

## *Exercise Sheet 1*

**Exercise 1 (Navier–Stokes equation)** We consider the Navier–Stokes equation

$$\partial_t u + (u \cdot \nabla)u + \nabla p = \Delta u, \quad \operatorname{div} u = 0 \quad \text{for } (t, x) \in (0, T) \times \Omega$$

in the two-dimensional strip  $\Omega = \mathbb{R} \times (0, h)$ . Here  $u(t, x) \in \mathbb{R}^2$  describes the velocity of the fluid particles, and  $p(t, x) \in \mathbb{R}$  is the pressure field. By  $\Delta u$  we denote the componentwise Laplace-Operator, and  $(u \cdot \nabla)v = \sum_1^d u_j \partial_j v$ . Moreover, we impose the boundary conditions  $u = 0$  on the boundary  $\partial\Omega$  of the flow domain  $\Omega$ .

- (a) Show that for all  $k \in \mathbb{N}$  there are solutions of the form  $u(t, x) = e^{\alpha t} \sin(k\pi x_2/h) \Phi$  for suitable  $\alpha \in \mathbb{R}$  and  $\Phi \in \mathbb{R}^2$ .
- (b) Determine all solutions  $(u, p)$ , for which  $u$  does not depend on  $(t, x_1)$ .
- (c) For general solutions show that the mass flux in  $x_1$  direction is independent of  $x_1$ , i.e.,  $\int_0^h u_1(t, x_1, \eta) d\eta = c(t)$ . (Hint: Apply Gauß's theorem to a suitable subdomain of  $\Omega$ .)
- (d) Find a simple partial differential equation characterizing all solutions  $(u, p)$ , for which  $u$  is independent of  $x_1 \in \mathbb{R}$ .

**Exercise 2 (Schrödinger equation)** A free quantummechanical particle is governed by the complex-valued Schrödinger equation without potential. The wave function  $\Psi : \mathbb{R} \times \mathbb{R}^d \rightarrow \mathbb{C}$  has to satisfy

$$i \frac{\partial}{\partial t} \Psi = \Delta \Psi. \tag{SE}$$

- (a) We look for solutions  $\Psi$  of (SE) in the form

$$\Psi(t, x) = \gamma e^{-a(t)|x|^2 + b(t) \cdot x + c(t)} \quad \text{for } t \in \mathbb{R} \text{ and } x \in \mathbb{R}^d.$$

Derive ordinary differential equations

$$\dot{a} = f(a, b, c), \quad \dot{b} = g(a, b, c), \quad \text{and } \dot{c} = h(a, b, c),$$

for  $a : \mathbb{R} \rightarrow \mathbb{C}$ ,  $b : \mathbb{R} \rightarrow \mathbb{C}^d$ , and  $c : \mathbb{R} \rightarrow \mathbb{C}$  such that this  $\Psi$  solves (SE).

- (b) Write  $a(t) = \alpha(t) + i\beta(t)$  and show that  $\dot{a} = f(a, b, c)$  gives the ODE

$$\dot{\alpha} = r(\alpha, \beta) \quad \text{and} \quad \dot{\beta} = s(\alpha, \beta).$$

Conclude that  $\alpha(0) > 0$  implies that  $\alpha(t) > 0$  for all  $t$ . Show that all solutions can be given explicitly. (Hint: For  $a$  one obtains a complex Bernoulli equation.)

- (c) Provide the general solution for the ODE system in (a).
- (d) The function  $\rho(t, \cdot) = |\Psi(t, \cdot)|^2 \geq 0$  has the physical interpretation of a density distribution of the particle's position in  $\mathbb{R}^d$  at time  $t \in \mathbb{R}$ . Discuss the motion of the average position  $X(t) \in \mathbb{R}^d$  and the variance  $\sigma(t) > 0$  of the distribution.

please turn

**Exercise 3 (Separation of variables)** Consider the three PDEs

$$(i) \ u_{tt} = u_{xx}, \quad (ii) \ u_{tt} + u_{xx} = 0, \quad (iii) \ u_t = u_{xx}.$$

(a) For all three PDEs construct all possible solutions with *separated variables*, i.e. they have the form  $u(t, x) = V(t)W(x)$  for some real-valued, twice differentiable functions  $V$  and  $W$ . (Hint: After inserting the ansatz into the PDE, separate the variables to the different sides of “=”. Argue then that both sides must be constant.)

(b) For the heat equation (iii) find all solutions of (a) satisfying additionally  $W(0) = 0 = W(1)$ . From this, construct a solution of (iii) that additionally satisfies the boundary conditions  $u(t, 0) = u(t, 1) = 0$  and the initial condition  $u(0, x) = \frac{\pi^2}{17} \sin(\pi x) - 1.2 \sin(4\pi x)$ .

(c) For the wave equation (i) find all solution of (a) satisfying additionally  $W(0) = 0 = W(\ell)$ , where  $\ell > 0$  is given. Show that all these solutions satisfy  $u(t+2n\ell, x) = u(t, x)$  for  $n \in \mathbb{Z}$ .

**Exercise 4 [written] (Heat equation)** We want to investigate the heat equation on  $\Omega = (t, x) \in (0, \infty) \times \mathbb{R}^d$ :

$$u_t = \Delta u = \sum_{j=1}^d \frac{\partial^2}{\partial x_j^2} u. \tag{HE}$$

(a) The function  $H(t, x) = \frac{1}{(4\pi t)^{d/2}} \exp(-|x|^2/(4t))$  is called *heat kernel*. Show that  $u(t, x) = H(t, x)$  is a solution of (HE).

(b) For  $f \in C_c^0(\mathbb{R}^d) = \{f \in C^0(\mathbb{R}^d) \mid \text{sppt}(f) \text{ compact in } \mathbb{R}^d\}$  check that

$$u(t, x) = \int_{y \in \mathbb{R}^d} H(t, x-y) f(y) dy \tag{S}$$

is a solution of (HE), too. (Please, justify the interchange of limit processes  $\int$  and  $\partial$ !)

(c) For the solutions  $u$  in (S) and  $t > 0$  derive the estimates

$$\inf\{f(x) \mid x \in \mathbb{R}^d\} \leq \inf\{u(t, x) \mid x \in \mathbb{R}^d\} \leq \sup\{u(t, x) \mid x \in \mathbb{R}^d\} \leq \sup\{f(x) \mid x \in \mathbb{R}^d\}.$$

Moreover, establish the conservation of the total heat energy, i.e.  $\int_{\mathbb{R}^d} u(t, x) dx$  is independent of the time  $t$ .

(d) Establish for  $u$  in (S) and all  $x \in \mathbb{R}^d$  the limit  $\lim_{t \nearrow 0} u(t, x) = f(x)$ . (Hint: Substitute  $y = x + t^\alpha z$  for a suitable  $\alpha$ .)

**Turn in solutions for written exercises by Tuesday 19. April 2011, 12:00 h.**