nabla) \mathbf{d}, -\nabla \cdot (\tilde{\mathbf{d}} \cdot \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) - (\nabla ((\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}}); \mathbf{d} \cdot \Theta : \nabla \mathbf{d} \otimes \mathbf{d}) \\ - \langle\langle v_s, \mathbf{S} \otimes (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} : \Theta : \mathbf{S} \otimes \mathbf{h} \rangle\rangle \\ - \langle\langle v_s, ((\mathbf{S} - \nabla \tilde{\mathbf{d}}) \tilde{\mathbf{v}} \otimes (\nabla \tilde{\mathbf{d}} - \mathbf{S})^T + (\mathbf{S} - \nabla \tilde{\mathbf{d}}) \otimes ((\mathbf{S} - \nabla \tilde{\mathbf{d}}) \tilde{\mathbf{v}})) : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} \rangle\rangle \\ + ((\nabla \mathbf{d} - \nabla \tilde{\mathbf{d}}) \tilde{\mathbf{v}} \otimes (\mathbf{d} - \tilde{\mathbf{d}}); \nabla \cdot (\Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) \\ - (\nabla ((\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}}); (\mathbf{d} - \tilde{\mathbf{d}}) \cdot \Theta : (\nabla \mathbf{d} \otimes \mathbf{d} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) \\ - \langle\langle v_s, (\mathbf{S} - \nabla \tilde{\mathbf{d}}) \otimes (\nabla \tilde{\mathbf{d}} \tilde{\mathbf{v}}) : \Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \rangle\rangle.$$

We observe that the last four lines can be estimated by

$$-\int_0^t \langle\langle v_s, ((\mathbf{S} - \nabla \tilde{\mathbf{d}}) \tilde{\mathbf{v}} \otimes (\nabla \tilde{\mathbf{d}} - \mathbf{S})^T + (\mathbf{S} - \nabla \tilde{\mathbf{d}}) \otimes ((\mathbf{S} - \nabla \tilde{\mathbf{d}}) \tilde{\mathbf{v}})) : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} \rangle\rangle ds \\ + \int_0^t ((\nabla \mathbf{d} - \nabla \tilde{\mathbf{d}}) \tilde{\mathbf{v}} \otimes (\mathbf{d} - \tilde{\mathbf{d}}); \nabla \cdot (\Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) ds - \int_0^t (\nabla ((\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}}) \otimes (\mathbf{d} - \tilde{\mathbf{d}}); \Theta : (\nabla \mathbf{d} \otimes \mathbf{d} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) ds \\ - \int_0^t \langle\langle v_s, (\mathbf{S} - \nabla \tilde{\mathbf{d}}) \otimes (\nabla \tilde{\mathbf{d}} \tilde{\mathbf{v}}) : \Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \rangle\rangle ds \\ \leq c \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)} \int_0^t \|\tilde{\mathbf{v}}\|_{L^\infty} \|\tilde{\mathbf{d}}\|_{W^{1,\infty}} \langle\langle v_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 \rangle\rangle ds \\ + c \int_0^t \|\tilde{\mathbf{v}}\|_{L^\infty} (\|\tilde{\mathbf{d}}\|_{W^{2,3}} \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)} + \|\tilde{\mathbf{d}}\|_{W^{1,6}}^2) (\|\mathbf{d} - \tilde{\mathbf{d}}\|_{\mathbf{L}^6}^2 + \langle\langle v_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 \rangle\rangle) ds \\ + c \int_0^t \|\tilde{\mathbf{v}}\|_{L^\infty} \|\tilde{\mathbf{d}}\|_{W^{2,3}} \left(\|\mathbf{d} - \tilde{\mathbf{d}}\|_{\mathbf{L}^6}^2 + \langle\langle v_s, |\Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})|^2 \rangle\rangle \right) ds \\ + c \int_0^t \|\tilde{\mathbf{v}}\|_{W^{1,3}} \|\tilde{\mathbf{d}}\|_{W^{1,\infty}} \left(\|\mathbf{d} - \tilde{\mathbf{d}}\|_{\mathbf{L}^6}^2 + \langle\langle v_s, |\Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})|^2 \rangle\rangle \right) ds \\ + c \int_0^t \|\tilde{\mathbf{v}}\|_{L^\infty} \|\tilde{\mathbf{d}}\|_{W^{1,\infty}} \left(\langle\langle v_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 \rangle\rangle + \langle\langle v_s, |\Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})|^2 \rangle\rangle \right) ds$$

such that with Lemma 3.3 the terms $K_{1,\infty}$ and $K_{2,\infty}$ can be estimated by

$$\int_0^t (K_{1,\infty} + K_{2,\infty}) ds \leq - \int_0^t [((\tilde{\mathbf{v}} \cdot \nabla) \mathbf{d}, -\nabla \cdot (\tilde{\mathbf{d}} \cdot \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) - (\nabla ((\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}}); \mathbf{d} \cdot \Theta : \nabla \mathbf{d} \otimes \mathbf{d})] ds \\ - \int_0^t \langle\langle v_s, \mathbf{S} \otimes (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} : \Theta : \mathbf{S} \otimes \mathbf{h} \rangle\rangle ds + \int_0^t \mathcal{K}(s) \mathcal{E}(s) ds.$$

The term $K_{3,\infty}$ can be transformed to

$$K_{3,\infty} = (\nabla \tilde{\mathbf{v}}; \nabla \mathbf{d}^T (\mathbf{d} \cdot \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) + \left(\nabla \tilde{\mathbf{v}}; \nabla \tilde{\mathbf{d}}^T (\tilde{\mathbf{d}} \cdot \Theta : (\nabla \mathbf{d} \otimes \mathbf{d} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) \right) \\ + \left(\nabla \tilde{\mathbf{v}}; \nabla \tilde{\mathbf{d}}^T (\tilde{\mathbf{d}} \cdot \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \right).$$

Adding the above terms to J_2 and rearranging the terms gives

$$\begin{aligned} J_2 + 0 &= J_2 + K_{1,\infty} + K_{2,\infty} + K_{3,\infty} \\ &\leq \int_0^t (\mathbf{d} \times ((\mathbf{v} - \tilde{\mathbf{v}}) \cdot \nabla) \mathbf{d}, -\tilde{\mathbf{d}} \times \nabla \cdot (\tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \tilde{\mathbf{d}} \times \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \, ds \end{aligned} \quad (6.8a)$$

$$+ \int_0^t (\mathbf{d} \times (\tilde{\mathbf{v}} \cdot \nabla) \mathbf{d}, -(\tilde{\mathbf{d}} - \mathbf{d}) \times \nabla \cdot (\tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \, ds \quad (6.8b)$$

$$- \int_0^t \langle \mathbf{v}_s, (\mathbf{S} - \nabla \tilde{\mathbf{d}}) \nabla \tilde{\mathbf{v}} \otimes (\mathbf{h} - \tilde{\mathbf{d}}) : \Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \rangle \, ds \quad (6.8c)$$

$$- \int_0^t \langle \mathbf{v}_s, (\mathbf{S} - \nabla \tilde{\mathbf{d}}) \nabla \tilde{\mathbf{v}} \otimes \tilde{\mathbf{d}} : \Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \rangle \, ds \quad (6.8d)$$

$$+ \int_0^t (([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X)^T \nabla ([\tilde{\mathbf{d}}]_X \nabla \tilde{\mathbf{d}} \tilde{\mathbf{v}}) \otimes \mathbf{d} : \Theta : \nabla \mathbf{d} \otimes \mathbf{d}) \, ds \quad (6.8e)$$

$$+ \int_0^t \langle \mathbf{v}_s, \mathbf{Y} : ([\tilde{\mathbf{d}}]_X (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} \otimes (\mathbf{S} - \nabla \tilde{\mathbf{d}})) \otimes \mathbf{h} : \Theta : \mathbf{S} \otimes \mathbf{h} \rangle \, ds \quad (6.8f)$$

$$+ \int_0^t \langle \mathbf{v}_s, \mathbf{S} \otimes (([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X)^T [\tilde{\mathbf{d}}]_X (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} : \Theta : \mathbf{S} \otimes \mathbf{h}) \rangle \, ds \quad (6.8g)$$

$$- \int_0^t (\nabla \mathbf{v} - \nabla \tilde{\mathbf{v}}; \nabla \tilde{\mathbf{d}}^T (\tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) \, ds + \int_0^t \mathcal{K}(s) \mathcal{E}(s) \, ds. \quad (6.8h)$$

Since \mathbf{v} and $\tilde{\mathbf{v}}$ are solenoidal and due to Lemma 3.3, there holds

$$\begin{aligned} (\nabla \mathbf{v} - \nabla \tilde{\mathbf{v}}; \nabla \tilde{\mathbf{d}}^T (\tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) &= (((\mathbf{v} - \tilde{\mathbf{v}}) \cdot \nabla) \tilde{\mathbf{d}}, -\nabla \cdot (\tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \\ &\quad - (\mathbf{v} - \tilde{\mathbf{v}}, \nabla (\nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) \\ &= ([\tilde{\mathbf{d}}]_X^T [\tilde{\mathbf{d}}]_X ((\mathbf{v} - \tilde{\mathbf{v}}) \cdot \nabla) \tilde{\mathbf{d}}, -\nabla \cdot (\tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}). \end{aligned}$$

The first (6.8a) and the last line (6.8h) of (6.8) can be transformed to

$$\begin{aligned} &\int_0^t ([\tilde{\mathbf{d}}]_X^T [\mathbf{d}]_X ((\mathbf{v} - \tilde{\mathbf{v}}) \cdot \nabla) \mathbf{d}, -\nabla \cdot (\tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \, ds \\ &- \int_0^t (\nabla \mathbf{v} - \nabla \tilde{\mathbf{v}}; \nabla \tilde{\mathbf{d}}^T (\tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) \, ds \\ &= \int_0^t ([\tilde{\mathbf{d}}]_X^T ([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X) ((\mathbf{v} - \tilde{\mathbf{v}}) \cdot \nabla) (\mathbf{d} - \tilde{\mathbf{d}}), -\nabla \cdot (\tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \, ds \\ &+ \int_0^t ([\tilde{\mathbf{d}}]_X^T [\tilde{\mathbf{d}}]_X ((\mathbf{v} - \tilde{\mathbf{v}}) \cdot \nabla) (\mathbf{d} - \tilde{\mathbf{d}}), -\nabla \cdot (\tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \, ds \\ &+ \int_0^t ([\tilde{\mathbf{d}}]_X^T ([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X) ((\mathbf{v} - \tilde{\mathbf{v}}) \cdot \nabla) \tilde{\mathbf{d}}, -\nabla \cdot (\tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \, ds, \end{aligned}$$

and estimated by

$$\begin{aligned} &\int_0^t ([\tilde{\mathbf{d}}]_X^T [\mathbf{d}]_X ((\mathbf{v} - \tilde{\mathbf{v}}) \cdot \nabla) \mathbf{d}, -\nabla \cdot (\tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \, ds \\ &- \int_0^t (\nabla \mathbf{v} - \nabla \tilde{\mathbf{v}}; \nabla \tilde{\mathbf{d}}^T (\tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) \, ds \\ &\leq \delta \int_0^t \|\mathbf{v} - \tilde{\mathbf{v}}\|_{L^6}^2 \, ds + C_\delta \int_0^t \left(\|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^3 \|\mathbf{d}\|_{W^{2,3}} + \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^2 \|\tilde{\mathbf{d}}\|_{W^{1,6}}^2 \right)^2 \langle \mathbf{v}_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 |\mathbf{h} - \tilde{\mathbf{d}}|^2 \rangle \, ds \\ &+ C_\delta \int_0^t \left(\|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^4 \|\mathbf{d}\|_{W^{2,3}} + \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^3 \|\tilde{\mathbf{d}}\|_{W^{1,6}}^2 \right)^2 \langle \mathbf{v}_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 \rangle \, ds \\ &+ C_\delta \int_0^t \|\tilde{\mathbf{d}}\|_{L^\infty(W^{1,3})} \left(\|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^3 \|\mathbf{d}\|_{W^{2,3}} + \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^2 \|\tilde{\mathbf{d}}\|_{W^{1,6}}^2 \right)^2 \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2 \, ds \end{aligned}$$

The line (6.8c) can be estimated by

$$\begin{aligned} &- \int_0^t \langle \mathbf{v}_s, (\mathbf{S} - \nabla \tilde{\mathbf{d}}) \nabla \tilde{\mathbf{v}} \otimes (\mathbf{h} - \tilde{\mathbf{d}}) : \Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \rangle \, ds \\ &\leq c \int_0^t \|\tilde{\mathbf{v}}\|_{W^{1,\infty}} \left(\langle \mathbf{v}_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 |\mathbf{h} - \tilde{\mathbf{d}}|^2 \rangle + \langle \mathbf{v}_s, |\Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})|^2 \rangle \right) \, ds, \end{aligned}$$

and the line (6.8d) by

$$\begin{aligned} & \int_0^t \langle\langle v_s, (\mathbf{S} - \nabla \tilde{\mathbf{d}}) \nabla \tilde{\mathbf{v}} \otimes \tilde{\mathbf{d}} : \Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \rangle\rangle ds \\ & \leq c \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)} \int_0^t \|\tilde{\mathbf{v}}\|_{\mathbf{W}^{1,\infty}} \left(\langle\langle v_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 |\mathbf{h} - \tilde{\mathbf{d}}|^2 \rangle\rangle + \langle\langle v_s, |\Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})|^2 \rangle\rangle \right) ds. \end{aligned}$$

For the line (6.8b) we observe that it can be rearranged to

$$\begin{aligned} & \int_0^t (\mathbf{d} \times (\tilde{\mathbf{v}} \cdot \nabla) \mathbf{d}, -(\tilde{\mathbf{d}} - \mathbf{d}) \times \nabla \cdot (\tilde{\mathbf{d}} \cdot \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) ds \\ & = \int_0^t \left(([\tilde{\mathbf{d}} - \mathbf{d}]_X)^T ([\mathbf{d} - \tilde{\mathbf{d}}]_X) (\nabla \mathbf{d} - \nabla \tilde{\mathbf{d}}) \tilde{\mathbf{v}}, -\nabla \cdot (\tilde{\mathbf{d}} \cdot \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} \right) ds \\ & \quad + \int_0^t \left(([\tilde{\mathbf{d}} - \mathbf{d}]_X)^T [\tilde{\mathbf{d}}]_X (\tilde{\mathbf{v}} \cdot \nabla) (\mathbf{d} - \tilde{\mathbf{d}}), -\nabla \cdot (\tilde{\mathbf{d}} \cdot \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} \right) ds \\ & \quad + \int_0^t \left(([\tilde{\mathbf{d}} - \mathbf{d}]_X)^T ([\mathbf{d} - \tilde{\mathbf{d}}]_X) \nabla \tilde{\mathbf{d}} \tilde{\mathbf{v}}, -\nabla \cdot (\tilde{\mathbf{d}} \cdot \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} \right) ds \\ & \quad + \int_0^t \left(([\tilde{\mathbf{d}} - \mathbf{d}]_X)^T [\tilde{\mathbf{d}}]_X (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}}, -\nabla \cdot (\tilde{\mathbf{d}} \cdot \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} \right) ds. \end{aligned} \tag{6.9}$$

The first three lines on the right-hand side of (6.9) can be estimated by

$$\begin{aligned} & \int_0^t \left(([\tilde{\mathbf{d}}]_X - [\mathbf{d}]_X)^T ([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X) \nabla (\mathbf{d} - \tilde{\mathbf{d}}) \tilde{\mathbf{v}}, -\nabla \cdot (\tilde{\mathbf{d}} \cdot \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} \right) ds \\ & \quad + \int_0^t \left(([\tilde{\mathbf{d}}]_X - [\mathbf{d}]_X)^T [\tilde{\mathbf{d}}]_X \times (\tilde{\mathbf{v}} \cdot \nabla) (\mathbf{d} - \tilde{\mathbf{d}}), -\nabla \cdot (\tilde{\mathbf{d}} \cdot \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} \right) ds \\ & \quad + \int_0^t \left(([\tilde{\mathbf{d}} - \mathbf{d}]_X)^T ([\mathbf{d} - \tilde{\mathbf{d}}]_X) \nabla \tilde{\mathbf{d}} \tilde{\mathbf{v}}, -\nabla \cdot (\tilde{\mathbf{d}} \cdot \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} \right) ds \\ & \leq c \int_0^t \|\tilde{\mathbf{v}}\|_{L^\infty} \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^2 \left(\langle\langle v_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 |\mathbf{h} - \tilde{\mathbf{d}}|^2 \rangle\rangle + \|\mathbf{d} - \tilde{\mathbf{d}}\|_{\mathbf{L}^6}^2 \right) ds \\ & \quad + c \int_0^t \|\tilde{\mathbf{v}}\|_{L^\infty} \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)} \|\tilde{\mathbf{d}}\|_{\mathbf{W}^{1,6}}^2 \left(\langle\langle v_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 |\mathbf{h} - \tilde{\mathbf{d}}|^2 \rangle\rangle + \|\mathbf{d} - \tilde{\mathbf{d}}\|_{\mathbf{L}^6}^2 \right) ds \\ & \quad + c \int_0^t \|\tilde{\mathbf{v}}\|_{L^\infty} \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^3 \|\tilde{\mathbf{d}}\|_{\mathbf{W}^{2,3}} \left(\langle\langle v_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 \rangle\rangle + \|\mathbf{d} - \tilde{\mathbf{d}}\|_{\mathbf{L}^6}^2 \right) ds \\ & \quad + c \int_0^t \|\tilde{\mathbf{v}}\|_{L^\infty} \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^2 \|\tilde{\mathbf{d}}\|_{\mathbf{W}^{1,6}}^2 \left(\langle\langle v_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 \rangle\rangle + \|\mathbf{d} - \tilde{\mathbf{d}}\|_{\mathbf{L}^6}^2 \right) ds \\ & \quad + c \int_0^t \|\tilde{\mathbf{v}}\|_{L^\infty} \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{W}^{1,3})} \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^2 \|\tilde{\mathbf{d}}\|_{\mathbf{W}^{2,3}} \|\mathbf{d} - \tilde{\mathbf{d}}\|_{\mathbf{L}^6}^2 ds \\ & \quad + c \int_0^t \|\tilde{\mathbf{v}}\|_{L^\infty} \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{W}^{1,3})} \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)} \|\tilde{\mathbf{d}}\|_{\mathbf{W}^{1,6}}^2 \|\mathbf{d} - \tilde{\mathbf{d}}\|_{\mathbf{L}^6}^2 ds. \end{aligned}$$

The last line of (6.9) can be transformed by an integration-by-parts into

$$\begin{aligned} & \int_0^t \left(([\tilde{\mathbf{d}}]_X - [\mathbf{d}]_X)^T [\tilde{\mathbf{d}}]_X (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}}, -\nabla \cdot (\tilde{\mathbf{d}} \cdot \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) + \nabla \tilde{\mathbf{d}} : \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} \right) ds \\ & = \int_0^t \left(([\tilde{\mathbf{d}}]_X - [\mathbf{d}]_X)^T \nabla ([\tilde{\mathbf{d}}]_X \nabla \tilde{\mathbf{d}} \tilde{\mathbf{v}}) \otimes \tilde{\mathbf{d}} ; \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}} \right) ds \end{aligned} \tag{6.10a}$$

$$+ \int_0^t (\Upsilon : ([\tilde{\mathbf{d}}]_X (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} \otimes (\nabla \tilde{\mathbf{d}} - \nabla \mathbf{d})) \otimes \tilde{\mathbf{d}} ; \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) ds \tag{6.10b}$$

$$+ \int_0^t (\nabla \tilde{\mathbf{d}} \otimes ([\tilde{\mathbf{d}}]_X - [\mathbf{d}]_X)^T [\tilde{\mathbf{d}}]_X (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} ; \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) ds. \tag{6.10c}$$

It remains to estimate the three terms of the lines (6.10a)-(6.10c) and the three remaining terms of the lines (6.8e)-(6.8g)

of (6.8). The sum of the lines (6.8e) and (6.10a) can be rearranged and estimated by

$$\begin{aligned}
& \int_0^t (([\tilde{\mathbf{d}}]_X - [\mathbf{d}]_X)^T \nabla ([\tilde{\mathbf{d}}]_X \nabla \tilde{\mathbf{d}} \tilde{\mathbf{v}}) \otimes \tilde{\mathbf{d}} ; \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \, ds + \int_0^t (([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X)^T \nabla ([\tilde{\mathbf{d}}]_X \nabla \tilde{\mathbf{d}} \tilde{\mathbf{v}}) \otimes \mathbf{d} ; \Theta : \nabla \mathbf{d} \otimes \mathbf{d}) \, ds \\
&= \int_0^t (([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X)^T \nabla ([\tilde{\mathbf{d}}]_X \nabla \tilde{\mathbf{d}} \tilde{\mathbf{v}}) \otimes (\mathbf{d} - \tilde{\mathbf{d}}) ; \Theta : (\nabla \mathbf{d} \otimes \mathbf{d} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) \, ds \\
&\quad + \int_0^t (([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X)^T \nabla ([\tilde{\mathbf{d}}]_X \nabla \tilde{\mathbf{d}} \tilde{\mathbf{v}}) \otimes \tilde{\mathbf{d}} ; \Theta : (\nabla \mathbf{d} \otimes \mathbf{d} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})) \, ds \\
&\quad + \int_0^t (([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X)^T \nabla ([\tilde{\mathbf{d}}]_X \nabla \tilde{\mathbf{d}} \tilde{\mathbf{v}}) \otimes (\mathbf{d} - \tilde{\mathbf{d}}) ; \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \, ds \\
&\leq c \int_0^t \|\nabla([\tilde{\mathbf{d}}]_X \nabla \tilde{\mathbf{d}} \tilde{\mathbf{v}})\|_{L^3} \left(\|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^{12}}^4 + \langle\langle v_s, |\Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})|^2 \rangle\rangle \right) \, ds \\
&\quad + c \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)} \int_0^t \|\nabla([\tilde{\mathbf{d}}]_X \nabla \tilde{\mathbf{d}} \tilde{\mathbf{v}})\|_{L^3} \left(\|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2 + \langle\langle v_s, |\Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})|^2 \rangle\rangle \right) \, ds \\
&\quad + c \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)} \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{W}^{1,3})} \int_0^t \|\nabla([\tilde{\mathbf{d}}]_X \nabla \tilde{\mathbf{d}} \tilde{\mathbf{v}})\|_{L^3} \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2 \, ds.
\end{aligned}$$

We see that the norm on the right hand side can be bounded by

$$\|\nabla([\tilde{\mathbf{d}}]_X \nabla \tilde{\mathbf{d}} \tilde{\mathbf{v}})\|_{L^3} \leq \|\tilde{\mathbf{v}}\|_{L^\infty} \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)} \|\tilde{\mathbf{d}}\|_{W^{2,3}} + \|\tilde{\mathbf{v}}\|_{L^\infty} \|\tilde{\mathbf{d}}\|_{W^{1,6}}^2 + \|\tilde{\mathbf{v}}\|_{W^{1,3}} \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)} \|\tilde{\mathbf{d}}\|_{W^{1,\infty}}$$

Regarding (6.8f) and (6.10b), we observe

$$\begin{aligned}
& \int_0^t \langle\langle v_s, \Upsilon : ([\tilde{\mathbf{d}}]_X (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} \otimes (\mathbf{S} - \nabla \tilde{\mathbf{d}})) \otimes \mathbf{h} ; \Theta : \mathbf{S} \otimes \mathbf{h} \rangle\rangle \, ds \\
&+ \int_0^t ((\Upsilon : ([\tilde{\mathbf{d}}]_X (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} \otimes (\nabla \tilde{\mathbf{d}} - \nabla \mathbf{d}))) \otimes \tilde{\mathbf{d}} ; \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \, ds \\
&= \int_0^t \langle\langle v_s, (\Upsilon : ([\tilde{\mathbf{d}}]_X (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} \otimes (\mathbf{S} - \nabla \tilde{\mathbf{d}}))) \otimes (\mathbf{h} - \tilde{\mathbf{d}}) ; \Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \rangle\rangle \, ds \\
&\quad + \int_0^t \langle\langle v_s, (\Upsilon : ([\tilde{\mathbf{d}}]_X (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} \otimes (\mathbf{S} - \nabla \tilde{\mathbf{d}}))) \otimes \tilde{\mathbf{d}} ; \Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \rangle\rangle \, ds \\
&\quad + \int_0^t ((\Upsilon : ([\tilde{\mathbf{d}}]_X (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}}) \otimes (\nabla \tilde{\mathbf{d}} - \nabla \mathbf{d})) \otimes (\tilde{\mathbf{d}} - \mathbf{d}) ; \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \, ds.
\end{aligned}$$

Thus, the right-hand side can be estimated by

$$\begin{aligned}
& \int_0^t \langle\langle v_s, \Upsilon : ([\tilde{\mathbf{d}}]_X (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} \otimes (\mathbf{S} - \nabla \tilde{\mathbf{d}})) \otimes \mathbf{h} ; \Theta : \mathbf{S} \otimes \mathbf{h} \rangle\rangle \, ds \\
&+ \int_0^t (\Upsilon : ([\tilde{\mathbf{d}}]_X (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} \otimes (\nabla \tilde{\mathbf{d}} - \nabla \mathbf{d})) \otimes \tilde{\mathbf{d}} ; \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \, ds \\
&\leq \int_0^t \|\tilde{\mathbf{v}}\|_{L^\infty} \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)} \|\tilde{\mathbf{d}}\|_{W^{1,\infty}} \left(\langle\langle v_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 |\mathbf{h} - \tilde{\mathbf{d}}|^2 \rangle\rangle + \langle\langle v_s, |\Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})|^2 \rangle\rangle \right) \, ds \\
&\quad + \int_0^t \|\tilde{\mathbf{v}}\|_{L^\infty} \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^2 \|\tilde{\mathbf{d}}\|_{W^{1,\infty}} \left(\langle\langle v_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 \rangle\rangle + \langle\langle v_s, |\Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}})|^2 \rangle\rangle \right) \, ds \\
&\quad + \int_0^t \|\tilde{\mathbf{v}}\|_{L^\infty} \|\tilde{\mathbf{d}}\|_{L^\infty(\mathbf{L}^\infty)}^3 \|\tilde{\mathbf{d}}\|_{W^{1,3}} \left(\langle\langle v_s, |\mathbf{S} - \nabla \tilde{\mathbf{d}}|^2 \rangle\rangle + \|\mathbf{d} - \tilde{\mathbf{d}}\|_{L^6}^2 \right) \, ds
\end{aligned}$$

Regarding the lines (6.8g) and (6.10c), we observe

$$\begin{aligned}
& \int_0^t \langle\langle v_s, \mathbf{S} \otimes (([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X)^T [\tilde{\mathbf{d}}]_X (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} : \Theta : \mathbf{S} \otimes \mathbf{h}) \rangle\rangle \, ds + \int_0^t (\nabla \tilde{\mathbf{d}} \otimes ([\tilde{\mathbf{d}}]_X - [\mathbf{d}]_X)^T [\tilde{\mathbf{d}}]_X (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}} ; \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \, ds \\
&= \int_0^t \langle\langle v_s, (\mathbf{S} - \nabla \tilde{\mathbf{d}}) \otimes (([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X)^T [\tilde{\mathbf{d}}]_X (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}}) : \Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \rangle\rangle \, ds \\
&\quad + \int_0^t \langle\langle v_s, \nabla \tilde{\mathbf{d}} \otimes (([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X)^T [\tilde{\mathbf{d}}]_X (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}}) : \Theta : (\mathbf{S} \otimes \mathbf{h} - \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \rangle\rangle \, ds \\
&\quad + \int_0^t ((\nabla \tilde{\mathbf{d}} - \nabla \mathbf{d}) \otimes (([\mathbf{d}]_X - [\tilde{\mathbf{d}}]_X)^T [\tilde{\mathbf{d}}]_X (\tilde{\mathbf{v}} \cdot \nabla) \tilde{\mathbf{d}}) ; \Theta : \nabla \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \, ds
\end{aligned}$$

Thus, the right-hand side can be estimated by

$$\begin{aligned} & \int_0^t \langle\langle v_s, S \otimes (([d]_X - [\tilde{d}]_X)^T [\tilde{d}]_X (\tilde{v} \cdot \nabla) \tilde{d}; \Theta : S \otimes h) \rangle\rangle ds + \int_0^t \left(\nabla \tilde{d} \otimes ([\tilde{d}]_X - [d]_X)^T [\tilde{d}]_X (\tilde{v} \cdot \nabla) \tilde{d}; \Theta : \nabla \tilde{d} \otimes \tilde{d} \right) ds \\ & \leq \|\tilde{d}\|_{L^\infty(\mathbf{L}^\infty)} \|\tilde{d}\|_{L^\infty(\mathbf{W}^{1,\infty})} \int_0^t \|\tilde{v}\|_{L^\infty} \left(\langle\langle v_s, |S - \nabla \tilde{d}|^2 |\tilde{h} - \tilde{d}|^2 \rangle\rangle + \langle\langle v_s, |\Theta : (S \otimes h - \nabla \tilde{d} \otimes \tilde{d})|^2 \rangle\rangle \right) ds \\ & \quad + \|\tilde{d}\|_{L^\infty(\mathbf{L}^\infty)} \int_0^t \|\tilde{v}\|_{L^\infty} \|\tilde{d}\|_{\mathbf{W}^{1,6}}^2 \left(\|d - \tilde{d}\|_{\mathbf{L}^6}^2 + \langle\langle v_s, |\Theta : (S \otimes h - \nabla \tilde{d} \otimes \tilde{d})|^2 \rangle\rangle \right) ds \\ & \quad + \|\tilde{d}\|_{L^\infty(\mathbf{L}^\infty)}^2 \int_0^t \|\tilde{v}\|_{L^\infty} \|\tilde{d}\|_{\mathbf{W}^{1,6}}^2 (\langle\langle v_s, |S - \nabla \tilde{d}|^2 \rangle\rangle + \|d - \tilde{d}\|_{\mathbf{L}^6}^2) ds \end{aligned}$$

Together we get for the term I_9

$$I_9 \leq \delta \int_0^t \mathcal{W}(s) ds + \int_0^t \mathcal{K}(s) \mathcal{E}(s) ds.$$

At last, the term I_{10} can be estimated by

$$|I_{10}| = \left| \int_0^t \langle\langle \mu_s, \Gamma : (\Gamma \cdot \nabla \tilde{v}) \rangle\rangle dt \right| \leq c \int_0^t \|(\nabla \tilde{v})_{\text{sym}}\|_{\mathbf{L}^\infty} \langle\langle \mu_s, 1 \rangle\rangle ds \leq \int_0^t \|\tilde{v}\|_{\mathbf{W}^{1,\infty}} \mathcal{E}(s) ds$$

Finally, we insert everything back into (6.3). We end up with

$$\begin{aligned} \mathcal{E}(t) + \int_0^t \mathcal{W}(s) ds & \leq \mathcal{E}(0) + \frac{1}{2} \mathcal{E}(t) + \delta \int_0^t \mathcal{W}(s) ds + C_\delta \int_0^t \mathcal{K}(s) \mathcal{E}(s) ds \\ & \quad + ((\nabla d(0) - \nabla \tilde{d}(0)) \otimes (d(0) - \tilde{d}(0))) : \Theta : (\nabla \tilde{d}(0) \otimes \tilde{d}(0)) \\ & \quad + c \|\nabla \tilde{d} \otimes \tilde{d}\|_{L^\infty(\mathbf{L}^\infty)}^2 \|d(0) - \tilde{d}(0)\|_{\mathbf{L}^2}^2 \\ & \quad + ((\mu_2 + \mu_3) - \lambda) \int_0^t (\tilde{d} \times (\nabla \tilde{v})_{\text{sym}} \tilde{d} - d \times (\nabla v)_{\text{sym}} d, \tilde{d} \times \tilde{q} - d \times q) ds \end{aligned}$$

Since the constants are assumed to fulfil the dissipative relation (2.2f) we can find a real number $\zeta \in (0, 1)$ such that

$$(\mu_2 + \mu_3) - \lambda \leq \zeta^2 4(\mu_5 + \mu_6 - \lambda(\mu_2 + \mu_3)). \quad (6.11)$$

we estimate the relative entropy further on with Youngs and Hölders inequality

$$\begin{aligned} \frac{1}{2} \mathcal{E}(t) + \int_0^t \mathcal{W}(s) ds & \leq \mathcal{E}(0) + ((\nabla d(0) - \nabla \tilde{d}(0)) \otimes (d(0) - \tilde{d}(0))) : \Theta : (\nabla \tilde{d}(0) \otimes \tilde{d}(0)) \\ & \quad + c \|d(0) - \tilde{d}(0)\|_{\mathbf{L}^2}^2 + \delta \int_0^t \mathcal{W}(s) ds + \int_0^t \mathcal{K}(s) \mathcal{E}(s) ds \\ & \quad + \zeta \int_0^t \|d \times q - \tilde{d} \times \tilde{q}\|_{\mathbf{L}^2}^2 ds \\ & \quad + \zeta \int_0^t (\mu_5 + \mu_6 - \lambda(\mu_2 + \mu_3)) \|d \times (\nabla v)_{\text{sym}} d - \tilde{d} \times (\nabla \tilde{v})_{\text{sym}} \tilde{d}\|_{\mathbf{L}^2}^2 ds. \end{aligned}$$

The last term on the right hand side can be written as

$$\begin{aligned} & \int_0^t \|d \times (\nabla v)_{\text{sym}} d - \tilde{d} \times (\nabla \tilde{v})_{\text{sym}} \tilde{d}\|_{\mathbf{L}^2}^2 ds \\ & \leq \int_0^t [\|d \times ((\nabla v)_{\text{sym}} d - (\nabla \tilde{v})_{\text{sym}} \tilde{d})\|_{\mathbf{L}^2}^2 + \|(d - \tilde{d}) \times (\nabla \tilde{v})_{\text{sym}} \tilde{d}\|_{\mathbf{L}^2}^2] ds \\ & \leq \int_0^t [\|(\nabla v)_{\text{sym}} d - (\nabla \tilde{v})_{\text{sym}} \tilde{d}\|_{\mathbf{L}^2}^2 - \|d \cdot ((\nabla v)_{\text{sym}} d - (\nabla \tilde{v})_{\text{sym}} \tilde{d})\|_{\mathbf{L}^2}^2] ds \\ & \quad + \|\tilde{d}\|_{L^\infty(\mathbf{L}^\infty)}^2 \int_0^t \|\tilde{v}\|_{\mathbf{W}^{1,3}}^2 \|d - \tilde{d}\|_{\mathbf{L}^6}^2 ds \\ & \leq \int_0^t \|(\nabla v)_{\text{sym}} d - (\nabla \tilde{v})_{\text{sym}} \tilde{d}\|_{\mathbf{L}^2}^2 ds + \|\tilde{d}\|_{L^\infty(\mathbf{L}^\infty)}^2 \int_0^t \|\tilde{v}\|_{\mathbf{W}^{1,3}}^2 \|d - \tilde{d}\|_{\mathbf{L}^6}^2 ds, \end{aligned}$$

such that we get the following estimate

$$\begin{aligned} \frac{1}{2} \mathcal{E}(t) + \int_0^t \mathcal{W}(s) ds & \leq \mathcal{E}(0) + ((\nabla d(0) - \nabla \tilde{d}(0)) \otimes (d(0) - \tilde{d}(0))) : \Theta : (\nabla \tilde{d}(0) \otimes \tilde{d}(0)) \\ & \quad + c \|\nabla \tilde{d} \otimes \tilde{d}\|_{L^\infty(\mathbf{L}^\infty)}^2 \|d(0) - \tilde{d}(0)\|_{\mathbf{L}^2}^2 + (\zeta + \delta) \int_0^t \mathcal{W}(s) ds + \int_0^t \mathcal{K}(s) \mathcal{E}(s) ds. \end{aligned}$$

We now choose δ sufficiently small, such that $\delta \leq (1 - \zeta)$. The assertion of Proposition 6.1 immediately follows from Gronwalls estimate.

□

Acknowledgments

The author gratefully acknowledges financial support from the Deutsche Forschungsgemeinschaft through the Collaborative Research Center 901 "Control of self-organizing nonlinear systems: Theoretical methods and concepts of application"(projects A8).

A Appendix

For $\mathbf{d} \in \mathcal{C}^1(\Omega)$ the terms of the quadratic Oseen–Frank free energy (see Section 2 and (2.4)) depending on $\nabla \mathbf{d}$ can be expressed using a tensor of fourth order:

$$\begin{aligned} |\nabla \mathbf{d}|^2 &= \nabla \mathbf{d} : \boldsymbol{\Lambda}^0 : \nabla \mathbf{d} \quad \text{mit } \boldsymbol{\Lambda}_{ijkl}^0 = \boldsymbol{\delta}_{ik} \boldsymbol{\delta}_{jl}, \\ (\nabla \cdot \mathbf{d})^2 &= \nabla \mathbf{d} : \boldsymbol{\Lambda}^1 : \nabla \mathbf{d} \quad \text{mit } \boldsymbol{\Lambda}_{ijkl}^1 = \boldsymbol{\delta}_{ij} \boldsymbol{\delta}_{kl}, \\ \text{tr}(\nabla \mathbf{d}^2) &= \nabla \mathbf{d} : \boldsymbol{\Lambda}^2 : \nabla \mathbf{d} \quad \text{mit } \boldsymbol{\Lambda}_{ijkl}^2 = \boldsymbol{\delta}_{il} \boldsymbol{\delta}_{jk}, \\ |\nabla \times \mathbf{d}|^2 &= \nabla \mathbf{d} : \boldsymbol{\Lambda}^3 : \nabla \mathbf{d} \quad \text{mit } \boldsymbol{\Lambda}^3 = \boldsymbol{\Lambda}^0 - \boldsymbol{\Lambda}^2. \end{aligned} \tag{A.1}$$

Here $\boldsymbol{\Lambda}^0$ is the identity in $\mathbb{R}^{3 \times 3}$. Similar, we see for the non-quadratic terms in the Oseen–Frank energy (see Section 2 and (2.4))

$$\begin{aligned} |\nabla \mathbf{d}|^2 |\mathbf{d}|^2 &= \nabla \mathbf{d} \otimes \mathbf{d} : \boldsymbol{\Theta}^0 : \nabla \mathbf{d} \otimes \mathbf{d} \quad \text{mit } \boldsymbol{\Theta}_{ijklmn}^0 = \boldsymbol{\delta}_{il} \boldsymbol{\delta}_{jm} \boldsymbol{\delta}_{kn}, \\ (\nabla \cdot \mathbf{d})^2 |\mathbf{d}|^2 &= \nabla \mathbf{d} \otimes \mathbf{d} : \boldsymbol{\Theta}^1 : \nabla \mathbf{d} \otimes \mathbf{d} \quad \text{mit } \boldsymbol{\Theta}_{ijklmn}^1 = \boldsymbol{\delta}_{ij} \boldsymbol{\delta}_{lm} \boldsymbol{\delta}_{kn}, \\ |\nabla \times \mathbf{d}|^2 |\mathbf{d}|^2 &= 2 |(\nabla \mathbf{d})_{\text{skw}}|^2 |\mathbf{d}|^2 \\ &= \nabla \mathbf{d} \otimes \mathbf{d} : \boldsymbol{\Theta}^2 : \nabla \mathbf{d} \otimes \mathbf{d} \quad \text{mit } \boldsymbol{\Theta}_{ijklmn}^2 = \boldsymbol{\delta}_{kn} (\boldsymbol{\delta}_{il} \boldsymbol{\delta}_{jm} - \boldsymbol{\delta}_{im} \boldsymbol{\delta}_{jl}), \\ |\mathbf{d} \cdot \nabla \times \mathbf{d}|^2 &= 4 |(\nabla \mathbf{d})_{\text{skw}} \mathbf{d}|^2 \\ &= \nabla \mathbf{d} \otimes \mathbf{d} : \boldsymbol{\Theta}^3 : \nabla \mathbf{d} \otimes \mathbf{d} \quad \text{mit } \boldsymbol{\Theta}_{ijklmn}^3 = \boldsymbol{\delta}_{il} \boldsymbol{\delta}_{mn} \boldsymbol{\delta}_{jk} - \boldsymbol{\delta}_{mi} \boldsymbol{\delta}_{ln} \boldsymbol{\delta}_{jk} \\ &\quad - \boldsymbol{\delta}_{lj} \boldsymbol{\delta}_{mn} \boldsymbol{\delta}_{ik} + \boldsymbol{\delta}_{jm} \boldsymbol{\delta}_{ln} \boldsymbol{\delta}_{ik}, \\ |\mathbf{d} \cdot \nabla \times \mathbf{d}|^2 &= ([\mathbf{d}]_X : (\nabla \mathbf{d})_{\text{skw}})^2 \\ &= \nabla \mathbf{d} \otimes \mathbf{d} : \boldsymbol{\Theta}_{ijklmn}^4 : \nabla \mathbf{d} \otimes \mathbf{d} \quad \text{mit } \boldsymbol{\Theta}^4 = \boldsymbol{\delta}_{kn} \boldsymbol{\delta}_{jm} \boldsymbol{\delta}_{il} + \boldsymbol{\delta}_{km} \boldsymbol{\delta}_{jl} \boldsymbol{\delta}_{in} + \boldsymbol{\delta}_{kl} \boldsymbol{\delta}_{jn} \boldsymbol{\delta}_{im} \\ &\quad - \boldsymbol{\delta}_{kn} \boldsymbol{\delta}_{jl} \boldsymbol{\delta}_{im} - \boldsymbol{\delta}_{km} \boldsymbol{\delta}_{jn} \boldsymbol{\delta}_{il} - \boldsymbol{\delta}_{kl} \boldsymbol{\delta}_{jm} \boldsymbol{\delta}_{in}. \end{aligned}$$

With this definitions, we prove the algebraic relations stated in Proposition 4.1.

Proposition A.1. For the tensor $\boldsymbol{\Theta} = k_3 \boldsymbol{\Theta}^1 + k_4 \boldsymbol{\Theta}^2 + k_5 \boldsymbol{\Theta}^3$ with $k_i > 0$, for all $i \in \{3, 4, 5\}$ (see (2.7)), there exists a constant $c > 0$ such that

$$|\boldsymbol{\Theta} : (\mathbf{S} \otimes \mathbf{h} - \tilde{\mathbf{S}} \otimes \tilde{\mathbf{h}})|^2 \leq c ((\mathbf{S} \otimes \mathbf{h} - \tilde{\mathbf{S}} \otimes \tilde{\mathbf{h}}) : \boldsymbol{\Theta} : (\mathbf{S} \otimes \mathbf{h} - \tilde{\mathbf{S}} \otimes \tilde{\mathbf{h}})), \quad \text{for all } \mathbf{S} \otimes \mathbf{h} \in \mathbb{R}^{3 \times 3} \times \mathbb{R}^3.$$

Proof. We consider the tensors $\boldsymbol{\Theta}^i$ with $i \in \{1, \dots, 3\}$ multiplied with $\mathbf{S} \otimes \mathbf{h}$

$$\begin{aligned} (\boldsymbol{\Theta}^1 : \mathbf{S} \otimes \mathbf{h})_{ijk} &:= \boldsymbol{\delta}_{ij} \text{tr}(\mathbf{S}) \mathbf{h}_k \\ (\boldsymbol{\Theta}^3 : \mathbf{S} \otimes \mathbf{h})_{ijk} &:= 2 \boldsymbol{\delta}_{jk} ((\mathbf{S})_{\text{skw}} \mathbf{h})_i - \boldsymbol{\delta}_{ik} ((\mathbf{S})_{\text{skw}} \mathbf{h})_j. \end{aligned}$$

The Levi–Civita-tensor is defined in Section 1. It allows to represent the cross product via

$$\mathbf{a} \times \mathbf{b} = \boldsymbol{\Upsilon} : (\mathbf{a} \otimes \mathbf{b}) = \boldsymbol{\Upsilon}_{ijk} \mathbf{a}_j \mathbf{b}_k \quad \text{for all } \mathbf{a}, \mathbf{b} \in \mathbb{R}^d$$

and the curl of a vectorfield via

$$\nabla \times \mathbf{d} = \Upsilon_{ijk} \partial_j \mathbf{d}_k \quad \text{for all } \mathbf{d} \in \mathcal{C}^1(\Omega).$$

Together, we get

$$(\mathbf{d} \cdot \nabla \times \mathbf{d}) = \Upsilon : \mathbf{d} \otimes \nabla \mathbf{d}^T = \Upsilon_{ijk} \mathbf{d}_i \mathbf{d}_{k,j}.$$

The definition of the Levi–Civita-tensor implies

$$\Upsilon_{kji} \Upsilon_{nml} = \delta_{kn} \delta_{jm} \delta_{il} + \delta_{km} \delta_{jl} \delta_{in} + \delta_{kl} \delta_{jn} \delta_{im} - \delta_{kn} \delta_{jl} \delta_{im} - \delta_{km} \delta_{jn} \delta_{il} - \delta_{kl} \delta_{jm} \delta_{in} = \Theta^2$$

and thus

$$(\Theta^2 : \mathbf{S} \otimes \mathbf{h})_{ijk} := \Upsilon_{kji} ((\mathbf{S})_{\text{skw}} : [\mathbf{h}]_X).$$

The different term multiplied with each other gives

$$\begin{aligned} |\Theta^1 : (\mathbf{S} \otimes \mathbf{h} - \tilde{\mathbf{S}} \otimes \tilde{\mathbf{h}})|^2 &= 3|\text{tr}(\mathbf{S})\mathbf{h} - \text{tr}(\tilde{\mathbf{S}})\tilde{\mathbf{h}}|^2, \\ |\Theta^3 : (\mathbf{S} \otimes \mathbf{h} - \tilde{\mathbf{S}} \otimes \tilde{\mathbf{h}})|^2 &= 4 \sum_{ijk=1}^3 \left(\delta_{jk} ((\mathbf{S})_{\text{skw}}\mathbf{h} - (\tilde{\mathbf{S}})_{\text{skw}}\tilde{\mathbf{h}})_i - \delta_{ik} ((\mathbf{S})_{\text{skw}}\mathbf{h} - (\tilde{\mathbf{S}})_{\text{skw}}\tilde{\mathbf{h}})_j \right)^2 \\ &= 4(6|(\mathbf{S})_{\text{skw}}\mathbf{h} - (\tilde{\mathbf{S}})_{\text{skw}}\tilde{\mathbf{h}}|^2 - 2|(\mathbf{S})_{\text{skw}}\mathbf{h} - (\tilde{\mathbf{S}})_{\text{skw}}\tilde{\mathbf{h}}|^2)^2 \\ &= 16|(\mathbf{S})_{\text{skw}}\mathbf{h} - (\tilde{\mathbf{S}})_{\text{skw}}\tilde{\mathbf{h}}|^2, \\ |\Theta^2 : \mathbf{S} \otimes \mathbf{h}|^2 &= \sum_{ijk=1}^3 \Upsilon_{kji} \Upsilon_{kji} ((\mathbf{S})_{\text{skw}} : [\mathbf{h}]_X - (\tilde{\mathbf{S}})_{\text{skw}} : [\tilde{\mathbf{h}}]_X)^2 \\ &= 6((\mathbf{S})_{\text{skw}} : [\mathbf{h}]_X - (\tilde{\mathbf{S}})_{\text{skw}} : [\tilde{\mathbf{h}}]_X)^2. \end{aligned}$$

For the products of different tensors, we see that

$$\begin{aligned} \mathbf{S} \otimes \mathbf{h} : \Theta^1 : \Theta^2 : \mathbf{S} \otimes \mathbf{h} &= \sum_{ijk=1}^3 (\delta_{ij} \text{tr}(\mathbf{S})\mathbf{h}_k) \left(2((\mathbf{S})_{\text{skw}})_{ij} \mathbf{h}_k \right) \\ &= 2\text{tr}(\mathbf{S})(\mathbf{I} : (\mathbf{S})_{\text{skw}})|\mathbf{h}|^2 = 0, \end{aligned}$$

The definition of the Levi–Civita-tensors gives

$$\Upsilon_{ijk} \delta_{ij} = \Upsilon_{ijk} \delta_{jk} = \Upsilon_{ijk} \delta_{ik} = 0.$$

And thus the products

$$\mathbf{S} \otimes \mathbf{h} : \Theta^2 : \Theta^i : \mathbf{S} \otimes \mathbf{h},$$

vanish for all $i \in \{1, 3\}$. This gives the assertion. \square

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